



Celestial Bodies as Dark Matter Detectors

Tim Linden

Celestial Bodies vs. Direct Detection



Xenon-1T

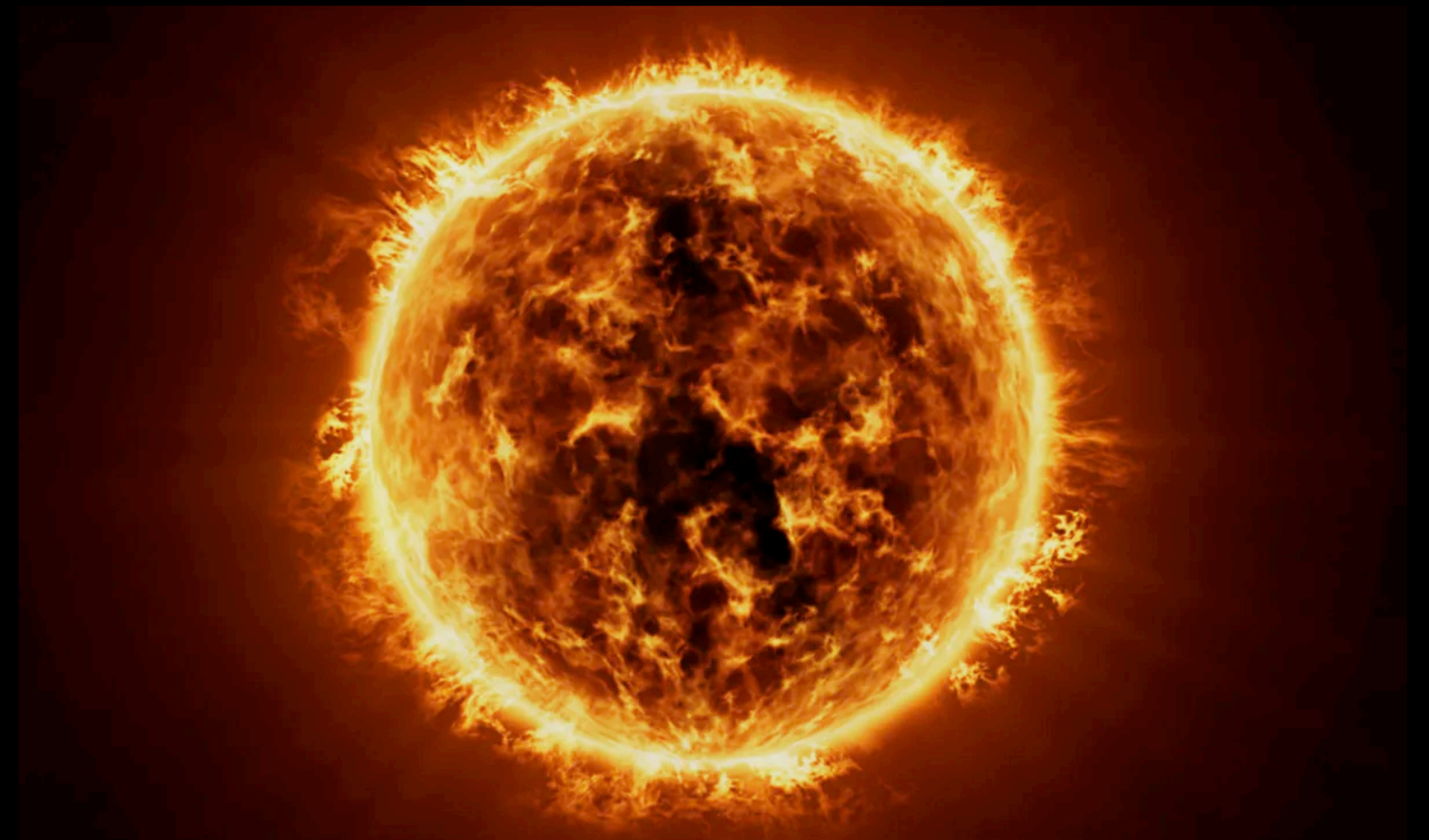
- 1000 kg
- 700 days

7×10^5 kg day

Celestial Body

- 3×10^{30} kg
- 2×10^{10} days

6×10^{40} kg day



Precision Physics is Possible

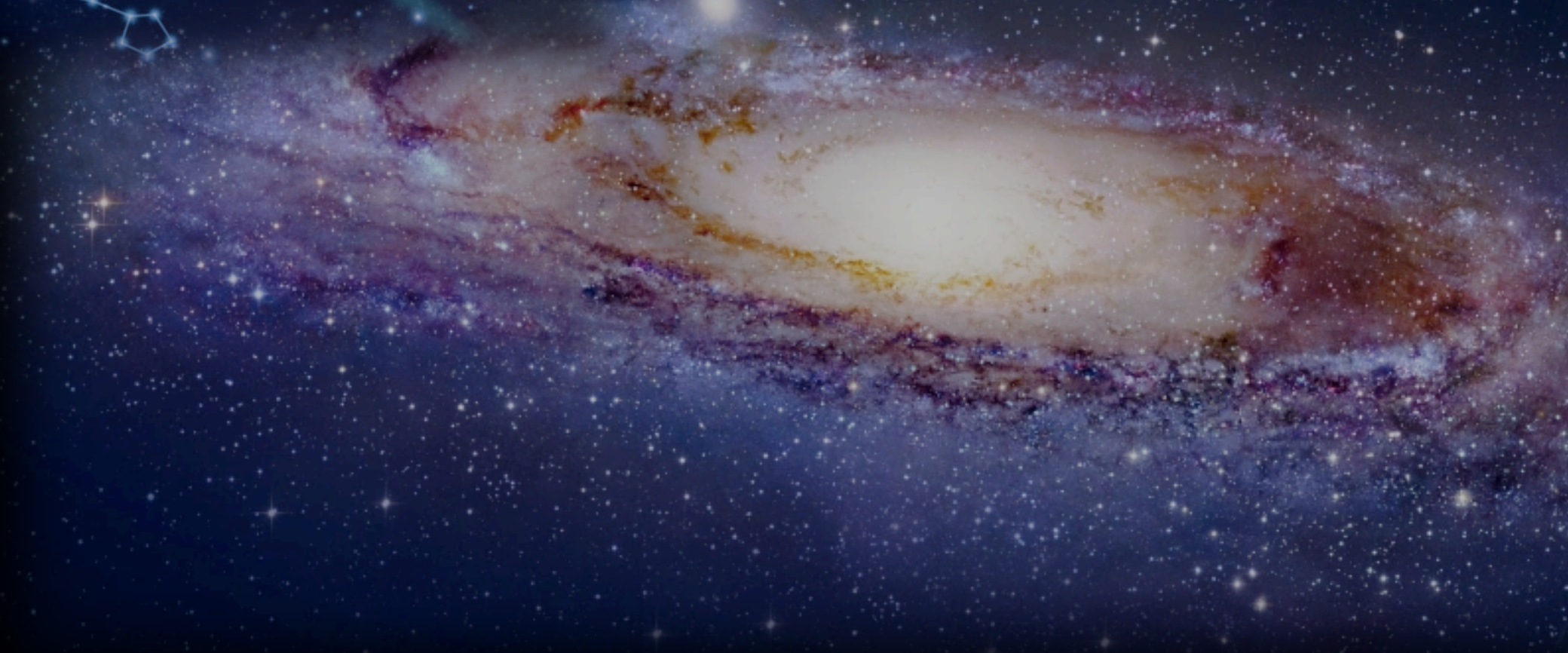
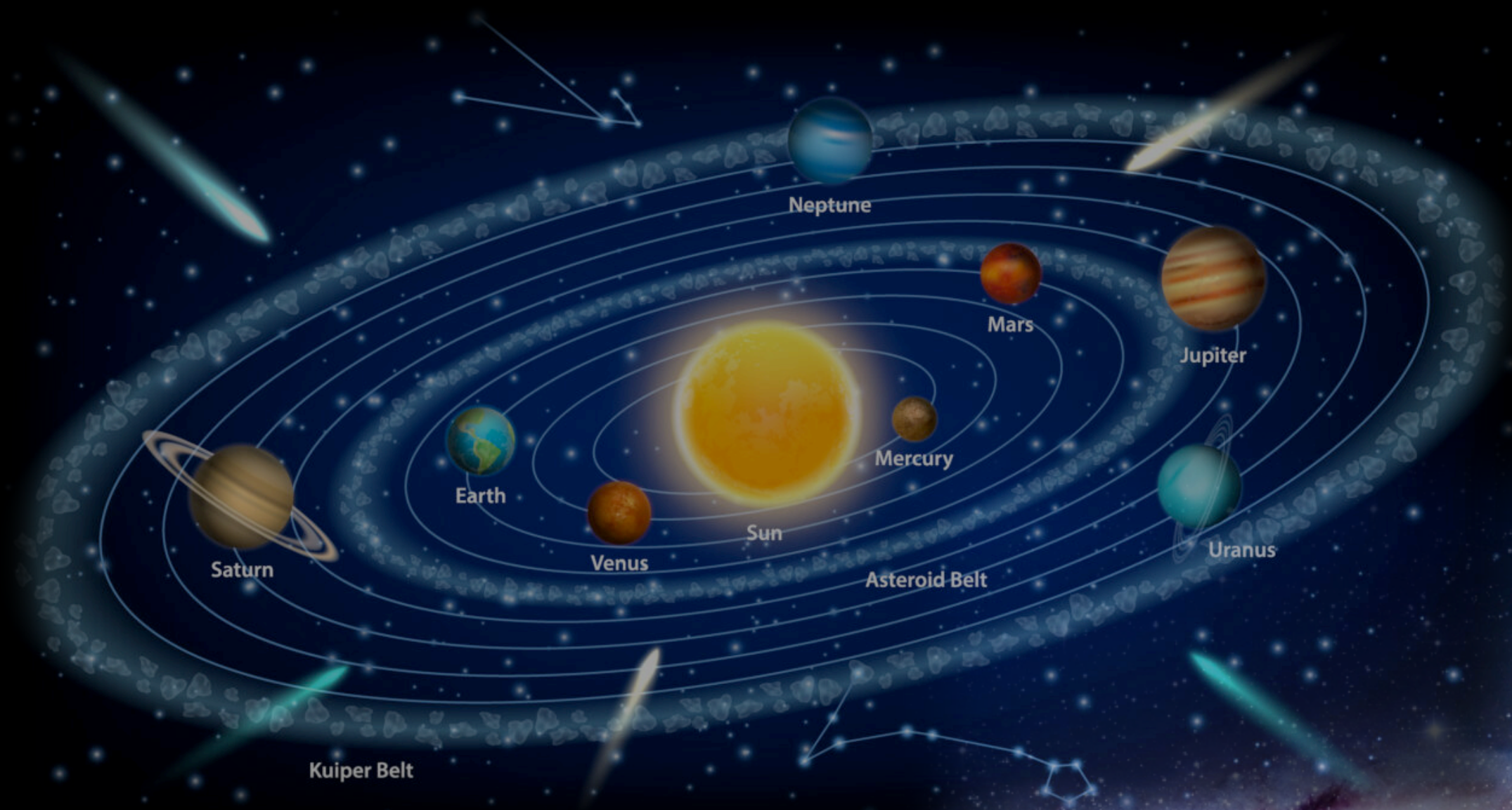
- Neutron star spin among the best measured quantities in physics.

PSR J1713+0747

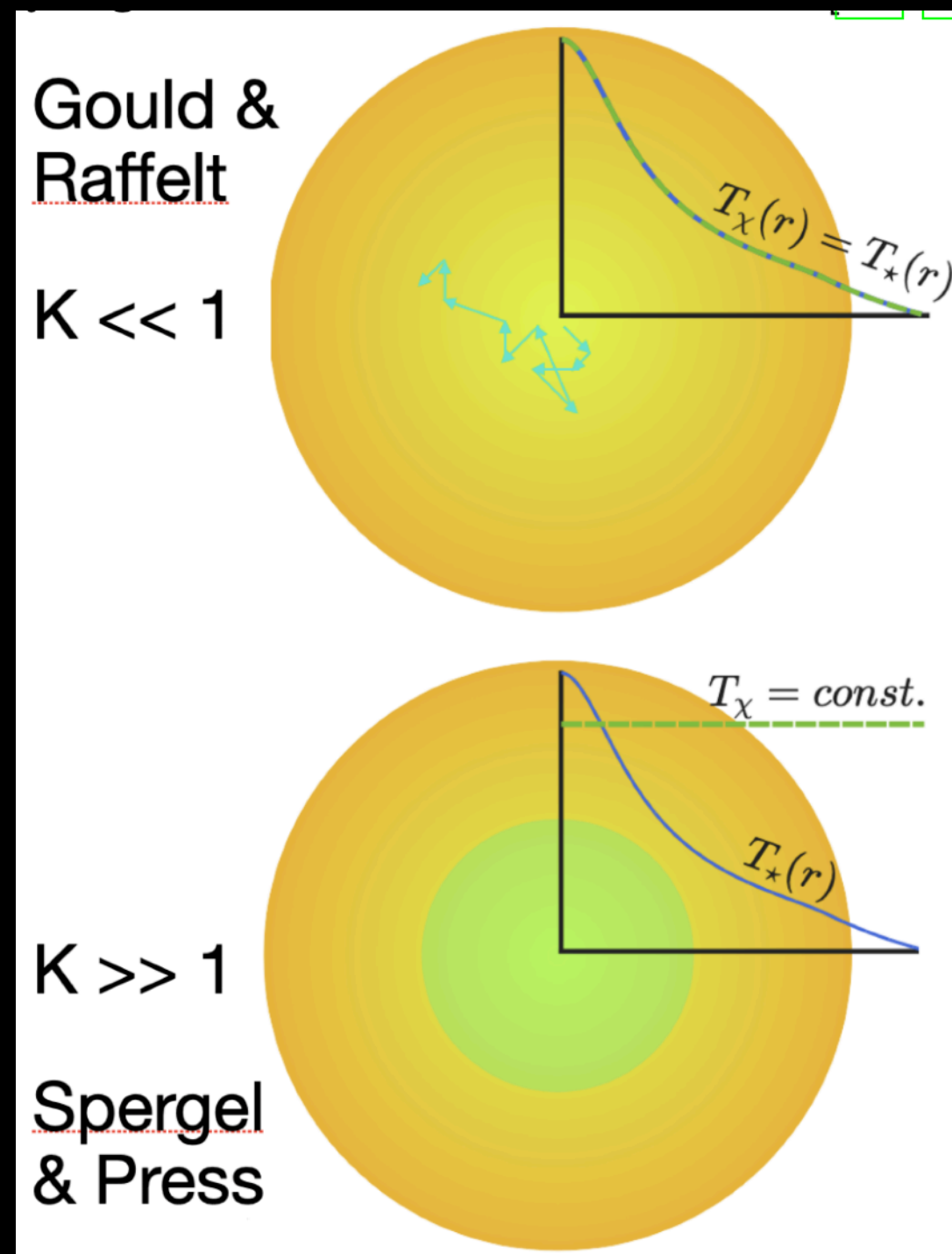
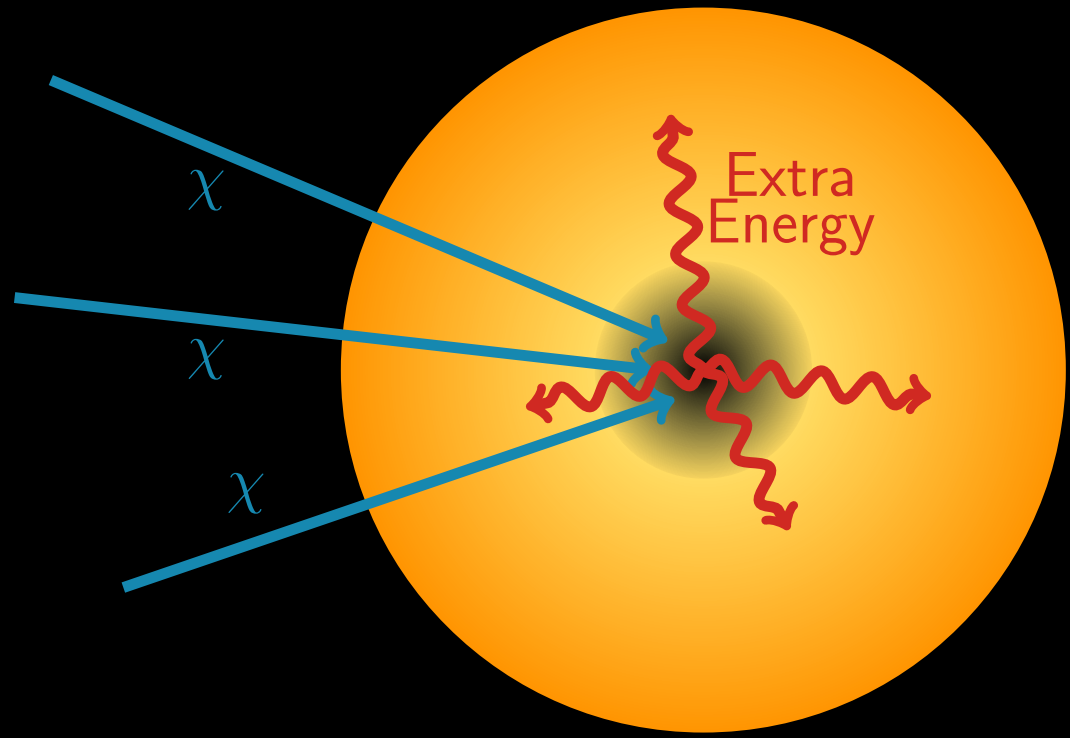
$$F = 218.8118437960826270 \pm 0.000000000000000988 \text{ s}^{-1}$$

$$F' = -4.083888637248 \pm 0.0000143324982645 \times 10^{-16} \text{ s}$$

A Multitude of Targets



A Multitude of Signatures



DM Heating

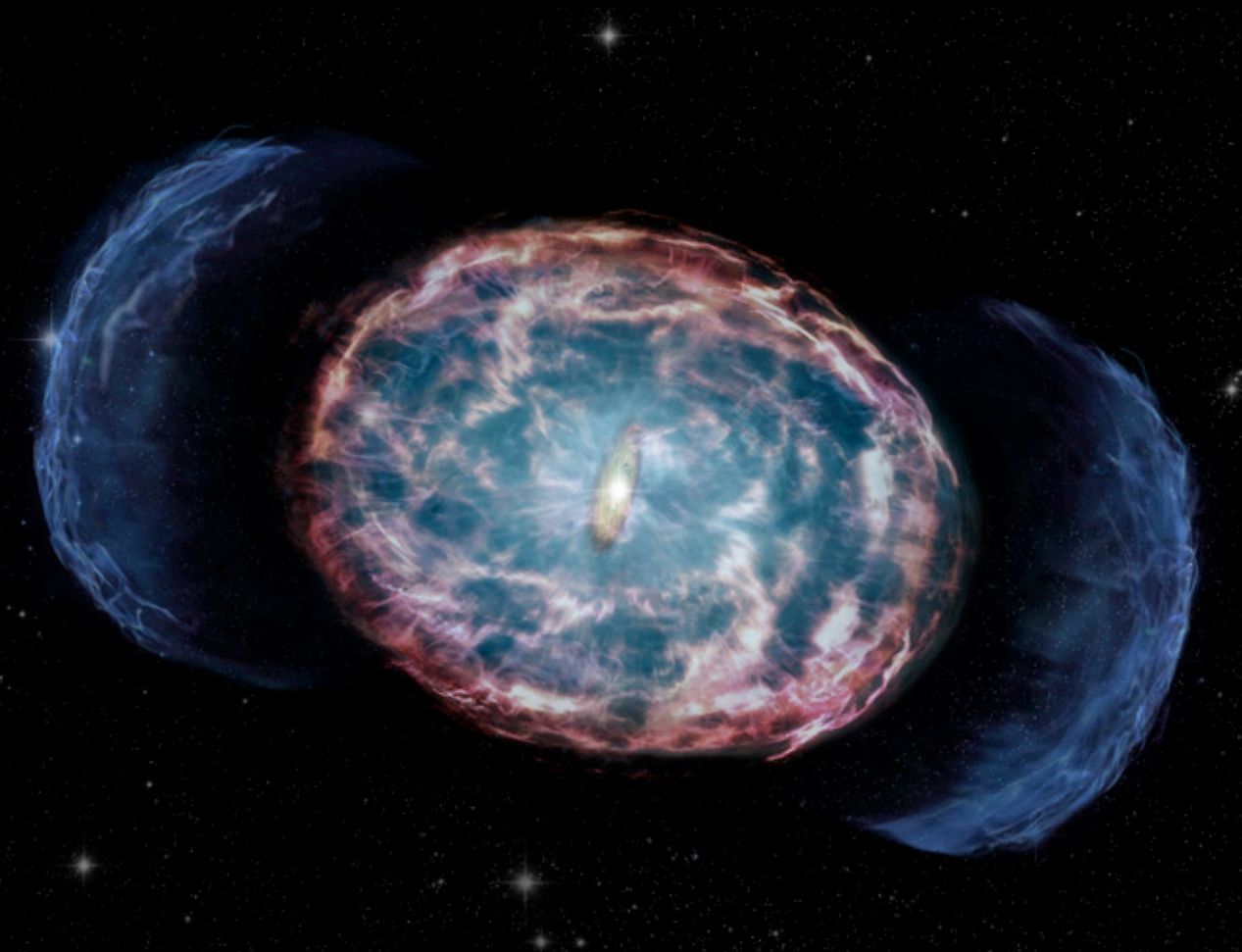
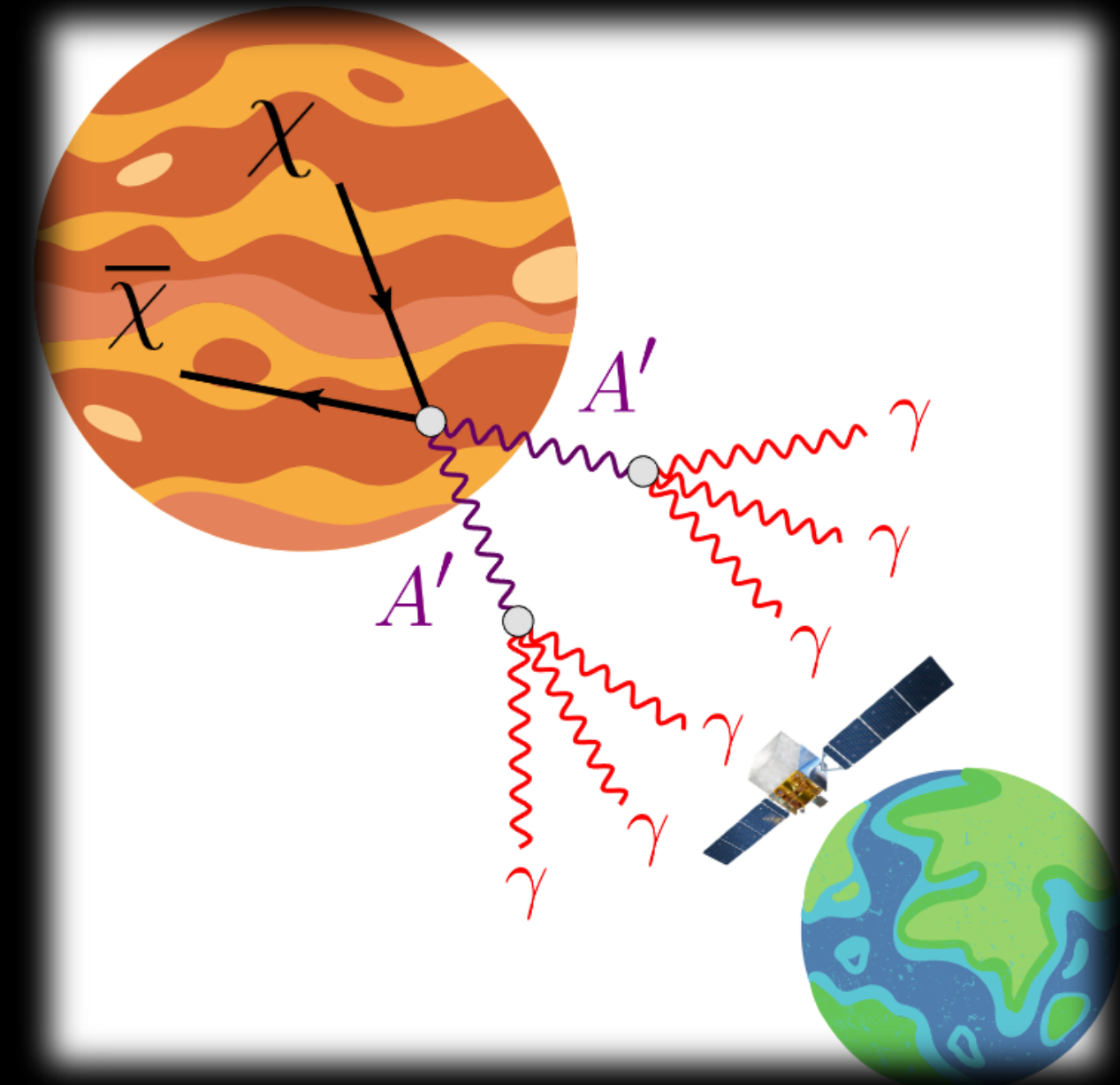
DM Signals

Energy

Transport

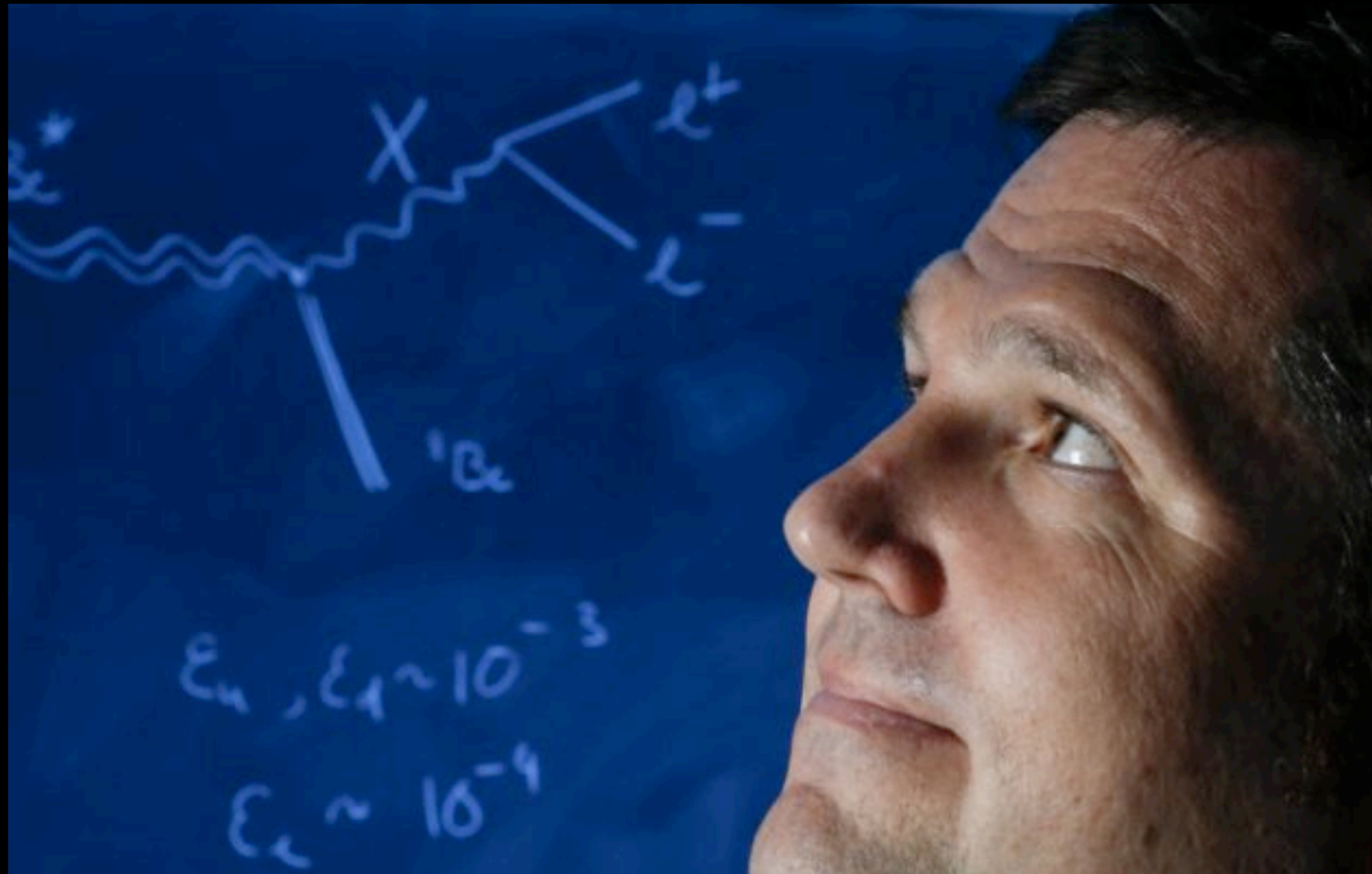
Explosive

Events










































A Multitude of Dark Matter Models

A Multitude of Dark Matter Models



A Cacophony Of Studies

DD	GeV Scale	Sun		     	Heating Signatures (Thermal Spectra)
	Asymmetric	Jupiter		 	Thermal Transport (Seismology)
	Self-Interacting	S-Cluster			Thermal Transport (Affect on Neutrino Flux)
	Light Mediator	G-Objects		  	Thermal Transport (Cooling Curves)
				 	Escaping Signals (DM Neutrinos)
no DD	<GeV Mass	GC Brown Dwarfs		   	Escaping Signals (Cosmic Rays/ γ -Rays)
	>TeV Mass	GC White Dwarfs		 	Structural Changes (Longer Lifespan of CBs)
	Spin-Dependent			   	Structural Changes (Destruction of CBs)
	Inelastic			   	Structural Changes (Spectral Changes)
	Velocity-Dependent	Dark Stars		   	Structural Changes (Population Level)

How to Do Science in the High-Risk High-Reward Regime

1.) Avoid Two-Miracle Studies

- **Standard model miracles cost half.**
- **Miracles can be correlated**

2.) Focus on observables

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3.) Attack the biggest uncertainty, and then move on.

- **Every individual study is individually unlikely.**

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$$A - M \geq 0$$

Possible detections

How to Do Science in the High-Risk High-Reward Regime

“Wear your **character** **theory** as lightly as a cap.”



A Few Recent Studies

1.) Stellar Heating at the Galactic Center (2311.16228; 2405.12267)

2.) SuperK Neutrino Searches in the Sun (2501.14864)

3.) Missing Pulsars at the Galactic Center (1405.1031)

4.) Unusual Supernovae (2211.00013)

Immortal Stars at the Galactic Center

See also: work on Dark Stars (e.g., Spolyar & Freese)

SLAC-PUB-17770

Dark Branches of Immortal Stars at the Galactic Center

Isabelle John,^{1,*} Rebecca K. Leane,^{2,3,†} and Tim Linden^{1,‡}

¹*Stockholm University and The Oskar Klein Centre for Cosmoparticle Physics, Alba Nova, 10691 Stockholm, Sweden*

²*Particle Theory Group, SLAC National Accelerator Laboratory, Stanford, CA 94035, USA*

³*Kavli Institute for Particle Astrophysics and Cosmology, Stanford University, Stanford, CA 94035, USA*

We show that stars in the inner parsec of the Milky Way can be significantly affected by dark matter annihilation, producing population-level effects that are visible in a Hertzsprung-Russell (H-R) diagram.

We establish the dark HR diagram, where stars lie on a new stable *dark main sequence* with lower luminosities, but lower temperatures, than the standard main sequence.

The dark main sequence is continuously replenished, granting these stars immortality.

Upcoming telescopes could detect the dark main sequence.

Dark Matter Scattering Constraints from Observations of Stars Surrounding Sgr A*

Isabelle John,^{1,*} Rebecca K. Leane,^{2,3,†} and Tim Linden^{1,‡}

¹*Stockholm University and The Oskar Klein Centre for Cosmoparticle Physics, Alba Nova, 10691 Stockholm, Sweden*
²*Particle Theory Group, SLAC National Accelerator Laboratory, Stanford, CA 94035, USA*
³*Kavli Institute for Particle Astrophysics and Cosmology, Stanford University, Stanford, CA 94035, USA*

High resolution infrared data has revealed several young stars in close proximity to the Galactic Center. We examine scenarios where dark matter annihilation may encounter extremely high dark matter densities, and then annihilate, producing population-level effects that are visible in a Hertzsprung-Russell (H-R) diagram. We establish the dark HR diagram, where stars lie on a new stable *dark main sequence* with lower luminosities, but lower temperatures, than the standard main sequence. The dark main sequence is continuously replenished, granting these stars immortality. Upcoming telescopes could detect the dark main sequence.

Dark stars at the Galactic centre – the main sequence

Pat Scott^{1*}, Malcolm Fairbairn^{2,3*} and Joakim Edsjö^{1*}

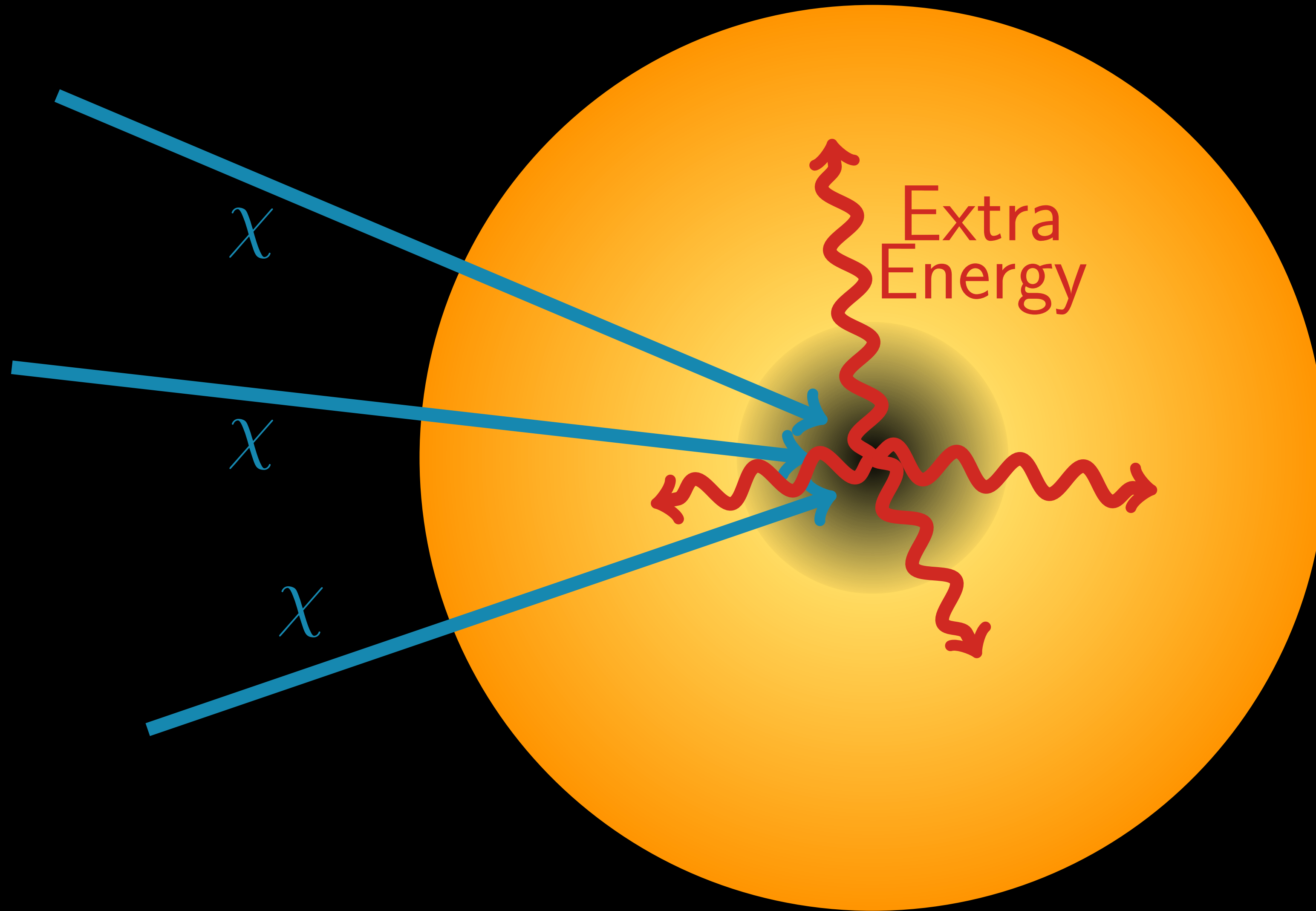
¹*Cosmology, Particle Astrophysics and String Theory, Department of Physics, Stockholm University & Oskar Klein Centre for Cosmoparticle Physics, AlbaNova University Centre, SE-106 91 Stockholm, Sweden*
²*Theory Division, CERN, CH-1211, Geneva 23, Switzerland*
³*Physics, Kings College London, Strand, London WC2R 2LS, UK*

2008 November 19. Submitted 2008 November 17; in original form 2008 October 5.

ABSTRACT

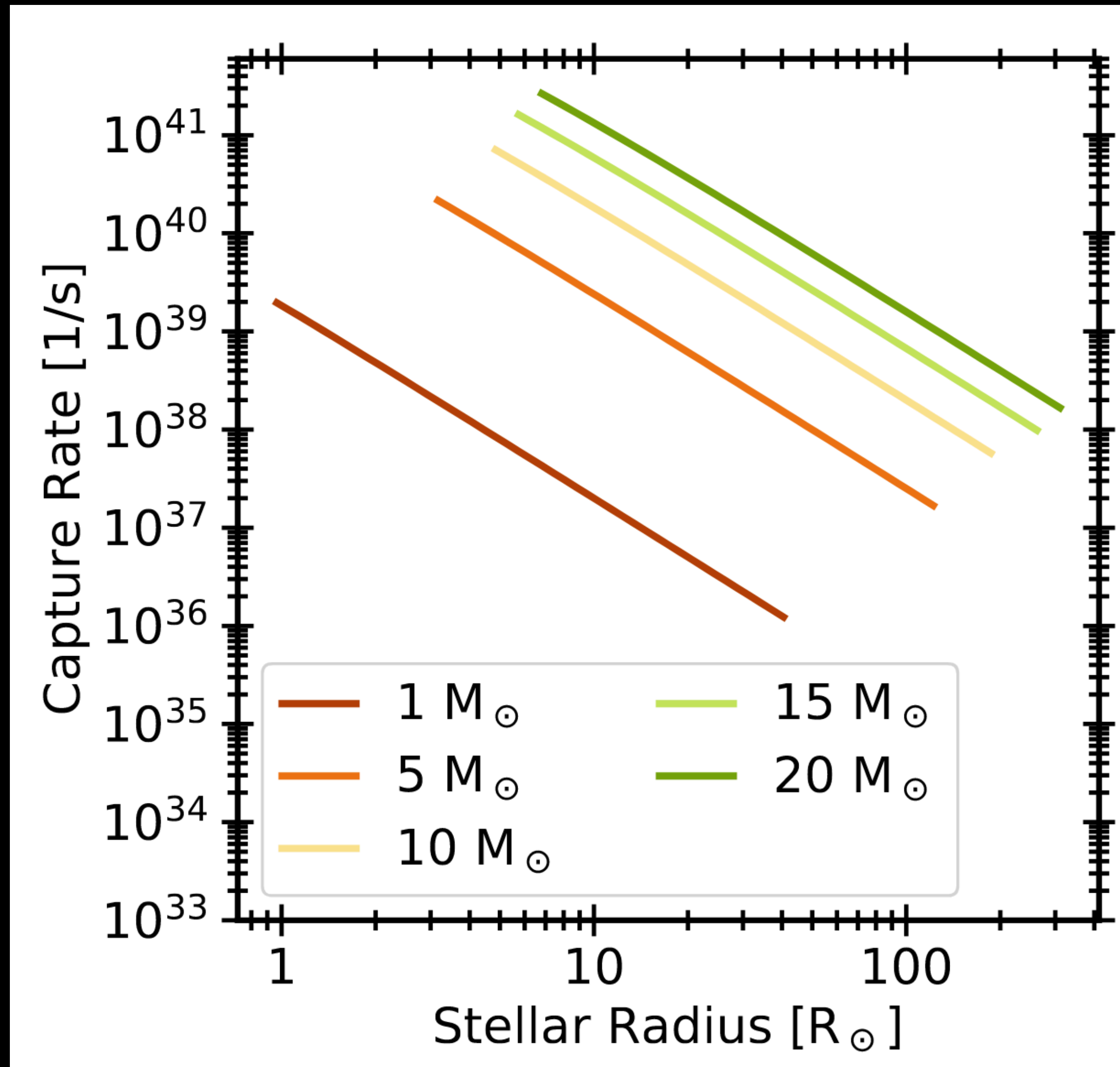
In regions of very high dark matter density such as the Galactic center, the annihilation of WIMP dark matter by stars has the potential to significantly influence the dark stellar evolution code. We describe the dark stellar evolution code and its influence on the main sequence evolution of stars with masses of 0.3–2.0 M_⊙. The results show that dark matter annihilation can significantly influence the main sequence evolution of stars, leading to the formation of a new stable *dark main sequence* with lower luminosities, but lower temperatures, than the standard main sequence. The dark main sequence is continuously replenished, granting these stars immortality. Upcoming telescopes could detect the dark main sequence.

Immortal Stars at the Galactic Center



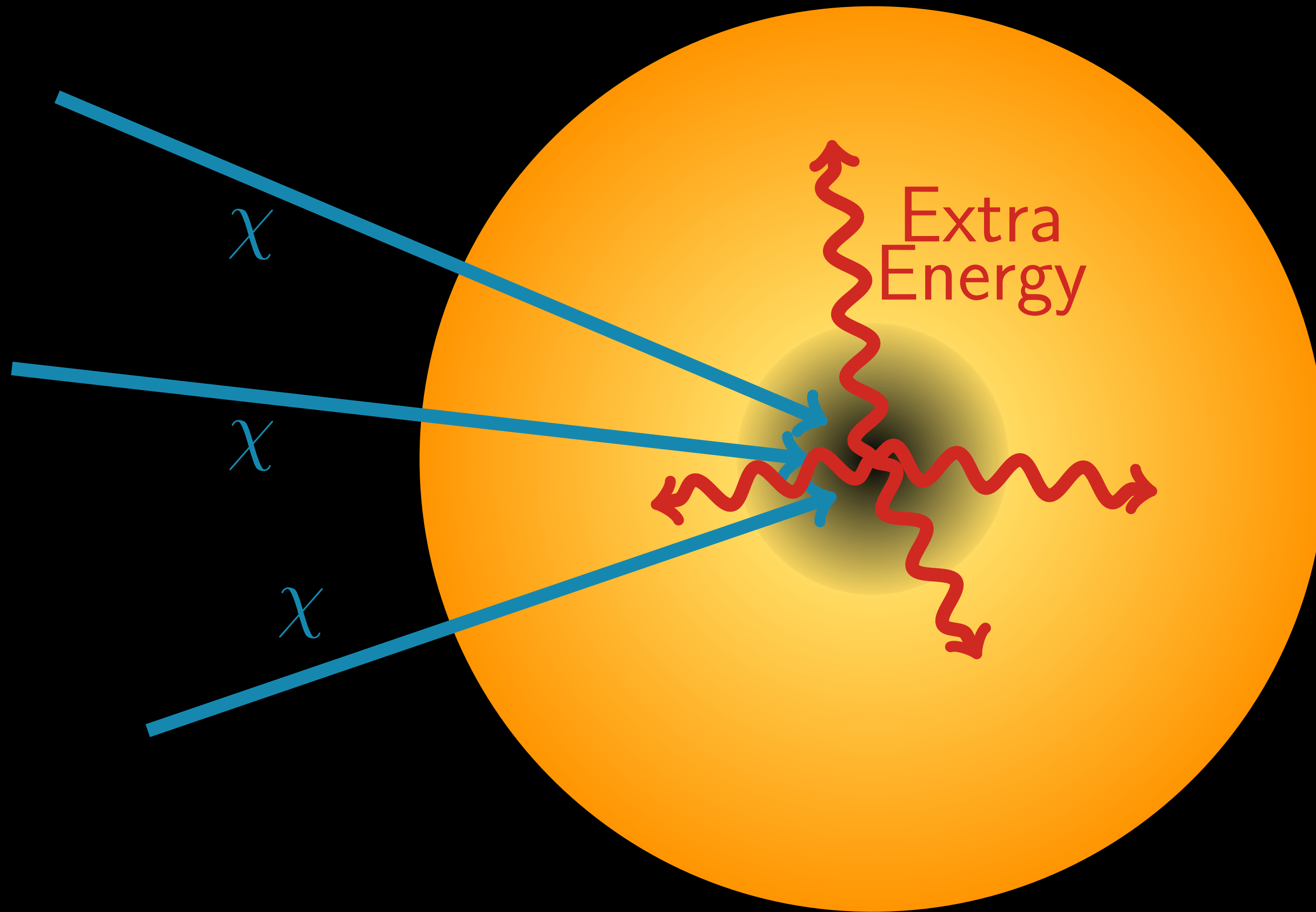
- Dark Matter annihilation provides an additional power source that heats the star.
- The star maintains equilibrium - it expands if too much power is injected.

Immortal Stars at the Galactic Center



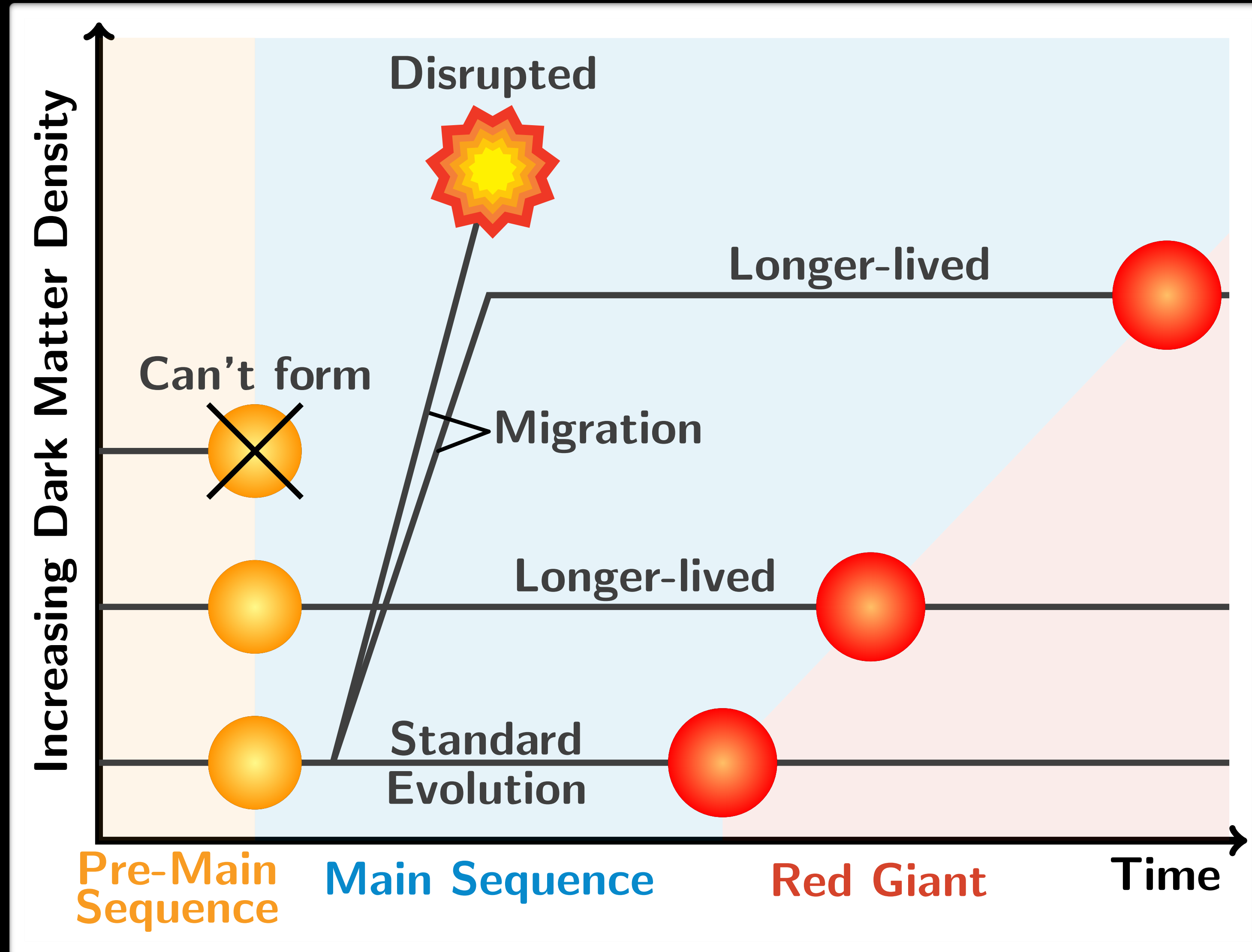
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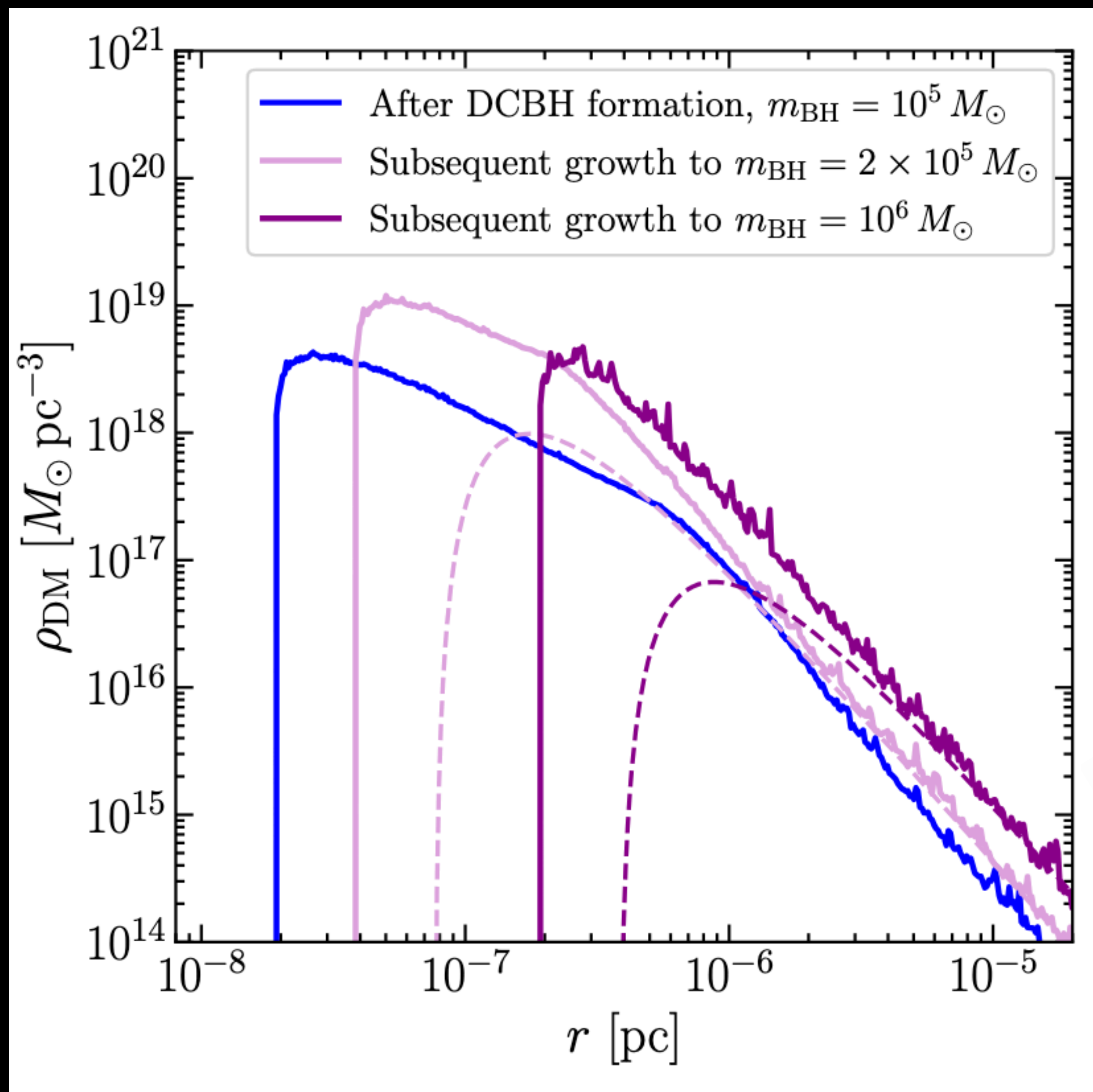
- “Miracle” - Very high dark matter densities at the galactic center.
- **Standard WIMP DM**
- **Standard (though relatively high) dark matter density (low mass WIMPs)**

Immortal Stars at the Galactic Center



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Immortal Stars at the Galactic Center



Dark Matter Mounds: towards a realistic description of dark matter overdensities around black holes

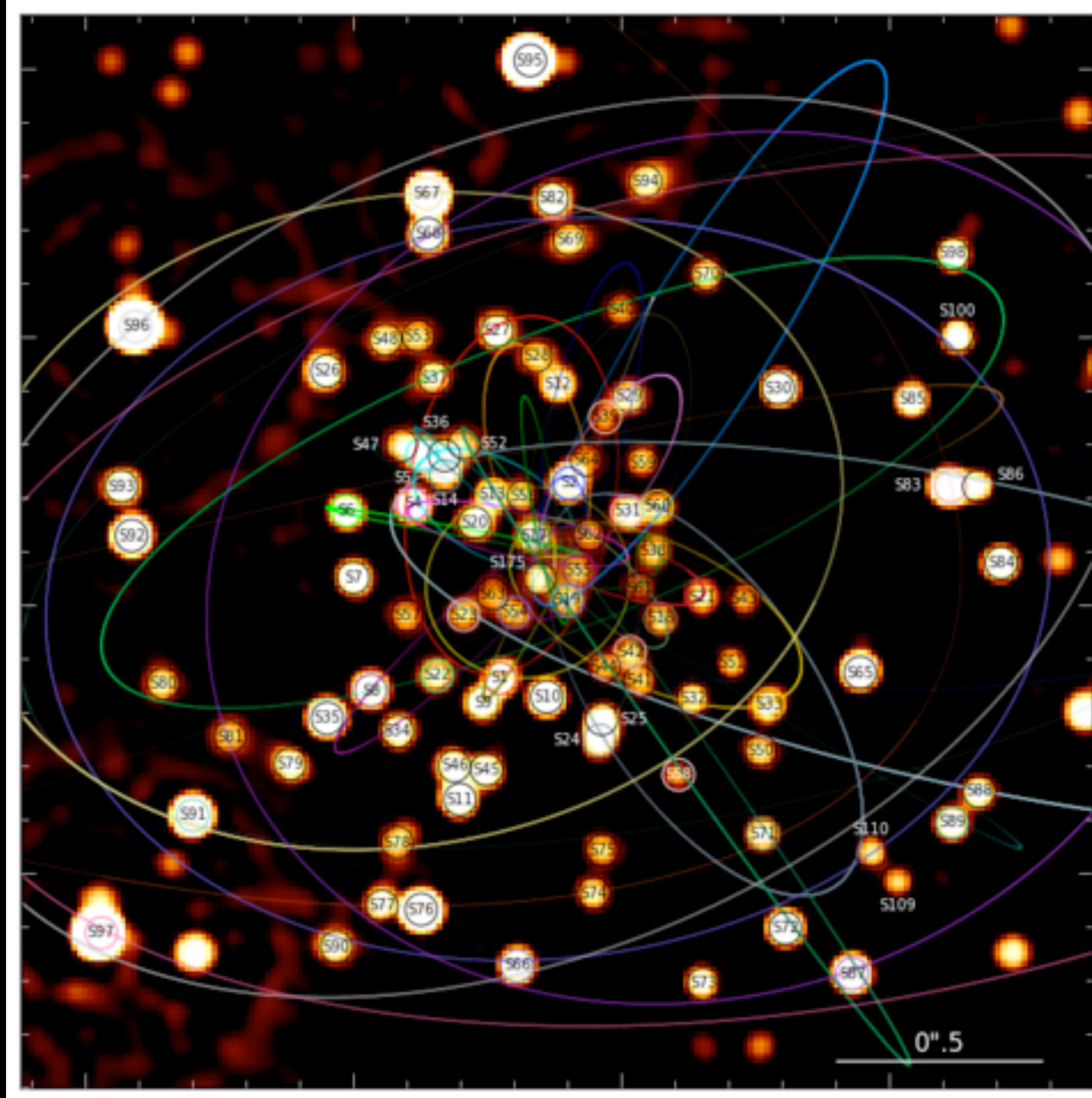
Gianfranco Bertone¹, A. Renske A. C. Wierda^{1,5}, Daniele Gaggero³,
Bradley J. Kavanagh⁴, Marta Volonteri^{1,5} and Naoki Yoshida^{6,7}

¹Gravitation Astroparticle Physics Amsterdam (GRAPPA),
University of Amsterdam, Amsterdam, 1098 XH, Netherlands
²Department of Physics, KTH Royal Institute of Technology,
The Oskar Klein Centre, AlbaNova, SE-106 91 Stockholm, Sweden
³INFN Sezione di Pisa, Polo Fibonacchi, Largo B. Pontecorvo 3, 56127 Pisa, Italy
⁴Instituto de Física de Cantabria (IFCA, UC-CSIC),
Avenida de Los Castros s/n, 39005 Santander, Spain
⁵Institut d'Astrophysique de Paris, Sorbonne Université,
CNRS, UMR 7095, 98 bis bd Arago, 75014 Paris, France
⁶Department of Physics, The University of Tokyo, Chiba 277-8583, Japan
⁷Kavli Institute for Advanced Study, The University of Tokyo, Chiba 277-8583, Japan

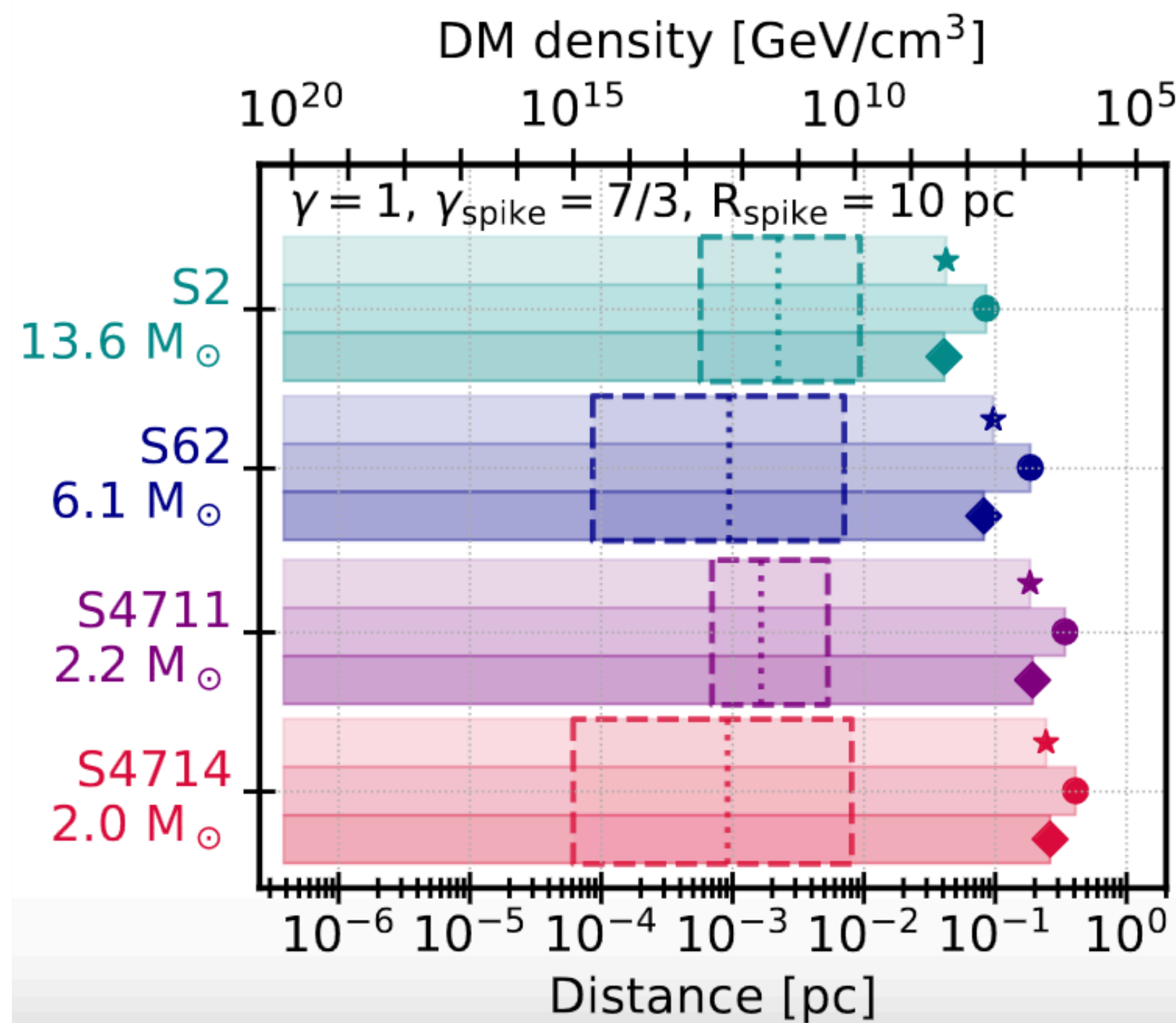
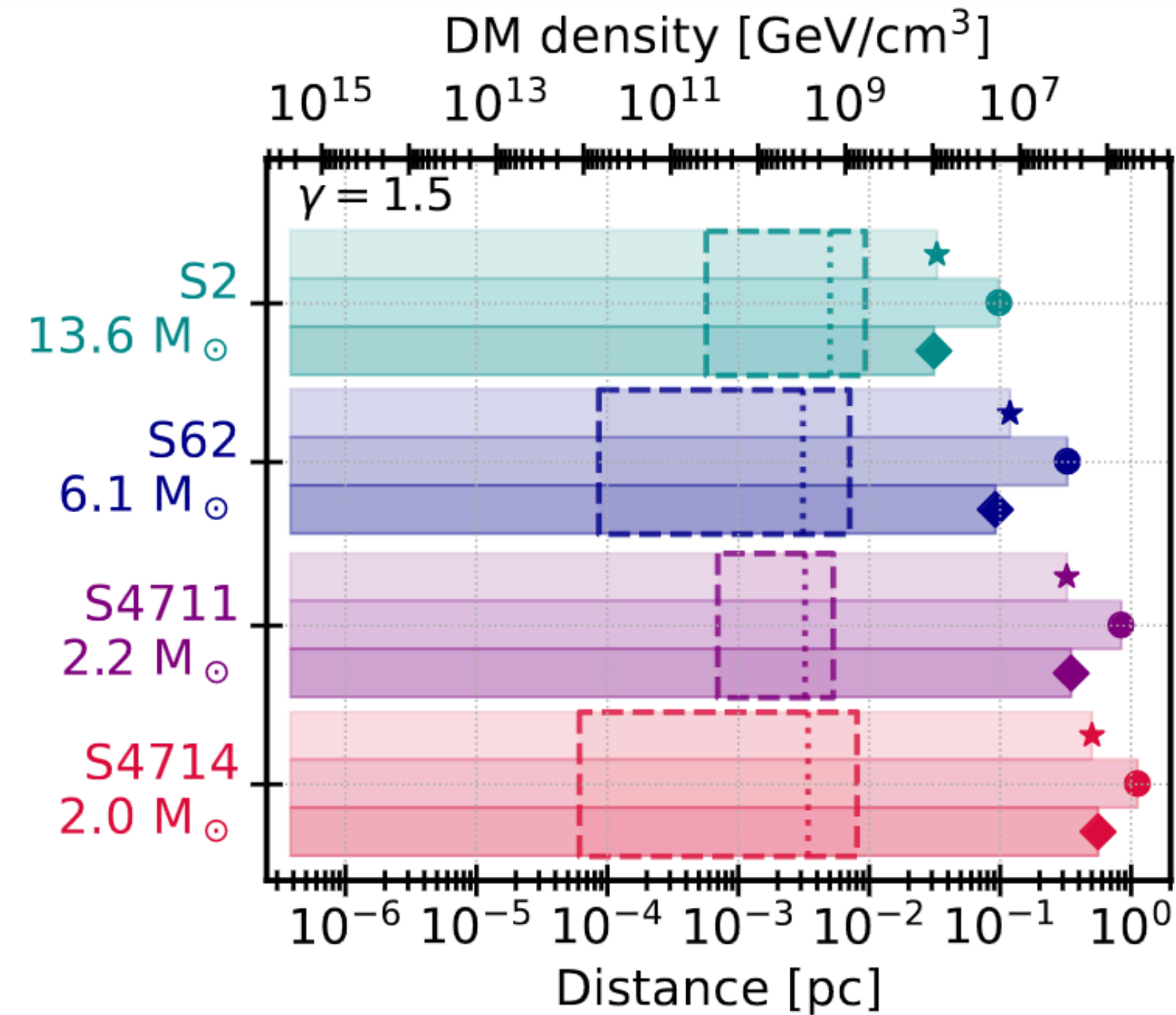
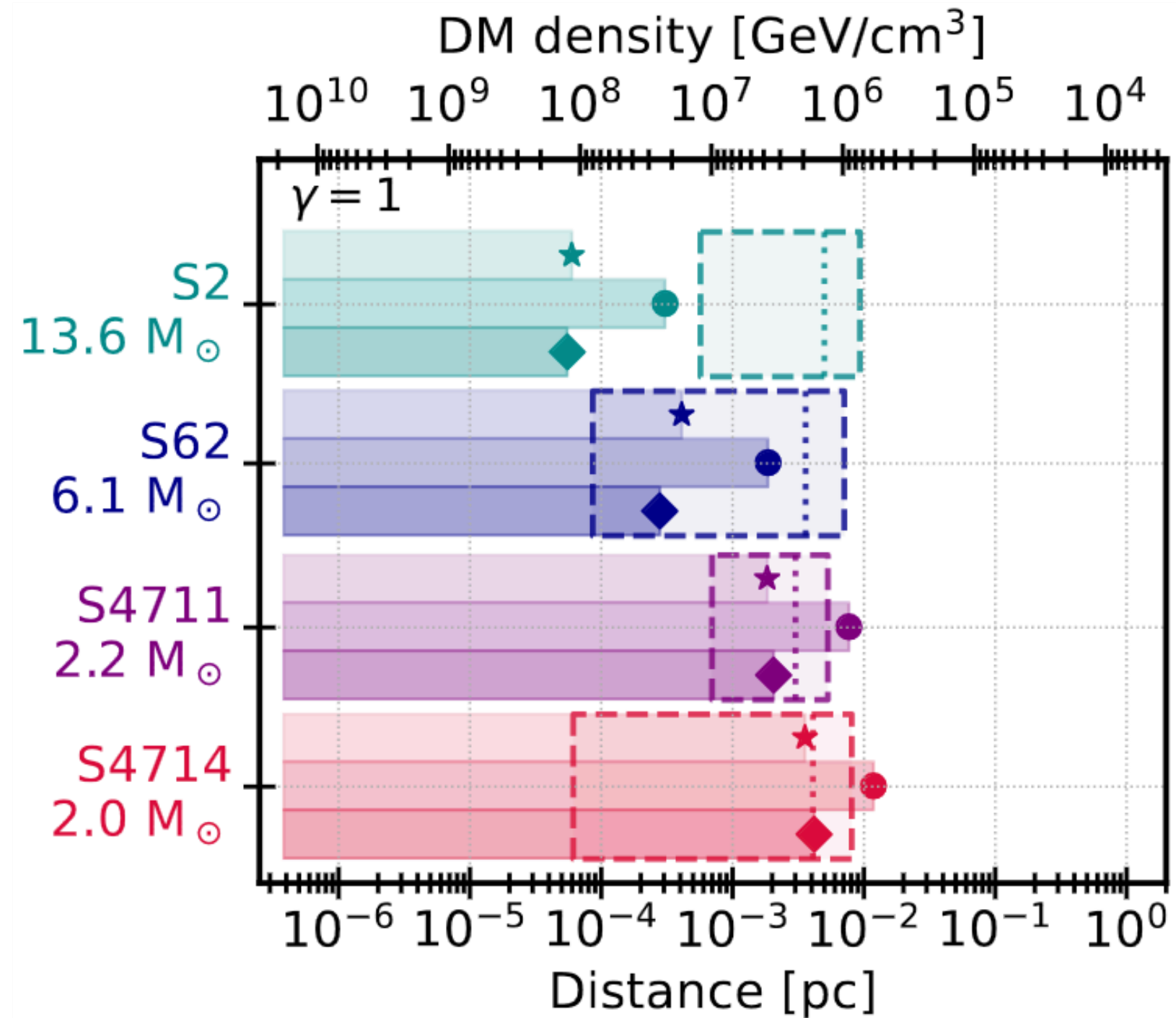
Dark matter overdensities around black holes can be searched for by looking at the characteristic imprint they leave on the gravitational waveform of binary black hole mergers. Current theoretical predictions of the density profile of dark matter cusps at the center of galactic halos are based on highly idealised formation scenarios, in which black holes are assumed to grow adiabatically from an infinitesimal seed mass, compressing dark matter into very high densities, but they fail to capture the dark matter distribution that would develop in more realistic scenarios where black holes form from the collapse of supermassive stars, or where the dark matter density is produced by indirect detection of supermassive stars, or where the dark matter forms shallower 'mounds' instead of cusps. We present here a realistic description of dark matter overdensities around black holes.






PR 2024

Immortal Stars at the Galactic Center

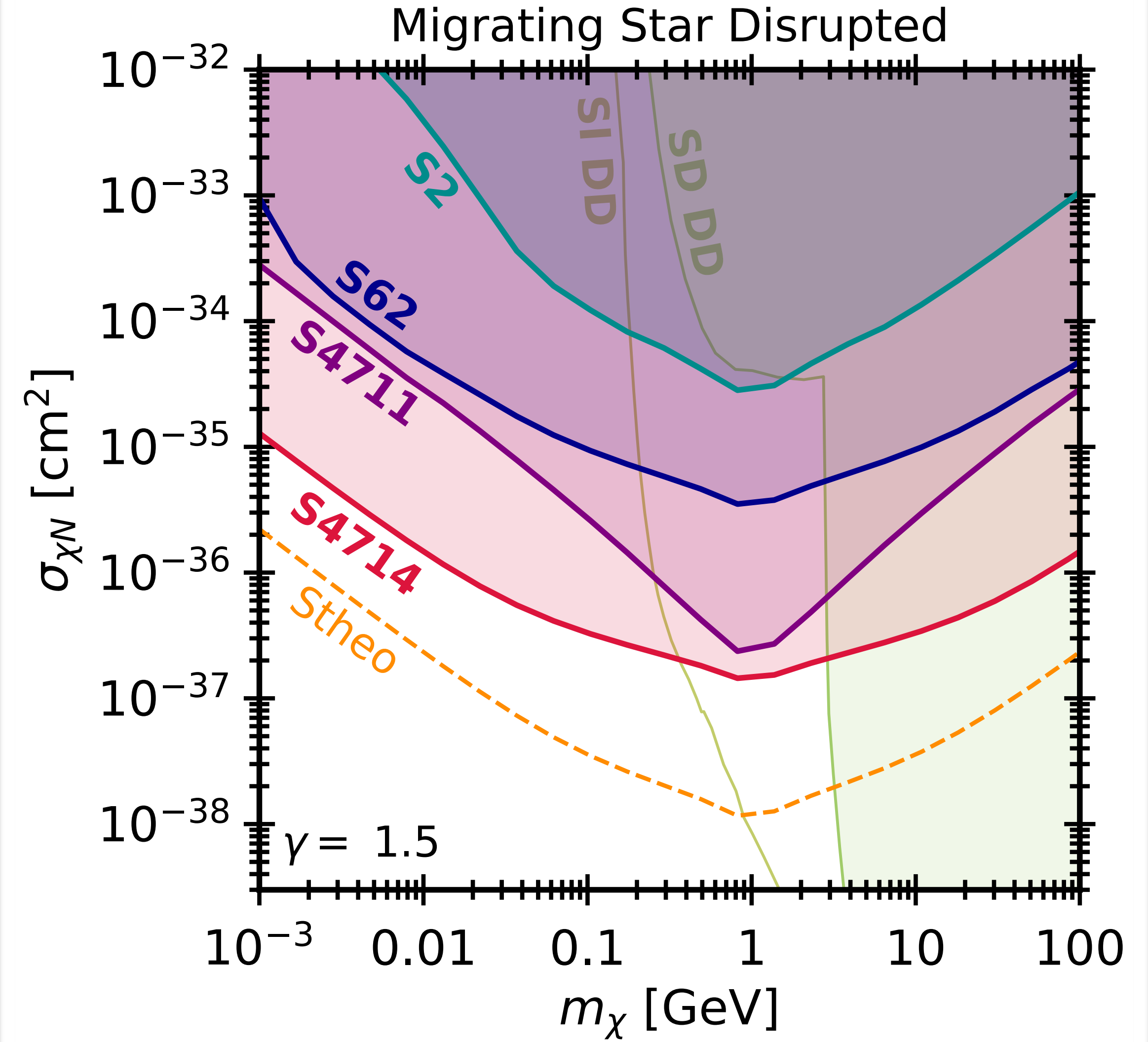
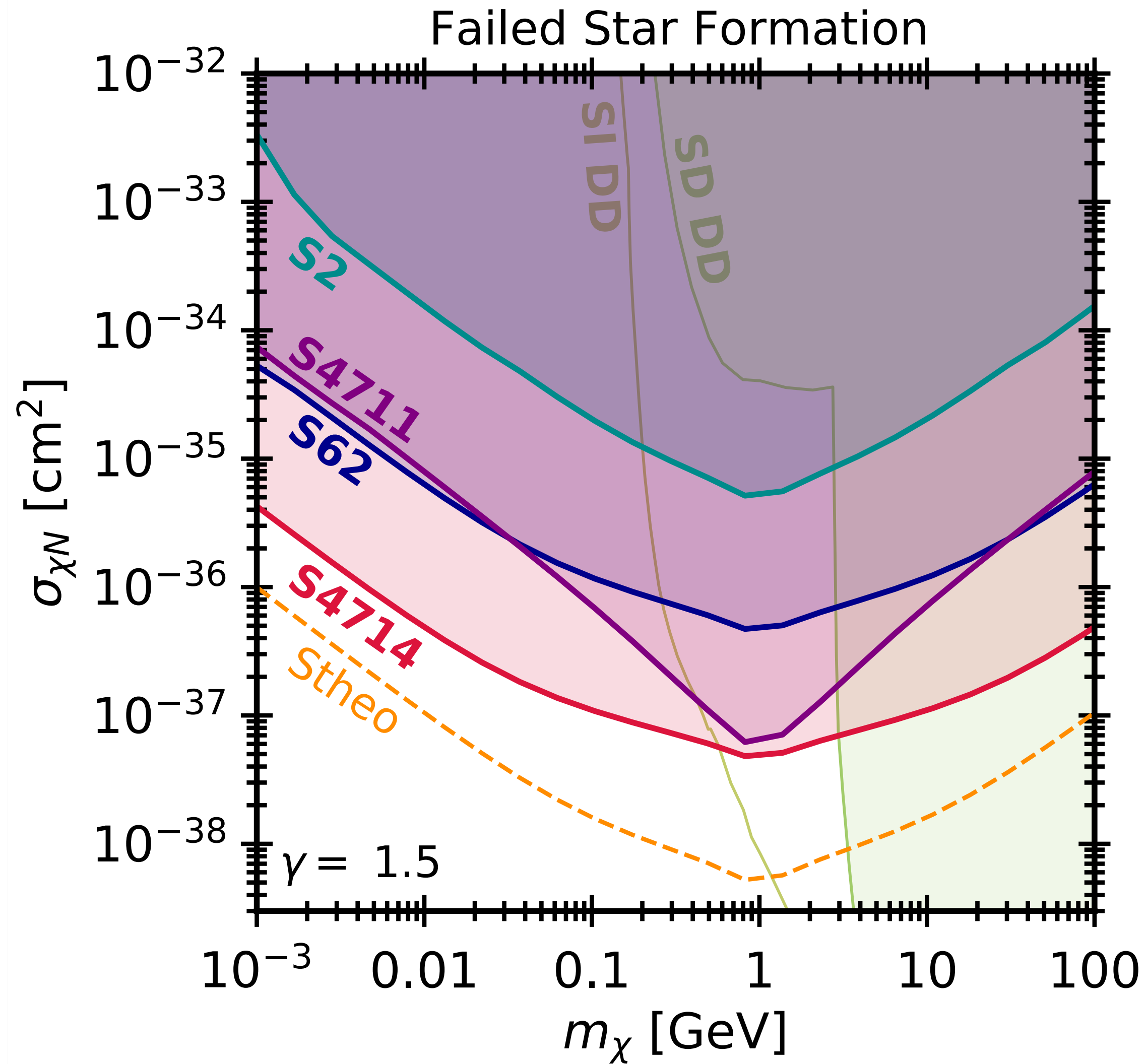


**Many stars
within 0.016 pc**



-  Distance covered by stellar orbit
-  Average dark matter density encountered along orbit
-  Dark matter annihilation power exceeds nuclear fusion power
-  Dark matter annihilation power prevents star from forming
-  Dark matter annihilation power disrupts star after migration

Immortal Stars at the Galactic Center



Immortal Stars at the Galactic Center



Origin not well understood: in situ formation or migration?

Paradox of Youth: Spectroscopically old but bright as young stars

Conundrum of Old Age: Lack of old stars



Top-heavy initial mass function: large abundance of massive stars

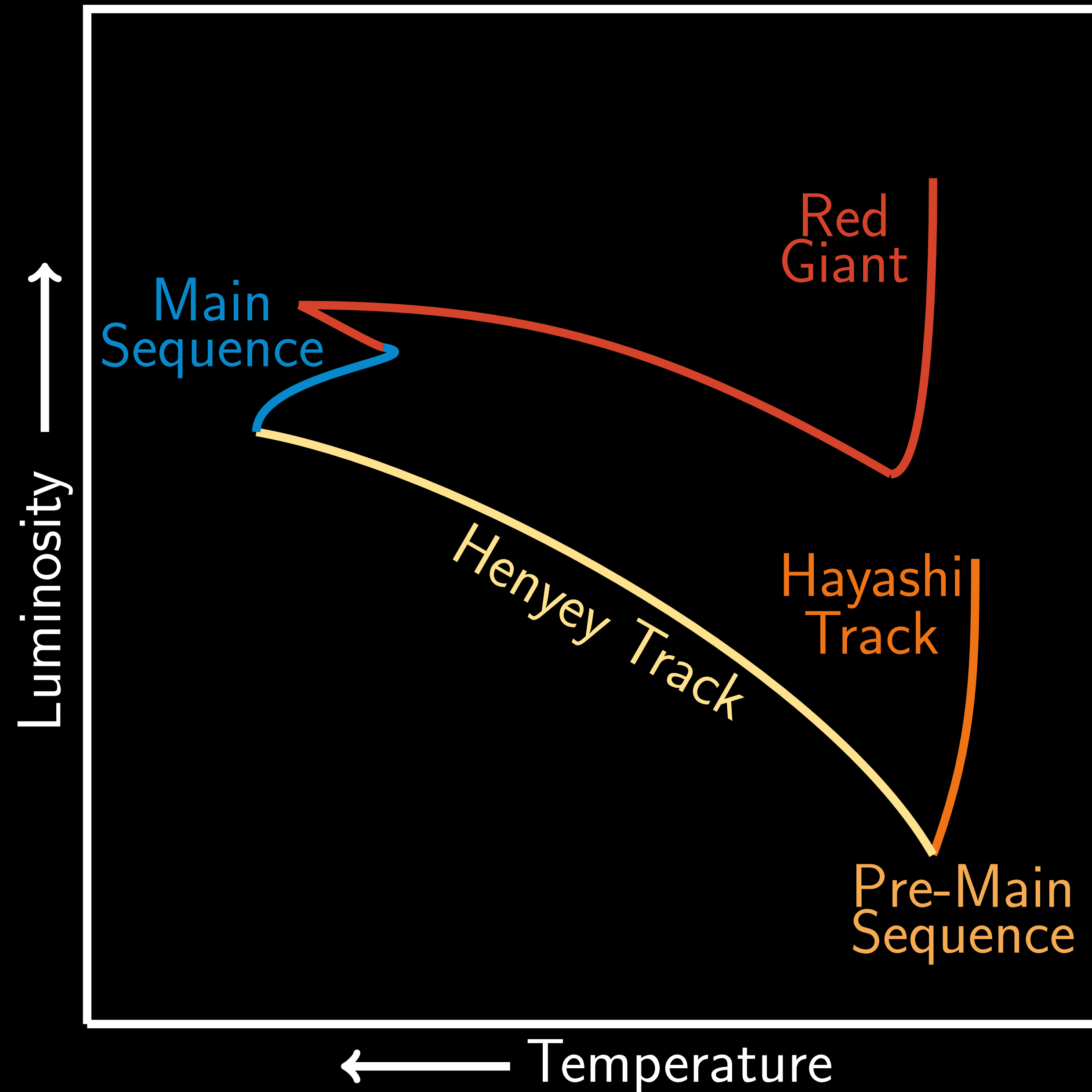
Massive Stars in Interacting Binaries
ASP Conference Series, Vol. 367, 2007
N. St-Louis & A.F.J. Moffat

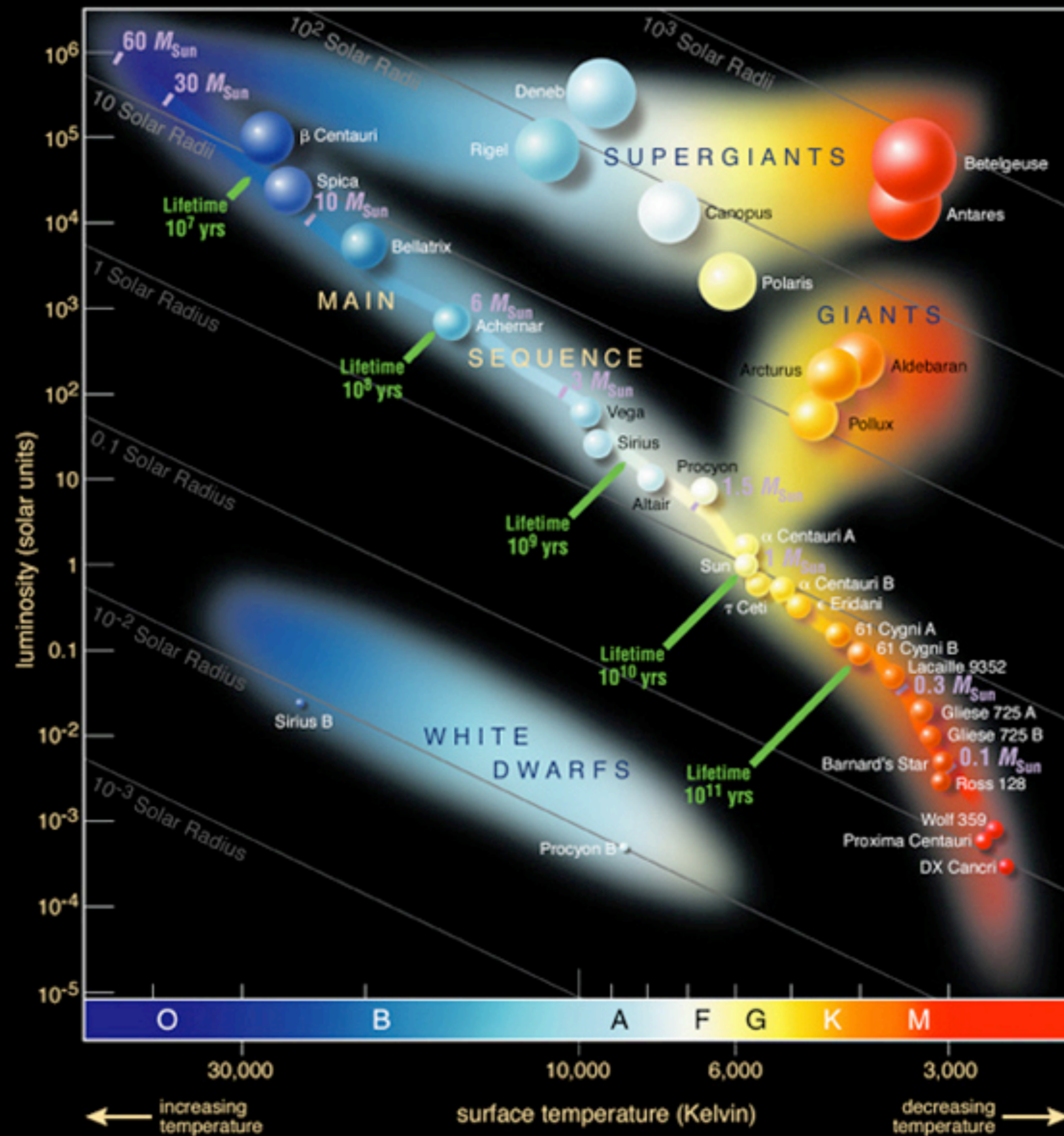
Massive Young Stars in the Vicinity of our Galaxy's
Supermassive Black Hole: A Paradox of Youth

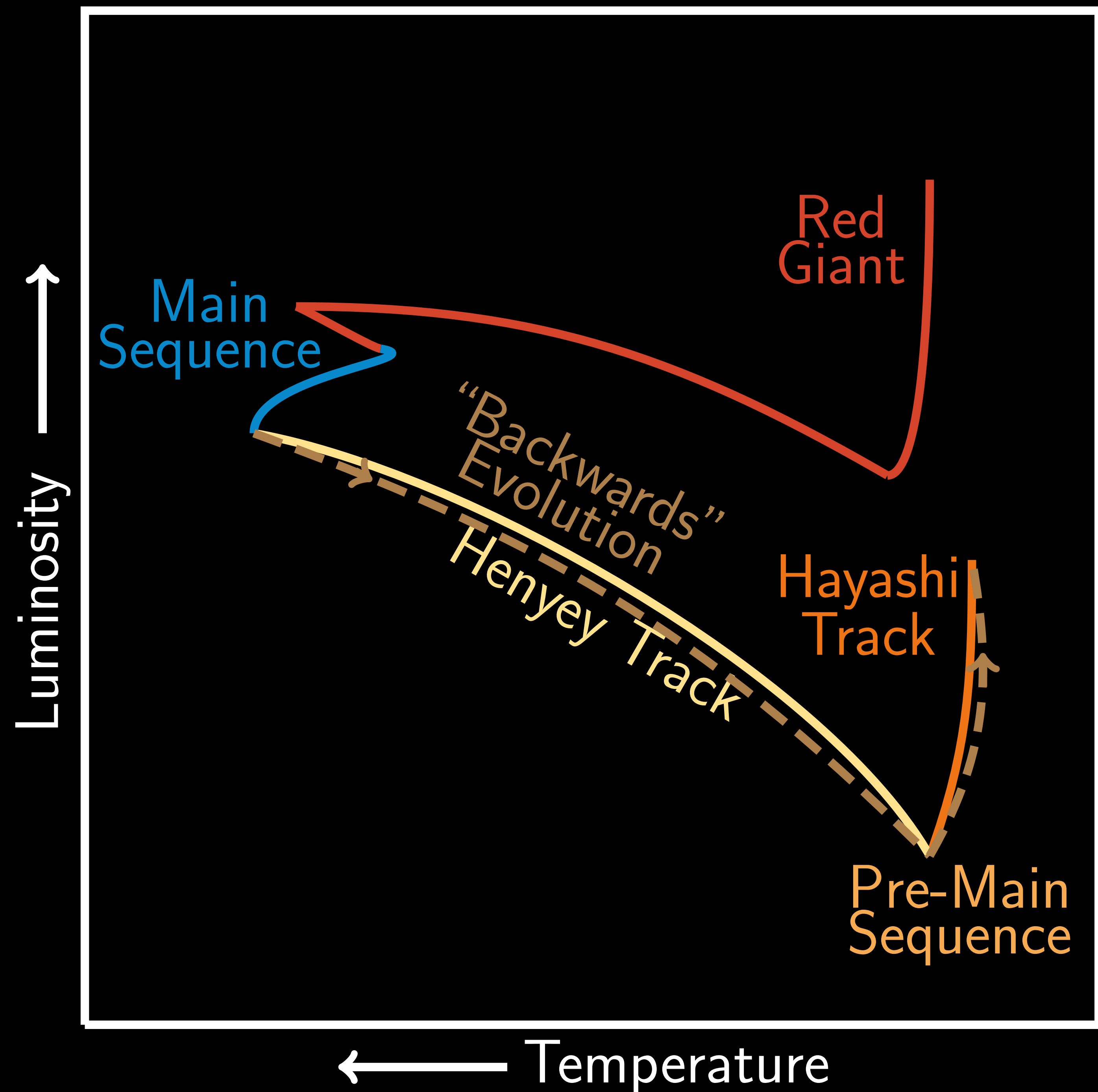
A. M. Ghez
Department of Physics & Astronomy and IGPP
University of California Los Angeles
12517 USA

of our Galaxy hav
contrib

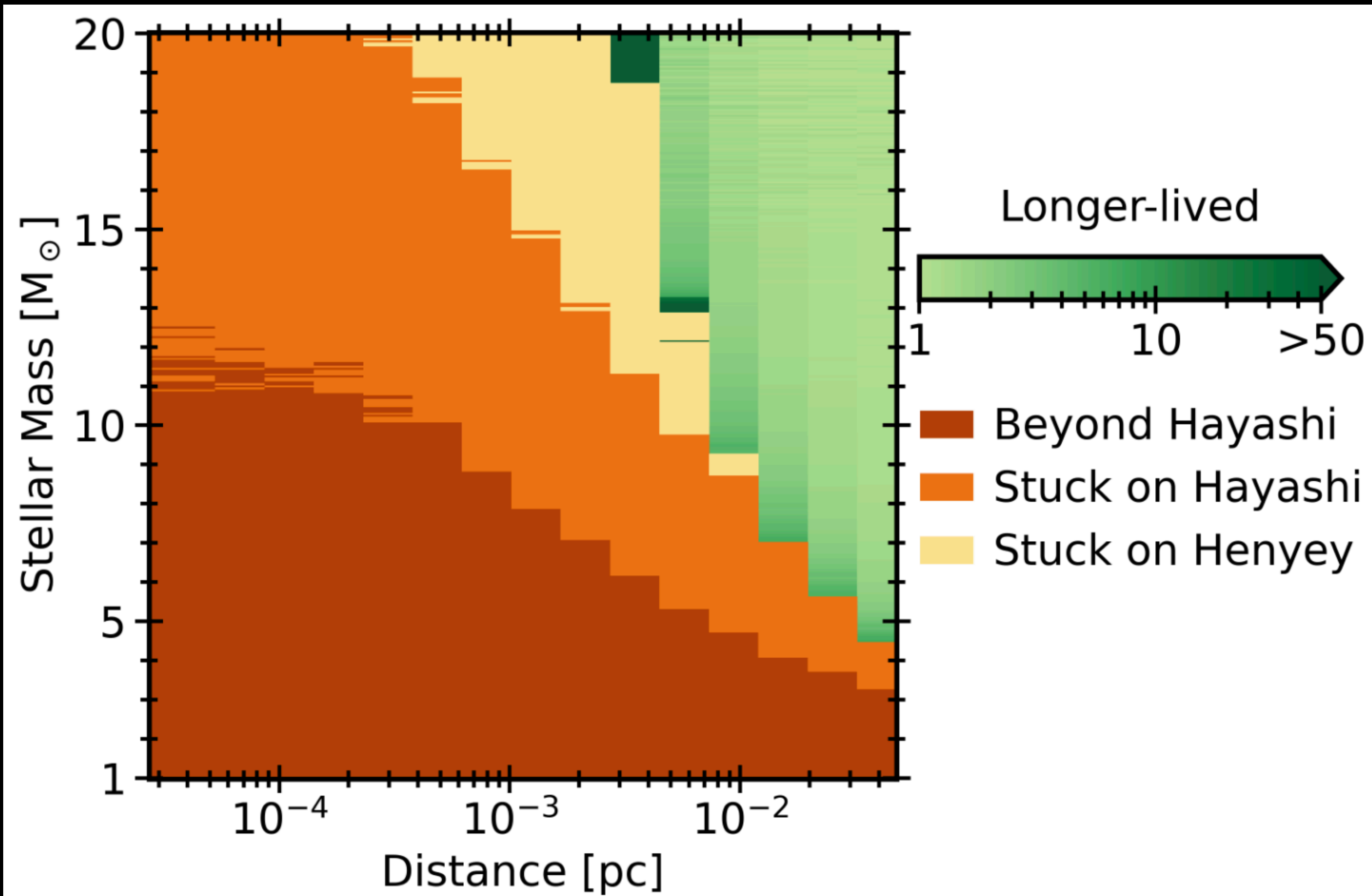
Immortal Stars



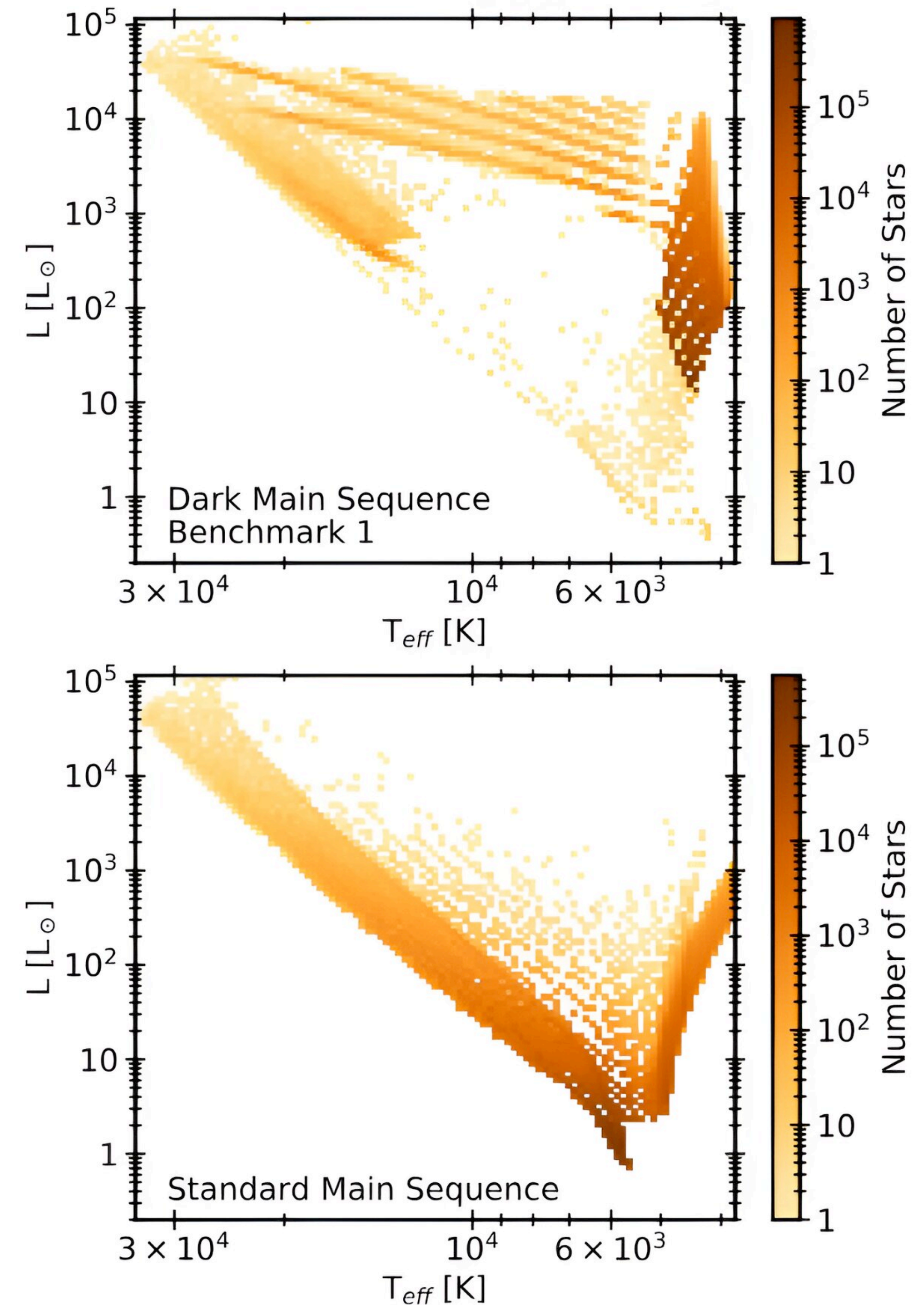




Immortal Stars



The type of signature observers can actually search for.



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Super-Kamiokande

Super-Kamiokande Strongly Constrains Leptophilic Dark Matter Capture in the Sun

Thong T.Q. Nguyen,^{1,*} Tim Linden,^{1,†} Pierluca Carenza,^{1,‡} and Axel Widmark^{1,2,§}

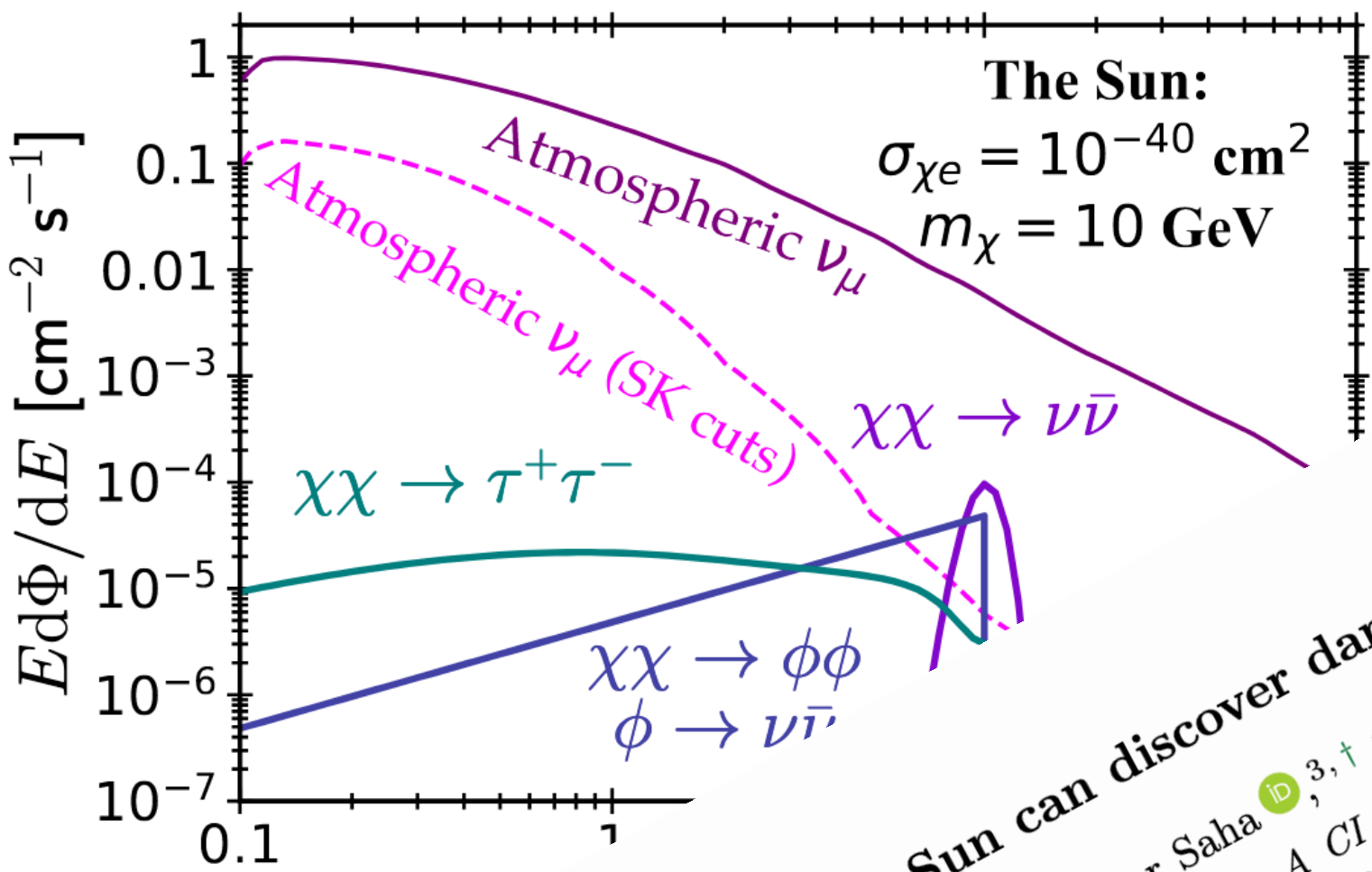
¹*Stockholm University and The Oskar Klein Centre for Cosmoparticle Physics, Alba Nova, 10691 Stockholm, Sweden*

²*Columbia University, 116th and Broadway, New York, NY 10027 USA*

The Sun can efficiently capture leptophilic dark matter that scatters with free electrons. If this dark matter subsequently annihilates into leptonic states, it can produce a detectable neutrino flux. Using 10 years of Super-Kamiokande observations, we set constraints on the dark-matter/electron scattering cross-section that exceed terrestrial direct detection searches by more than an order of magnitude for dark matter masses below 100 GeV, and reach cross-sections as low as $\sim 4 \times 10^{-41} \text{ cm}^{-2}$.

Introduction. — Detecting the particle interactions of dark matter is a cornerstone in our efforts to study beyond the standard model physics [1–3]. Many of the most sensitive constraints depend on searching for rare scattering interactions between the dark matter particle and standard model particles [4–11].

Current strategies motivate current searches. The standard model, used in terrestrial detectors, uses unimodal strategies to avoid astrophysical backgrounds by focusing on single dark matter interactions. However, to “go big”, use multimodal strategies to constrain rare scattering of dark



Neutrinos from the Sun can discover dark matter-electron scattering

Tarak Nath Maity^{1,2,3,*} Akash Kumar Saha^{3,†} Sagnik Mondal^{3,‡} and Ranjan Laha^{3,§}
¹Harish-Chandra Research Institute, A CI of Homi Bhabha National Institute, Chhatnag Road, Jhansi, Prayagraj (Allahabad) 211019, India
²Regional Centre for Accelerator-based Particle Physics, Harish-Chandra Research Institute, Prayagraj (Allahabad) 211019, India
³Indian Institute of Science, C.V. Raman Avenue, Bengaluru 560012, India
(Dated: August 25, 2023)

Dark matter-electron scattering using high-energy neutrino observations from the Super-Kamiokande (Super-K) with electrons can get captured inside the Sun. These captured dark matter particles (SM) particles. Neutrinos produced from the Standard Model (SM) particles. Although there is no excess of neutrinos from the Super-Kamiokande (Super-K) data-sets of IceCube and DeepCore set the limits on the DM mass range 10 GeV to 10⁵ GeV. These telescopes have the potential

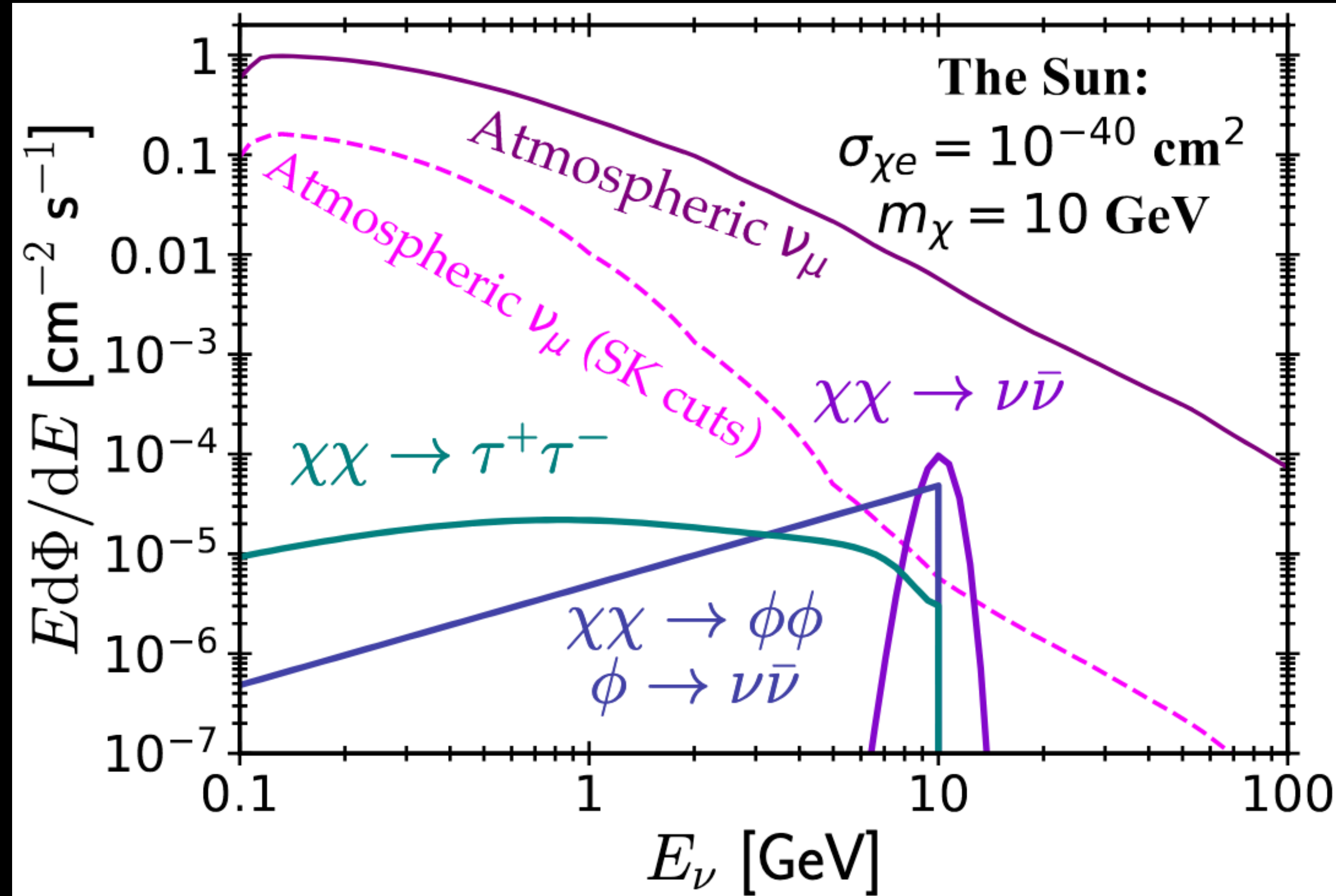
DAMA/LIBRA and leptonically interacting Dark Matter

Thim Kopp,^{1,*} Viviana Niro,^{1,†} Thomas Schwetz,^{1,‡} and Jure Zupan^{1,§}
¹Max-Planck-Institute for Nuclear Physics, P.O. Box 103980, 69029 Heidelberg, Germany
²Theory Division, Physics Department, CERN, CH-1211 Geneva 23, Switzerland

CERN-PH-TH/2009-116

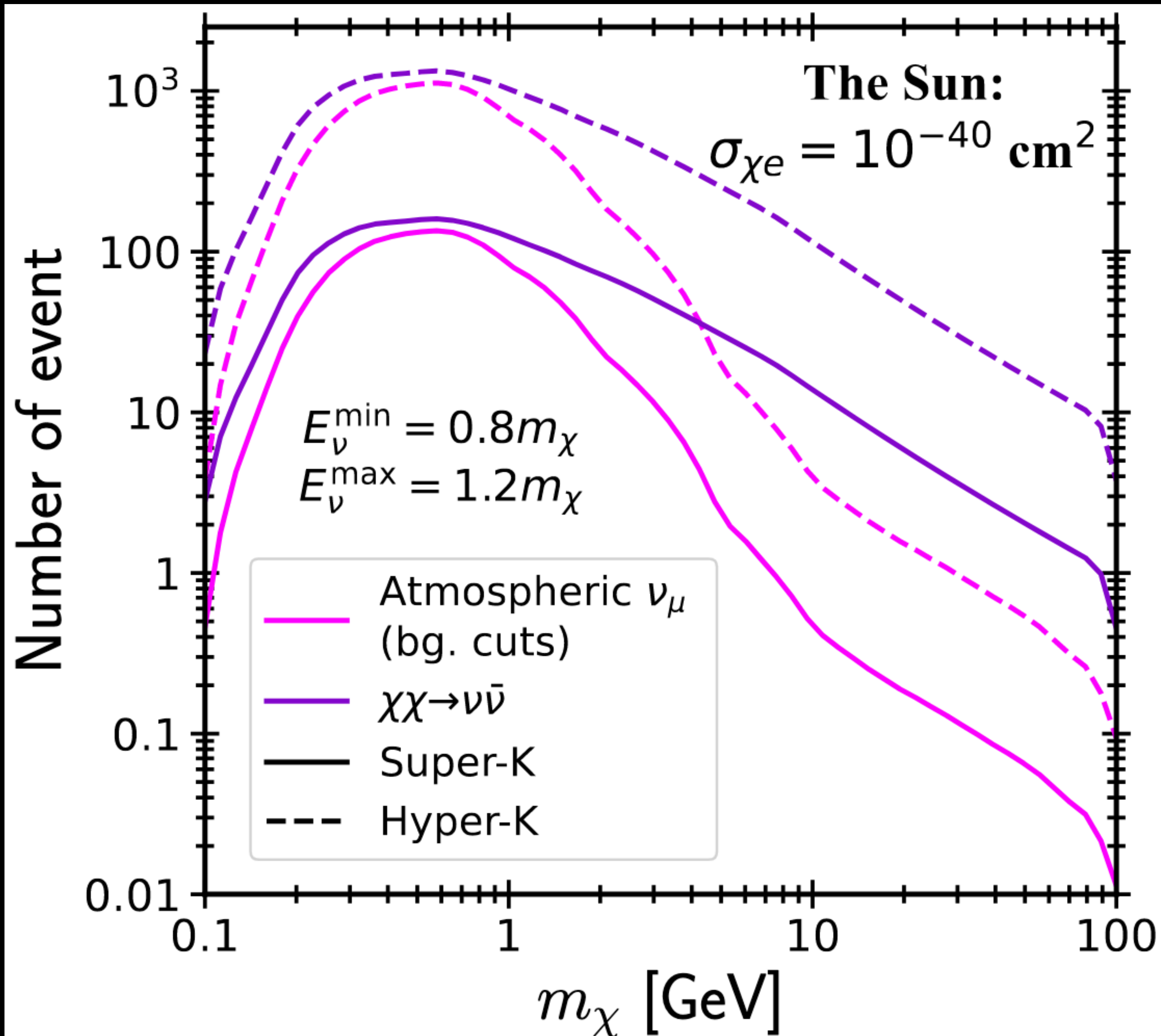
HRI-RECAPP-2023-03

Super-Kamiokande



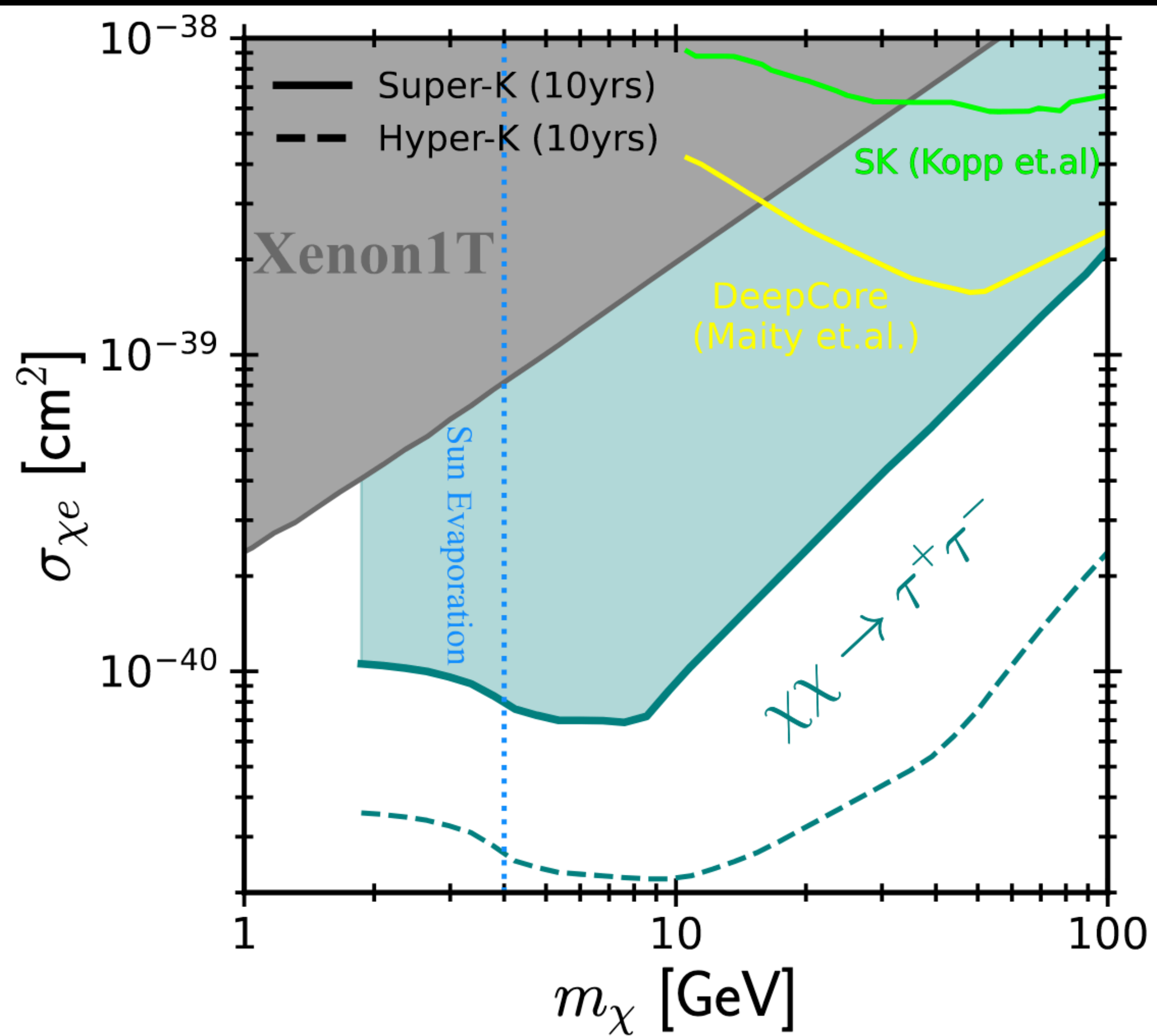
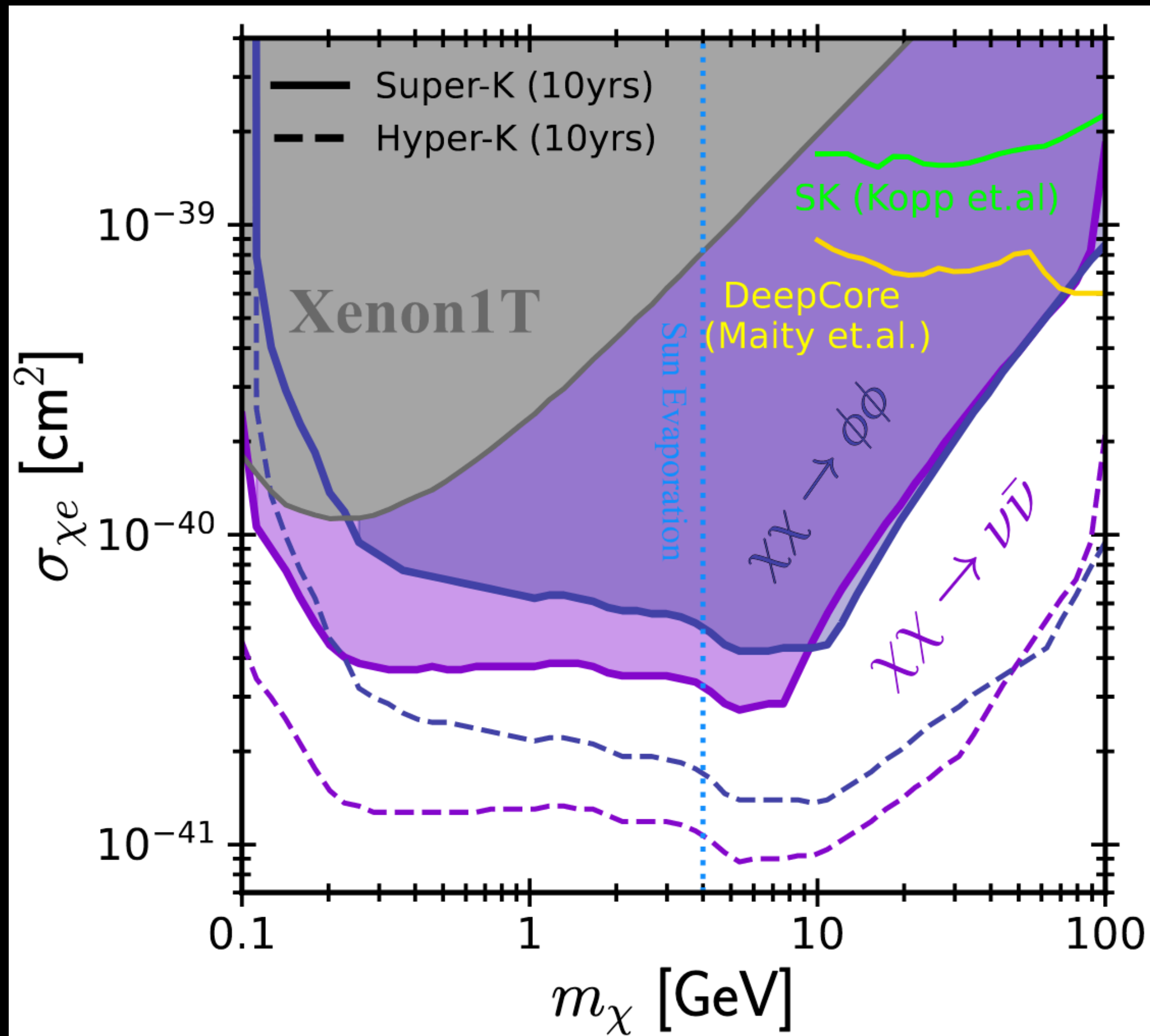
- **“Miracle” - Dark Matter must be leptophilic.**
- **Standard annihilation cross-sections and unconstrained scattering rate.**
- **Significant annihilation rate to neutrinos (or taus).**

Super-Kamiokande



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Detecting Dark Matter with Imploding Pulsars in the Galactic Center

Joseph Bramante

*Department of Physics, 225 Nieuwland Science Hall,
University of Notre Dame, Notre Dame, IN 46556, USA*

Tim Linden

*Kavli Institute for Cosmological Physics 5640 South Ellis Avenue
University of Chicago Chicago, IL 60637*

of old millisecond pulsars observed at the galactic center of the Milky Way. We consider the possibility of matter accumulating in and destroying neutron stars. In regions where the density of matter clumped in a pulsar can exceed the Schwarzschild limit, a black hole is formed which destroys the pulsar. We examine what parameters of a pulsar determine whether it survives and find regions of parameter space where a significant fraction of the neutron star population within the galactic center survives. Millisecond pulsars in globular clusters and the galactic center might cause the formation of black holes and supermassive black holes.

*Kavli Institute for Cosmology
University*

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THE PECULIAR PULSAR POPULATION OF THE CENTRAL PARSEC

JASON DEXTER
Department of Astronomy, University of California, Berkeley, CA 94720-3411

RYAN M. O'LEARY
Department of Astronomy, University of California, Berkeley, CA 94720-3411

Draft version April 14, 2018

ABSTRACT

the galactic center black hole, Sgr A*,
used to test general relativity
the explanation in
t dis

PULSAR POPULATION OF THE CENTRAL PA

JASON DEXTER
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Draft version April 14, 2018

ABSTRACT

Sgr A*, would be potentially used to test general relativity. Despite recent discovery of radio pulsars in the Galactic center, no explanation has been that the temporal variation in the ordinary pulsar population shows that the extraordinary pulsar population implies a new type of pulsar.

ABSTRACT

le, Sgr A*, would be potential probes of its mass, distance
Galactic center, none have been detected within 25
been that hyperstrong temporal scattering prevents
pulsations from a highly magnetized neutron
scattering is much weaker than predicted
populating the most likely reason for
the Galactic center. In contrast
etar formation in the region
rdinary pulsars, their
could be caused
al mass f
million

The Galactic centre pulsar population

Galactic centre pulsar population

Jayanth Chennamangalam^{1*} and D. R. Lorimer^{1,2†}

¹ Department of Physics and Astronomy, West Virginia University, PO Box 6315, Morgantown, WV 26506, USA

² National Radio Astronomy Observatory, PO Box 2, Green Bank, WV 24944, USA

ABSTRACT

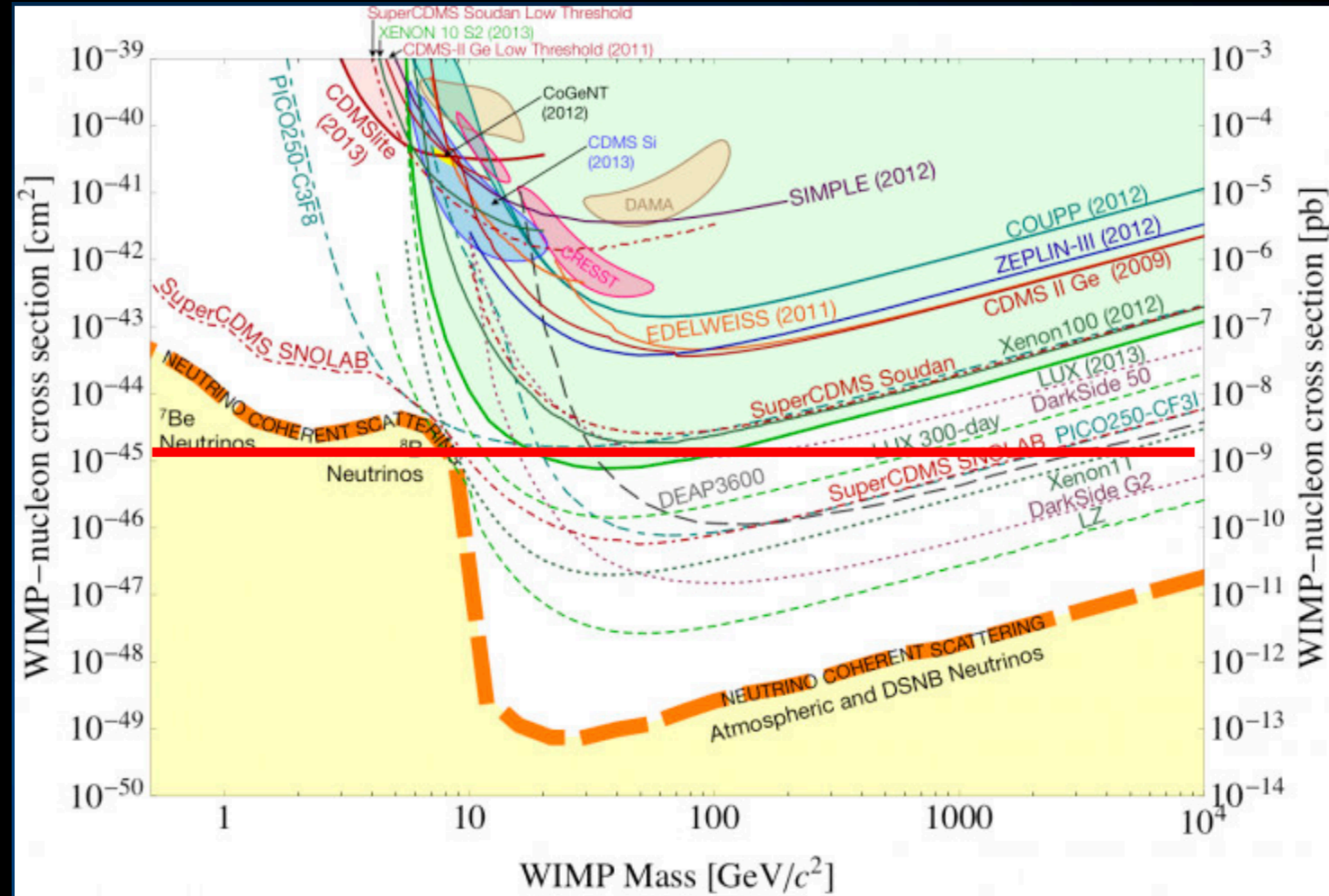
The recent discovery of a millisecond pulsar by Spitler et al. to characterize that the temporal broadening in that predicted by the number of what the pulsar is emitting re-

ABSTRACT

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The recent discovery of a magnetar in the Galactic centre region (Spitler et al. 2015) has led to a renewed interest in the possibility that the temporal broadening of the pulse profile of the magnetar is due to the number of pulsars beaming towards the Earth – in the inner part of the Galactic centre region is the same as that in the rest of the Galaxy. We consider the question of what the plausible limits for the electron density of the interstellar medium, using reasonable assumptions – namely, (i) the pulsar formation rate, (ii) the spin and luminosity distribution of pulsars, (iii) the scattering in the interstellar medium, and (iv) the scattering in the inner part of the Galaxy, are. We conclude that the number of pulsars in the inner part of the Galaxy is similar to the number of pulsars in the rest of the Galaxy, and that the pulsar formation rate is similar to the pulsar formation rate in the rest of the Galaxy. We conclude that the pulsar formation rate is similar to the pulsar formation rate in the rest of the Galaxy.

Neutron Stars as the Optimal Dark Matter Detectors

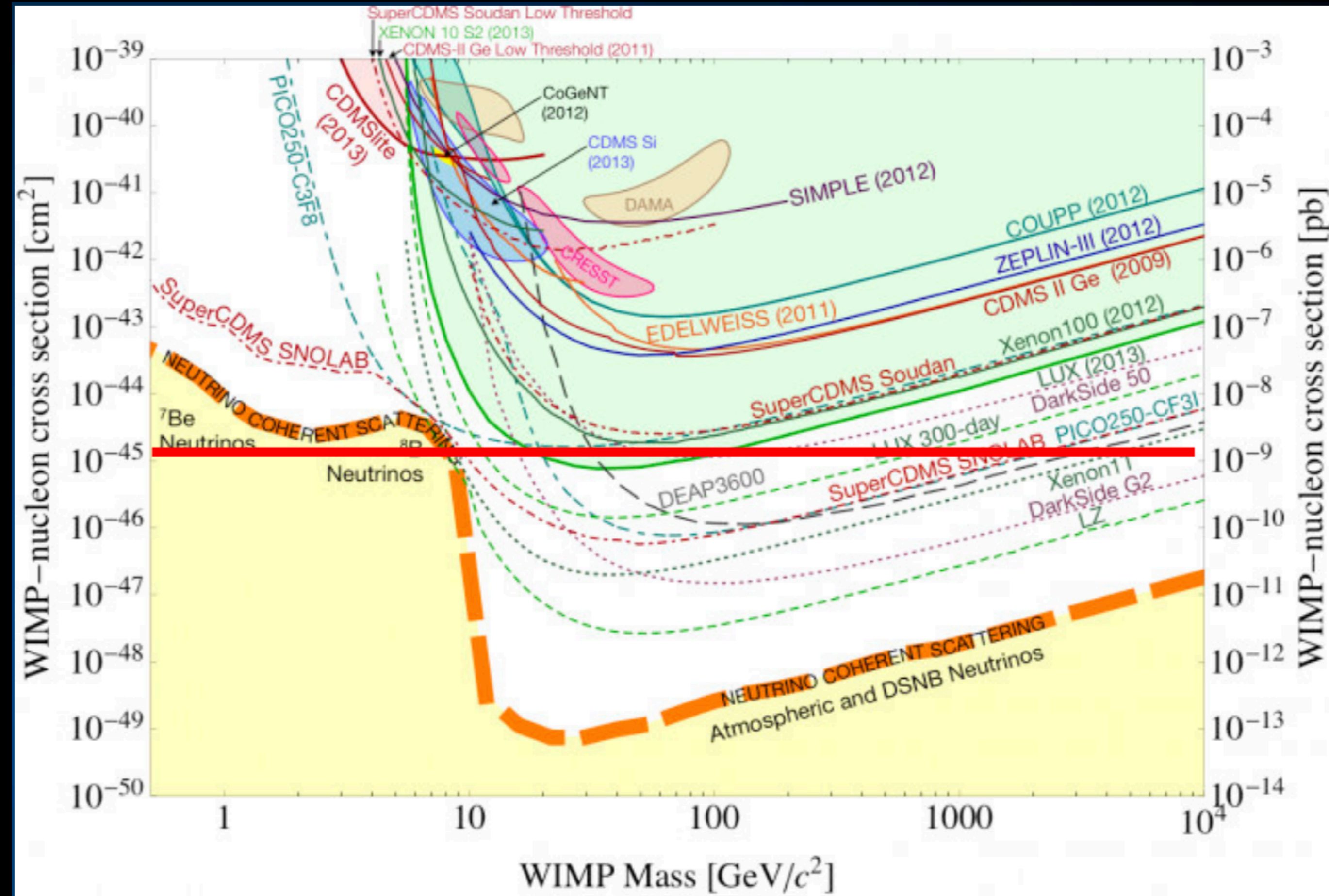
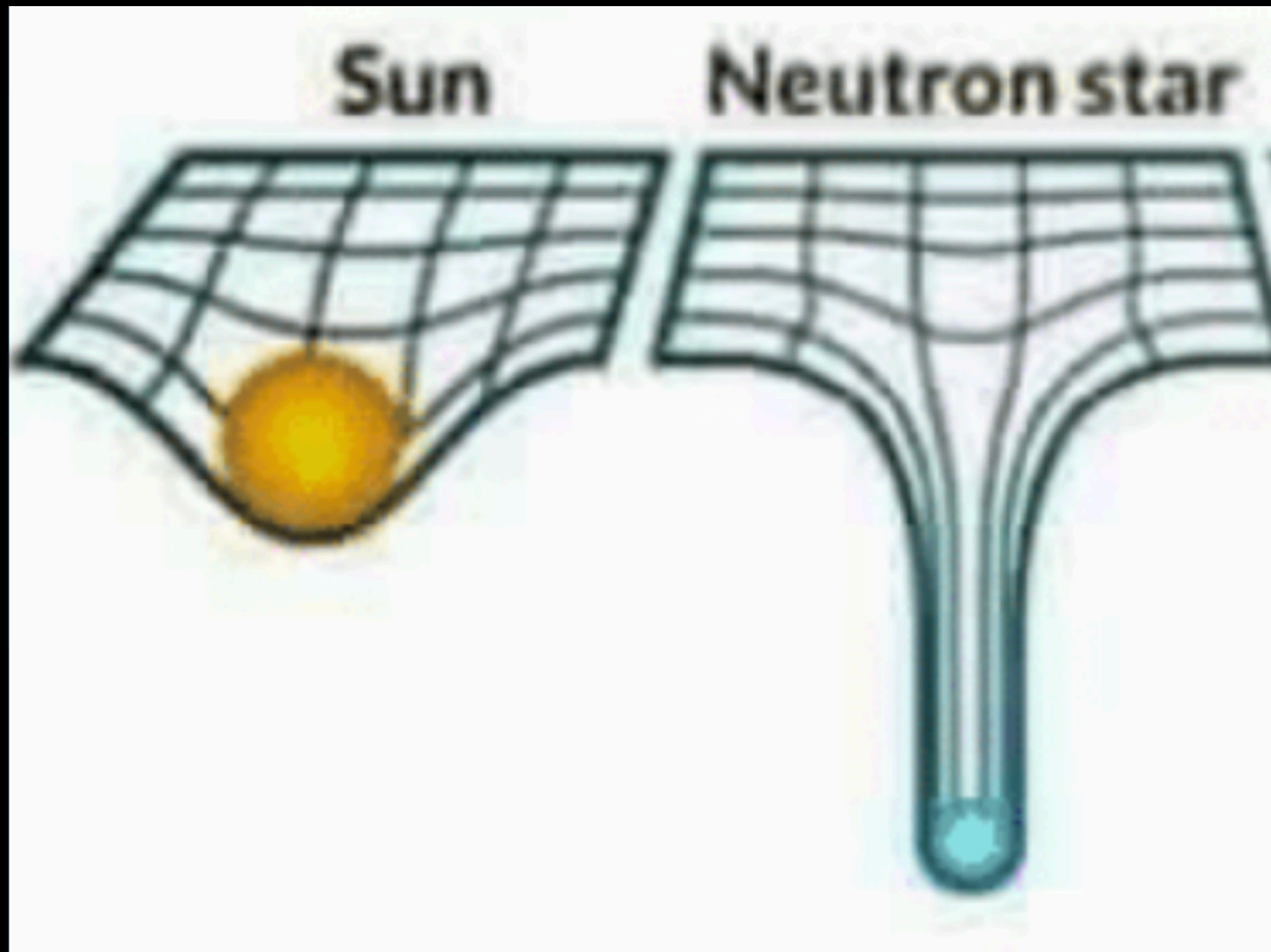
- Direct detection cross-sections near 10^{-46} cm^2 produce $\sim 1 \text{ event}/(\text{ton yr})$
- Sensitivity peaks near Xenon mass, and falls off significantly at higher or lower masses.



$$\sigma_{\text{sat}}^{\text{single}} \simeq \pi R^2 m_n / M \simeq 2 \times 10^{-45} \text{ cm}^2 \left(\frac{1.5 M_\odot}{M} \right) \left(\frac{R}{10 \text{ km}} \right)^2$$

Neutron Stars as the Optimal Dark Matter Detectors

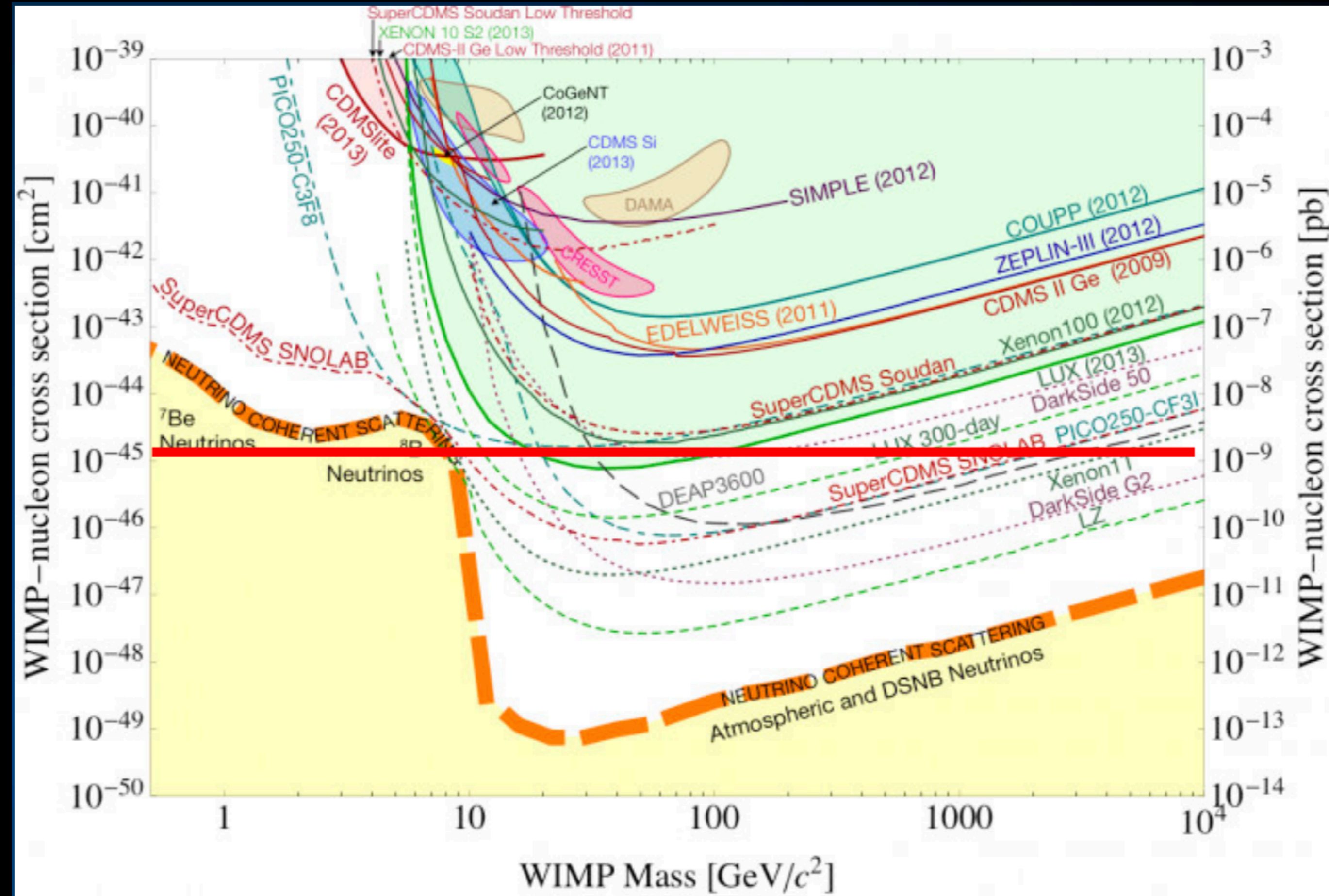
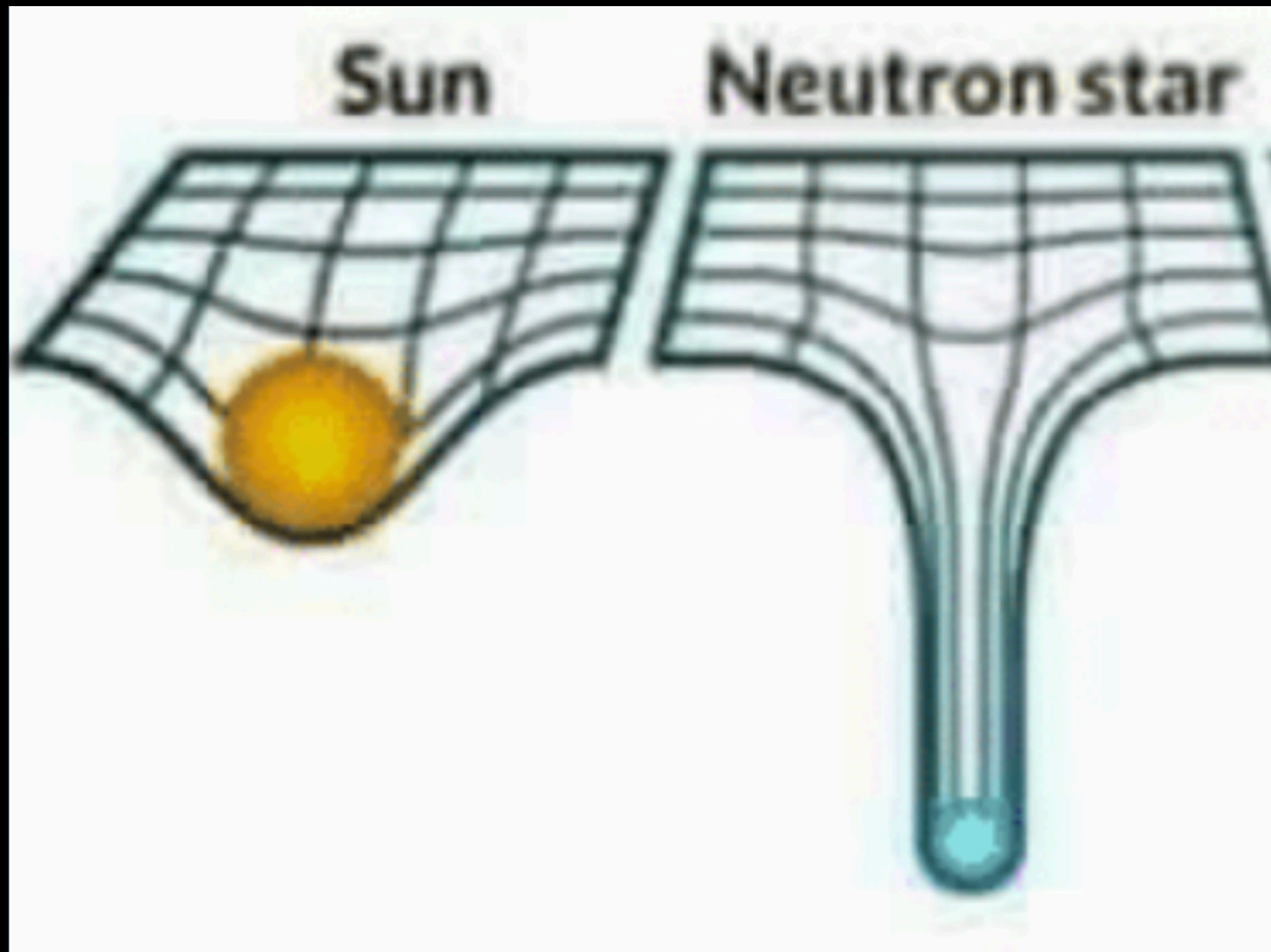
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$$b_{\text{max}} = \left(\frac{2GM R}{v_x^2} \right)^{1/2} \left(1 - \frac{2GM}{R} \right)^{-1/2}$$

Neutron Stars as the Optimal Dark Matter Detectors

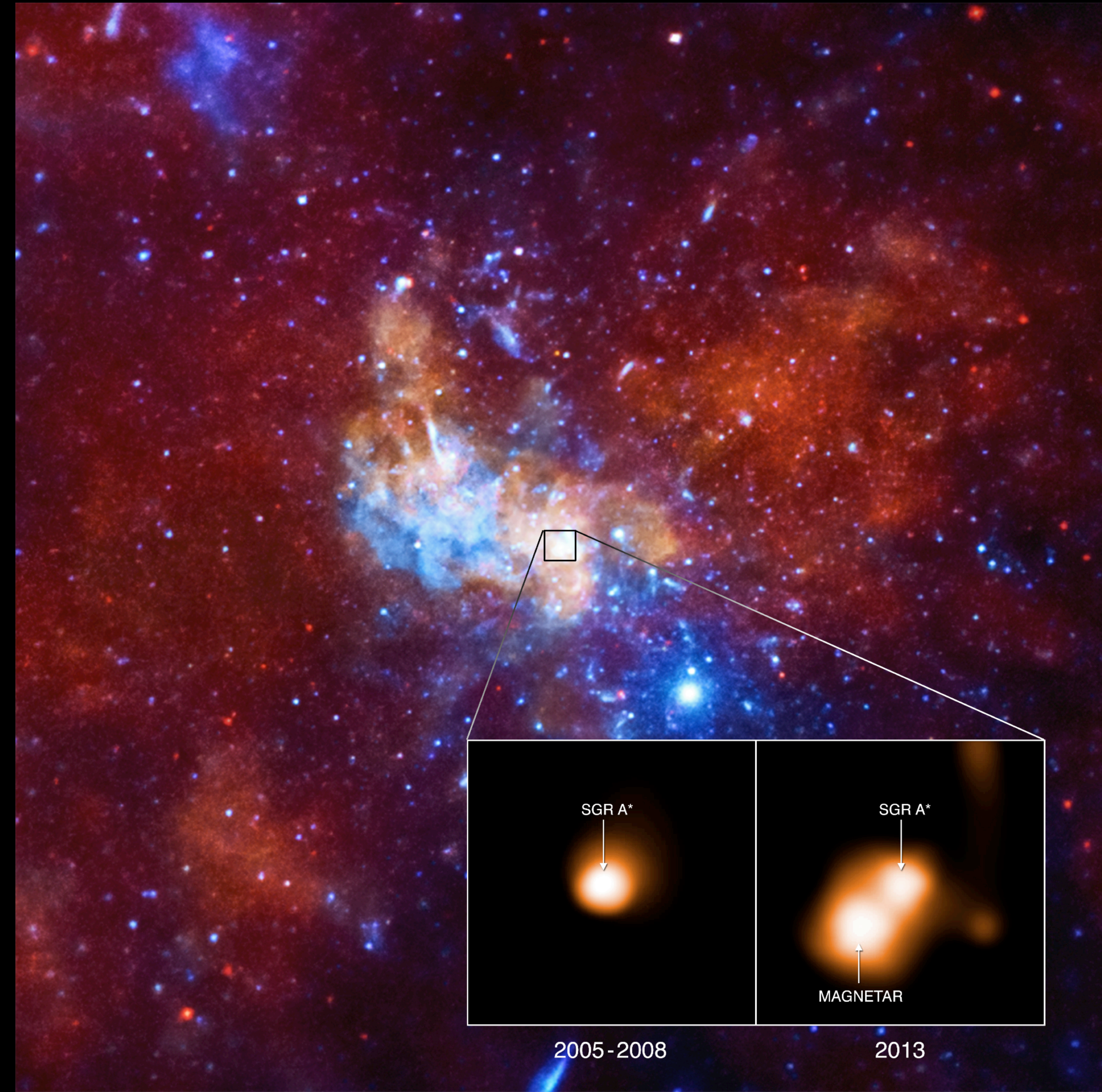
- Direct detection cross-sections near 10^{-46} cm^2 produce $\sim 1 \text{ event}/(\text{ton yr})$



$$v_{esc} = \sqrt{\frac{2GM}{r}} \sim 0.7c$$

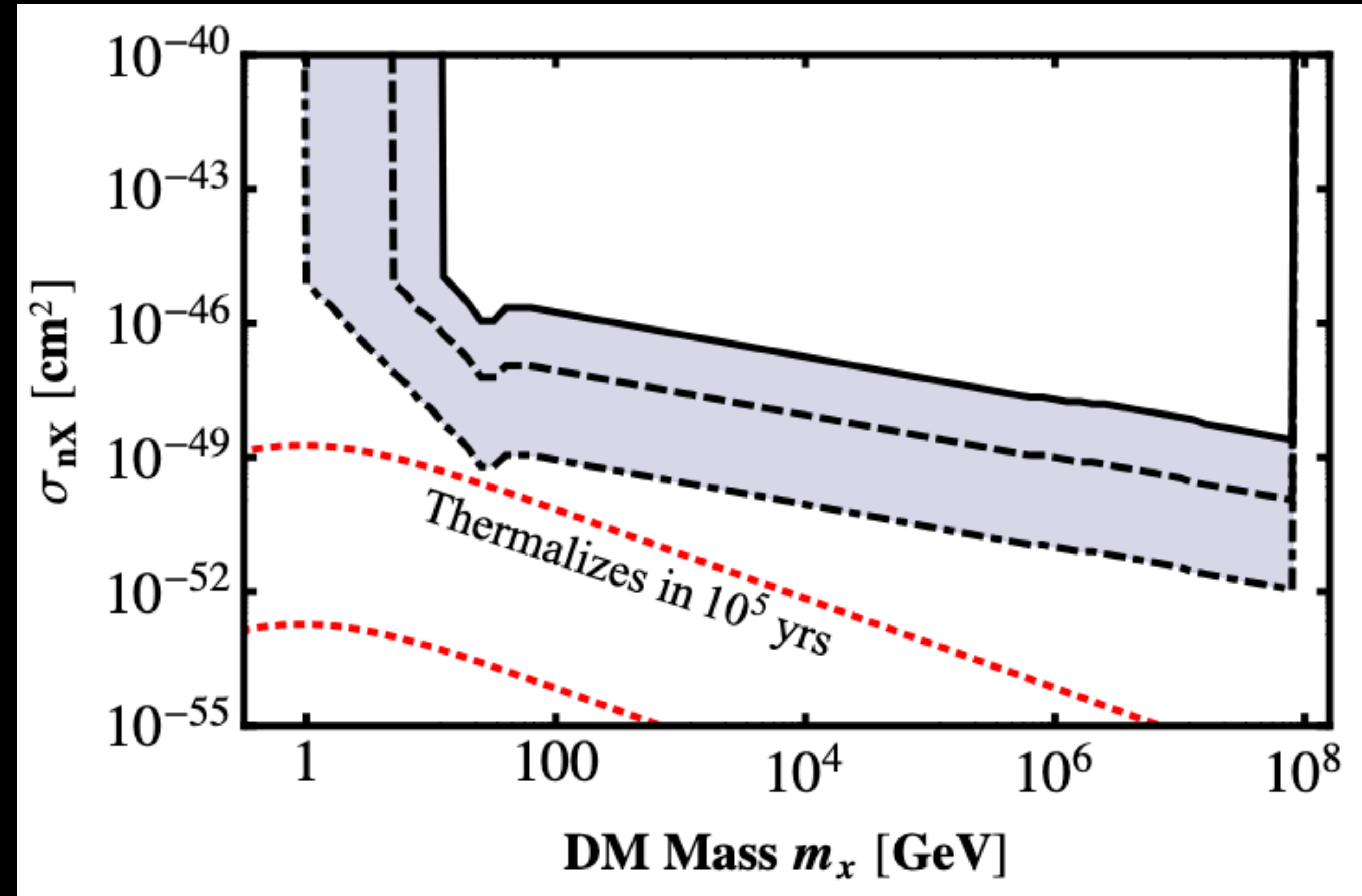
The 2013 GC Magnetar Detection

- Before 2013, it was thought that pulse dispersion made GC pulsars invisible.
- Magnetar observed (**First in X-rays!**) in 2013.
- Largest pulse dispersion of any existing pulsar, but not sufficient to make GC pulsars invisible.



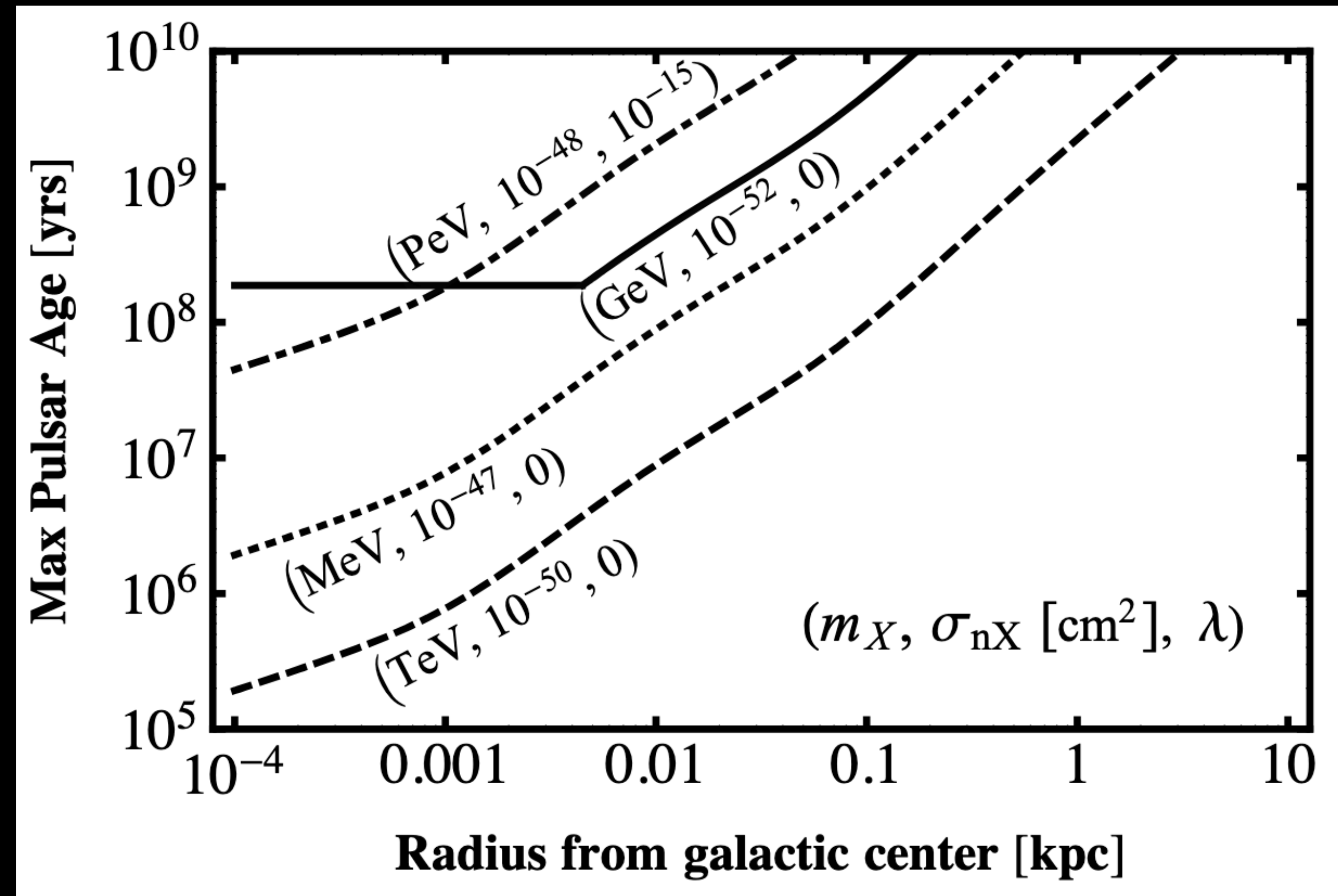
Dark Matter interactions in NS

- Dark Matter accumulation can eliminate these pulsars:
- **Need DM to either be non-annihilating massive fermions or bosons (to avoid Fermi degeneracy pressure)**



Dark Matter interactions in NS

- Dark Matter accumulation can eliminate these pulsars:
- **Need DM to either be non-annihilating massive fermions or bosons (to avoid Fermi degeneracy pressure)**
- **A New Astrophysical observable! A maximum age for pulsars that depends on GC radius.**



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- Standard model miracles cost half.
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White Dwarfs in Dwarf Spheroidal Galaxies: A New Class of Compact-Dark-Matter Detectors

Juri Smirnov,^{1,2,*} Ariel Goobar,^{2,†} Tim Linden,^{2,‡} and Edvard Mörtzell^{2,§}

¹*Department of Mathematical Sciences, University of Liverpool, Liverpool, L69 7ZL, United Kingdom*

²*The Oskar Klein Centre, Department of Physics, Stockholm University, AlbaNova, SE-10691 Stockholm, Sweden*

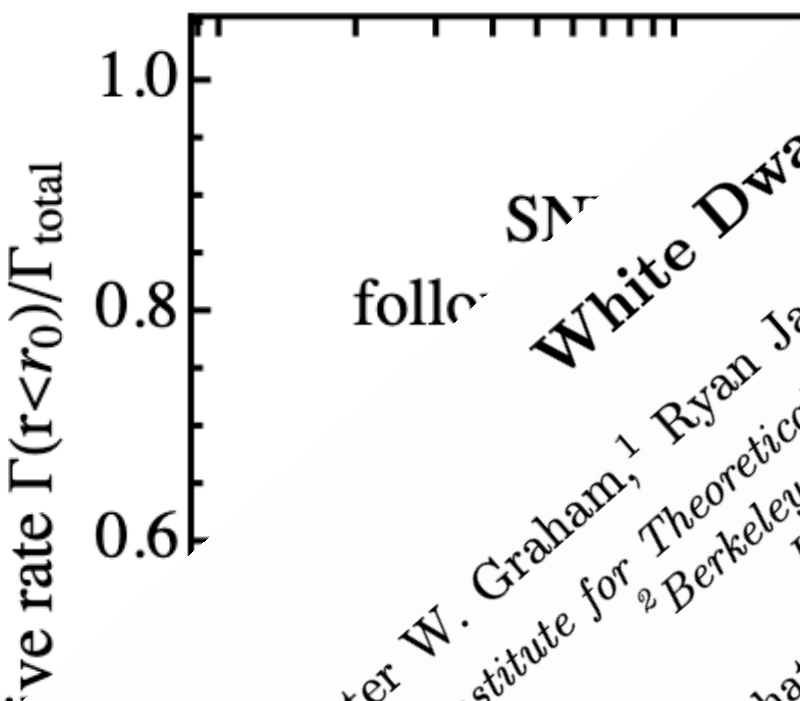
Recent surveys have discovered a population of faint supernovae, known as Ca-rich gap transients, inferred to originate from explosive ignitions of white dwarfs. In addition to their unique spectra and luminosities, these supernovae have an unusual spatial distribution and are predominantly found at large distances from their presumed host galaxies. We show that the locations of Ca-rich gap transients are well matched to the distribution of dwarf spheroidal galaxies surrounding large galaxies, in accordance with a scenario where dark matter interactions induce thermonuclear explosions among low-mass white dwarfs that may be otherwise difficult to explain by standard stellar or binary evolution mechanisms. A plausible candidate to explain the observed events is primordial black holes with masses above 10^{21} grams.

Dark Matter Triggers of Supernovae

Peter W. Graham,¹ Surjeet Rajendran,² and Jaime Varela²
¹*Stanford Institute for Theoretical Physics, Stanford, CA 94305*
²*Berkeley Center for Theoretical Physics, Department of Physics, University of California, Berkeley, CA 94720*

Abstract

We find a new population of super-
novae [1, 2], called Ca-Rich Gap
transients, which trace
the distribution of dark matter
at a much larger
scale than Type Ia SN, which trace
the distribution of stars. Additionally,
they may originate from
explosions near the Chandrasekhar limit,
or from the tidal disruption of
primordial black holes, or from
other mechanisms.

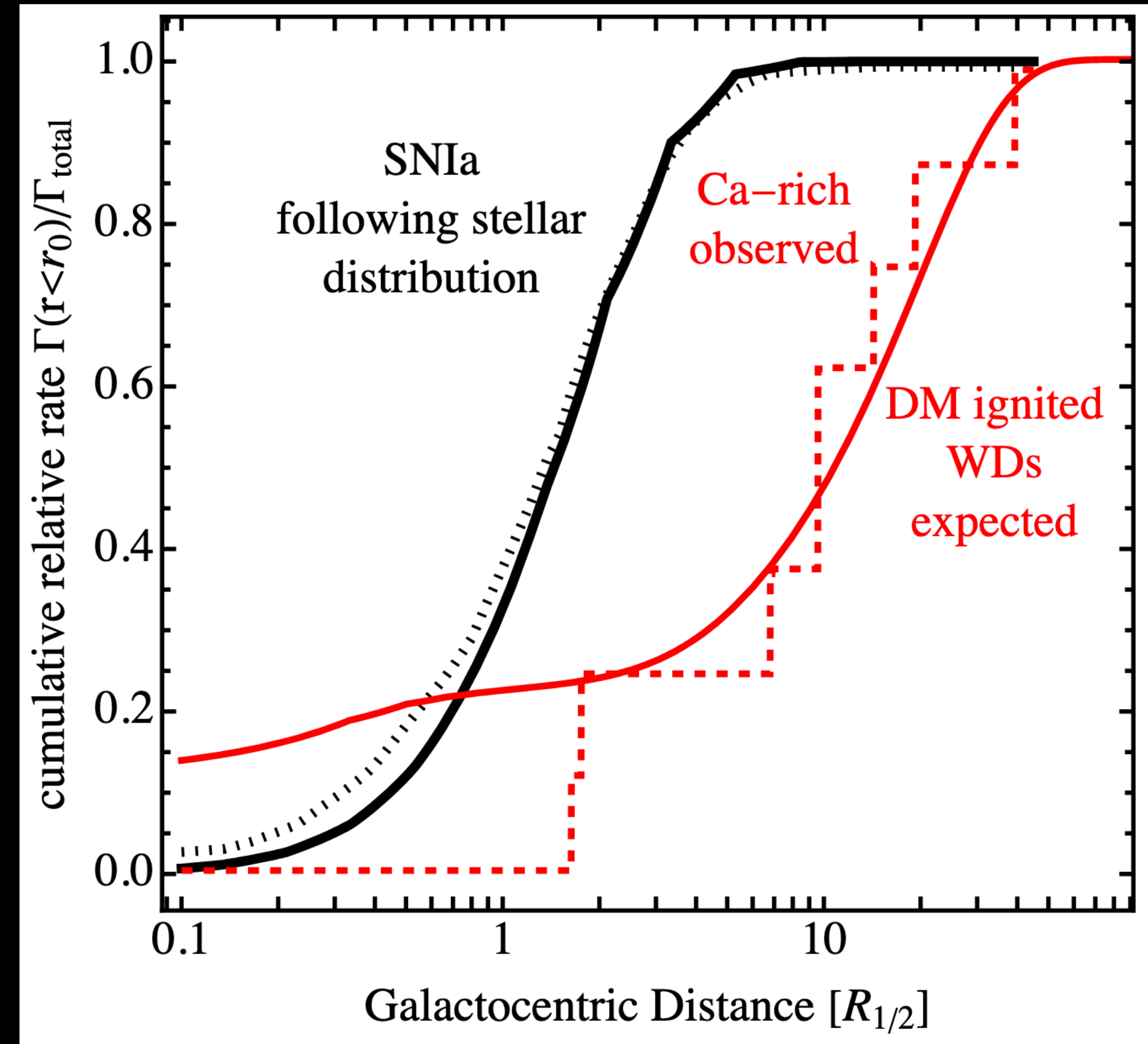


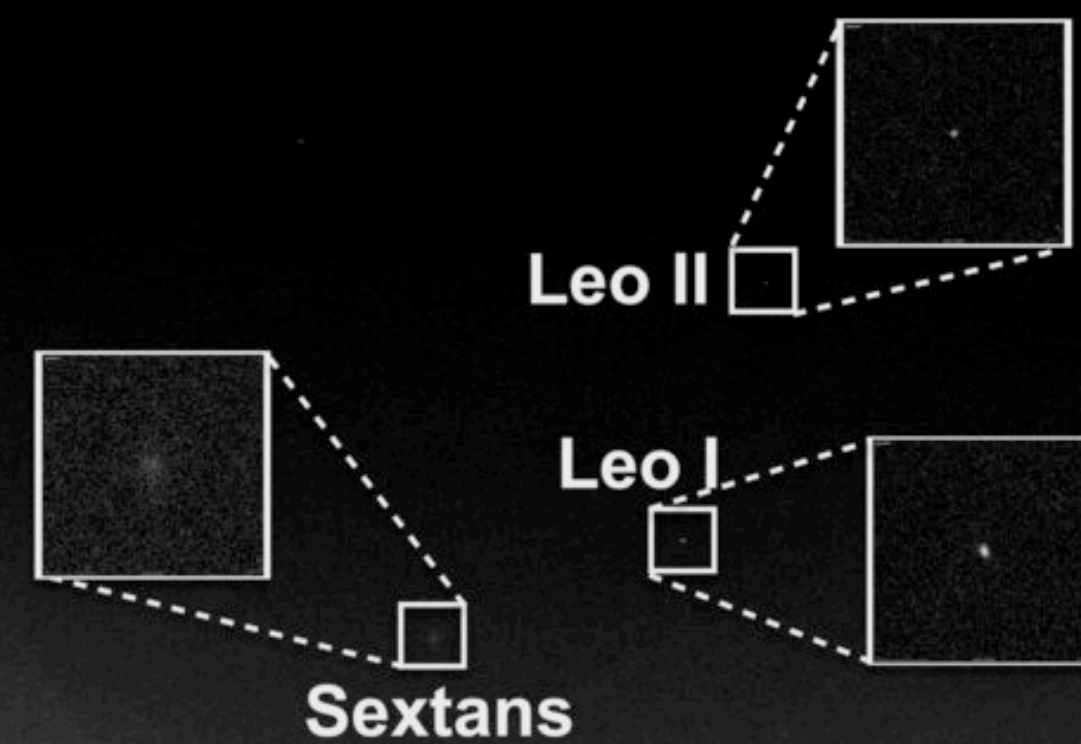
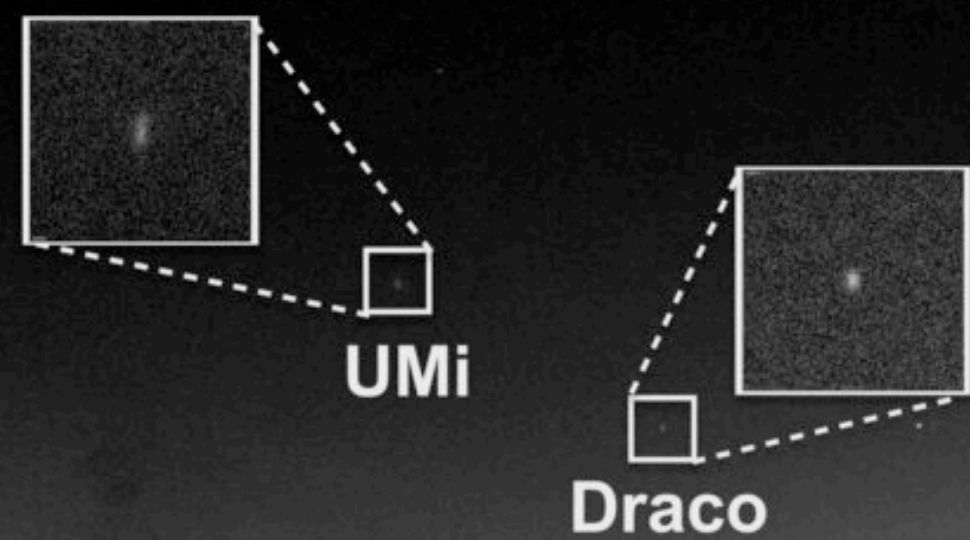
White Dwarfs as Dark Matter Detectors

Peter W. Graham,¹ Ryan Janish,² Vijay Narayan,² Surjeet Rajendran,² and Paul Riggins²
¹*Institute for Theoretical Physics, Department of Physics, Stanford University, Stanford, CA 94305*
²*Berkeley Center for Theoretical Physics, Department of Physics, University of California, Berkeley, CA 94720, USA*
We show that a population of white dwarfs can act as a local region in a white dwarf will trigger
the ignition of a type Ia supernova. This was originally proposed by Graham et al. and
is capable of sufficiently heating a local region in a white dwarf via dynamical friction.
Primordial black holes which transit and heat a white dwarf through the production of high-
energy dark matter (DM) candidates that heat through the production of high-
energy particles, and show that such particles will efficiently thermalize the
supernovae. Based on the existence of long-lived white dwarfs and
derive new constraints on ultra-heavy DM with masses greater
than those of the standard model. As a concrete example, we place bounds
on the formation and self-gravitating nature of white dwarfs, considering the formation and self-gravitating
decays and annihilations within the core of a white dwarf. This work provides an alternative
scenario for the formation of white dwarfs.

Fitting both the Number and Distribution of Ca-Rich Transients

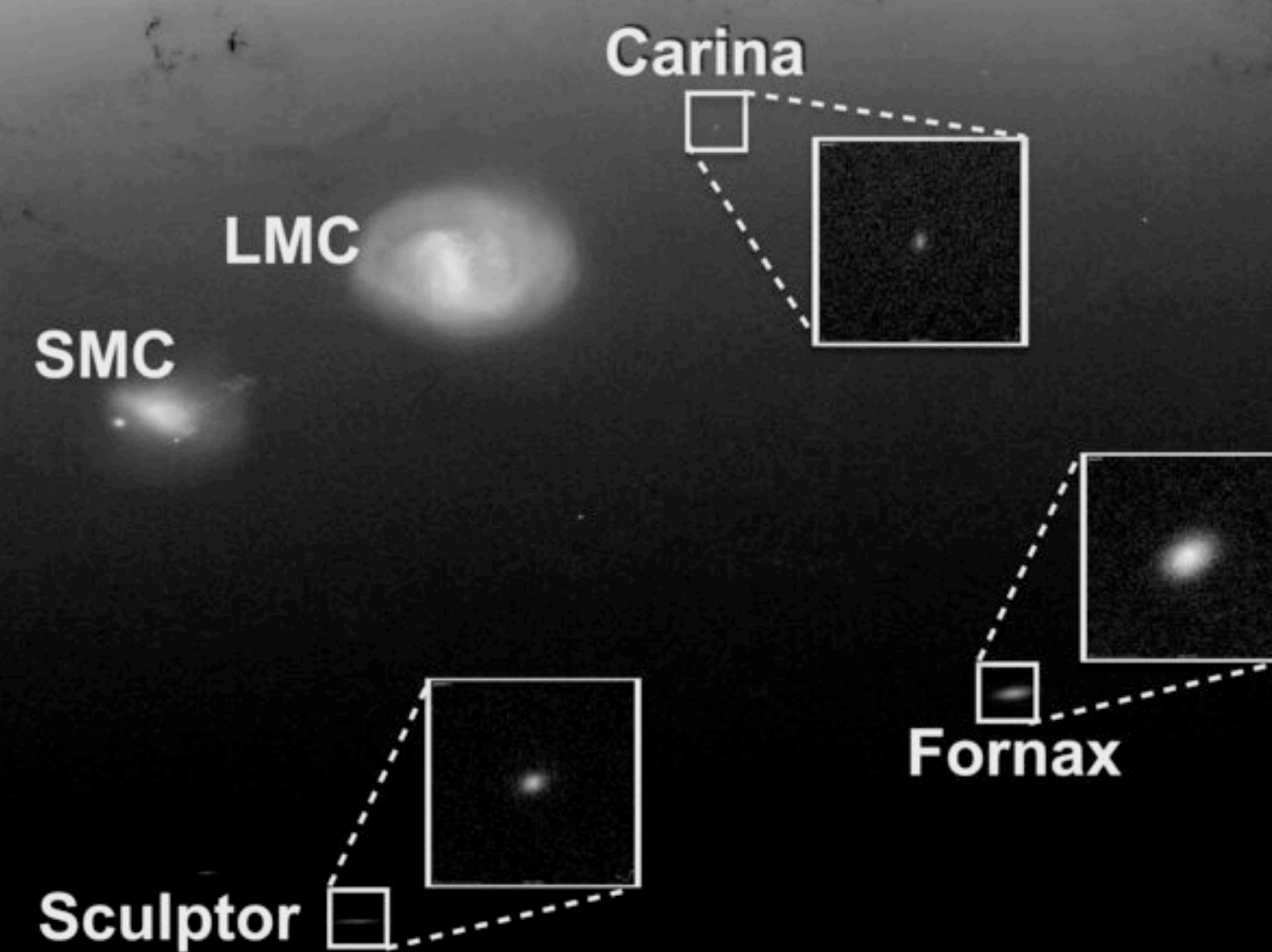
- **Low-luminosity and high-calcium content of Ca-rich population indicates low-mass progenitors ($\sim 0.6 M_{\odot}$, far below the Chandrasekhar mass)**
- **Distribution of events in galactic center radius is also unusual.**

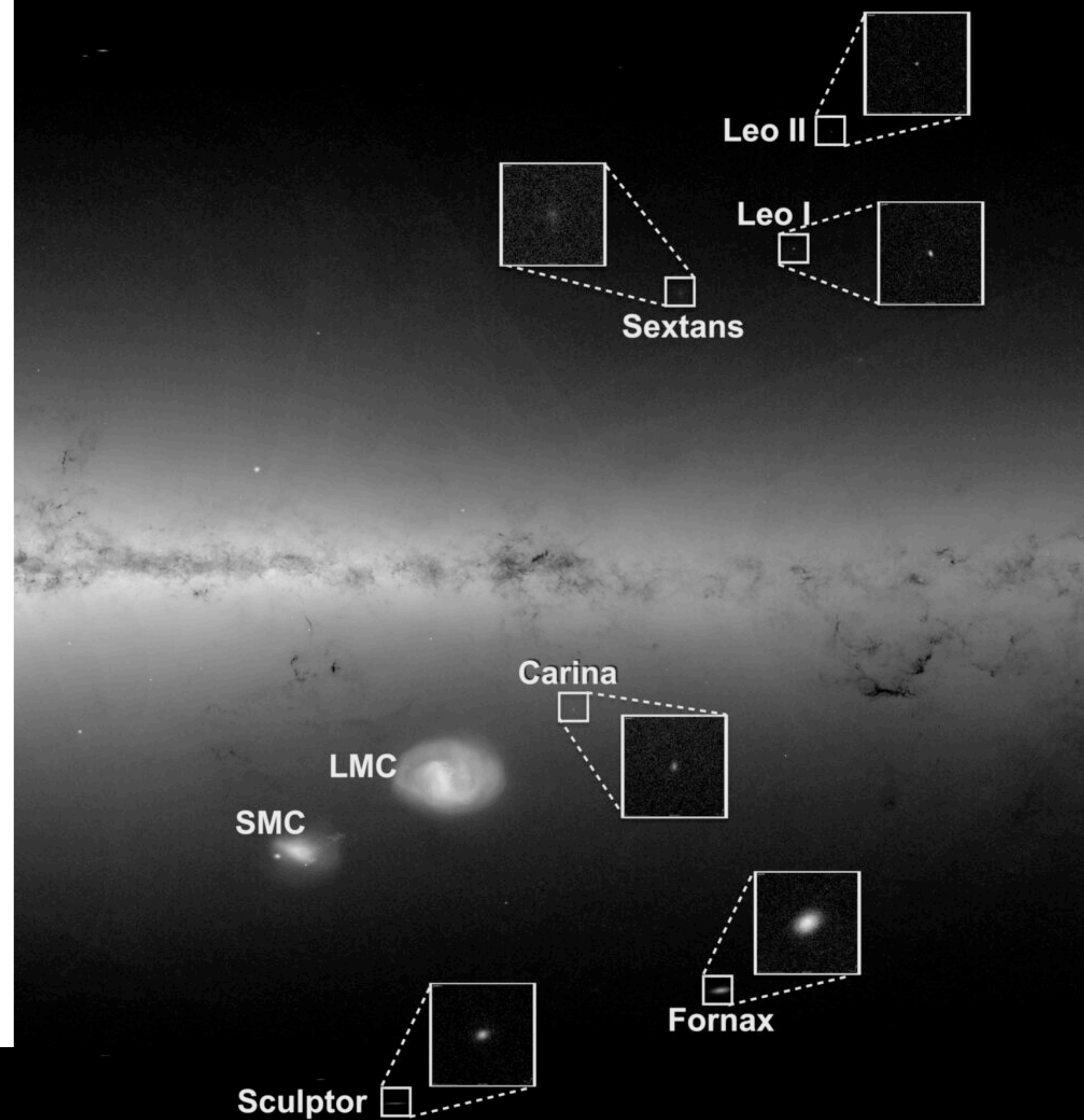
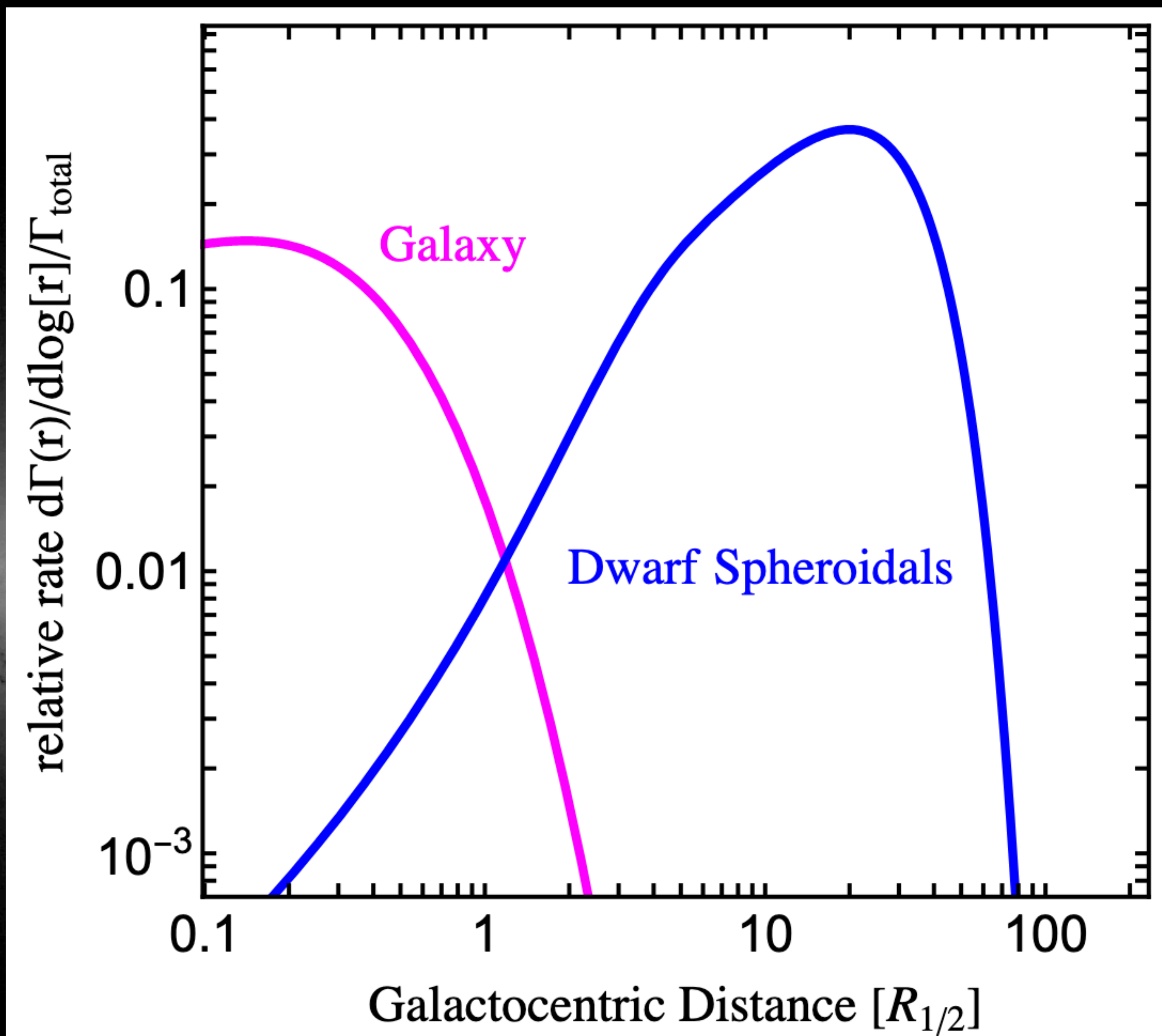




$$C_X \simeq \sqrt{\frac{6}{\pi}} \left(\frac{\rho_X}{\bar{v}_X} \right) \frac{\xi N_B v_{\text{esc}}^2}{m_X} \left[1 - \frac{1 - \exp(-B^2)}{B^2} \right] f(\sigma_{nX}) \quad (1)$$

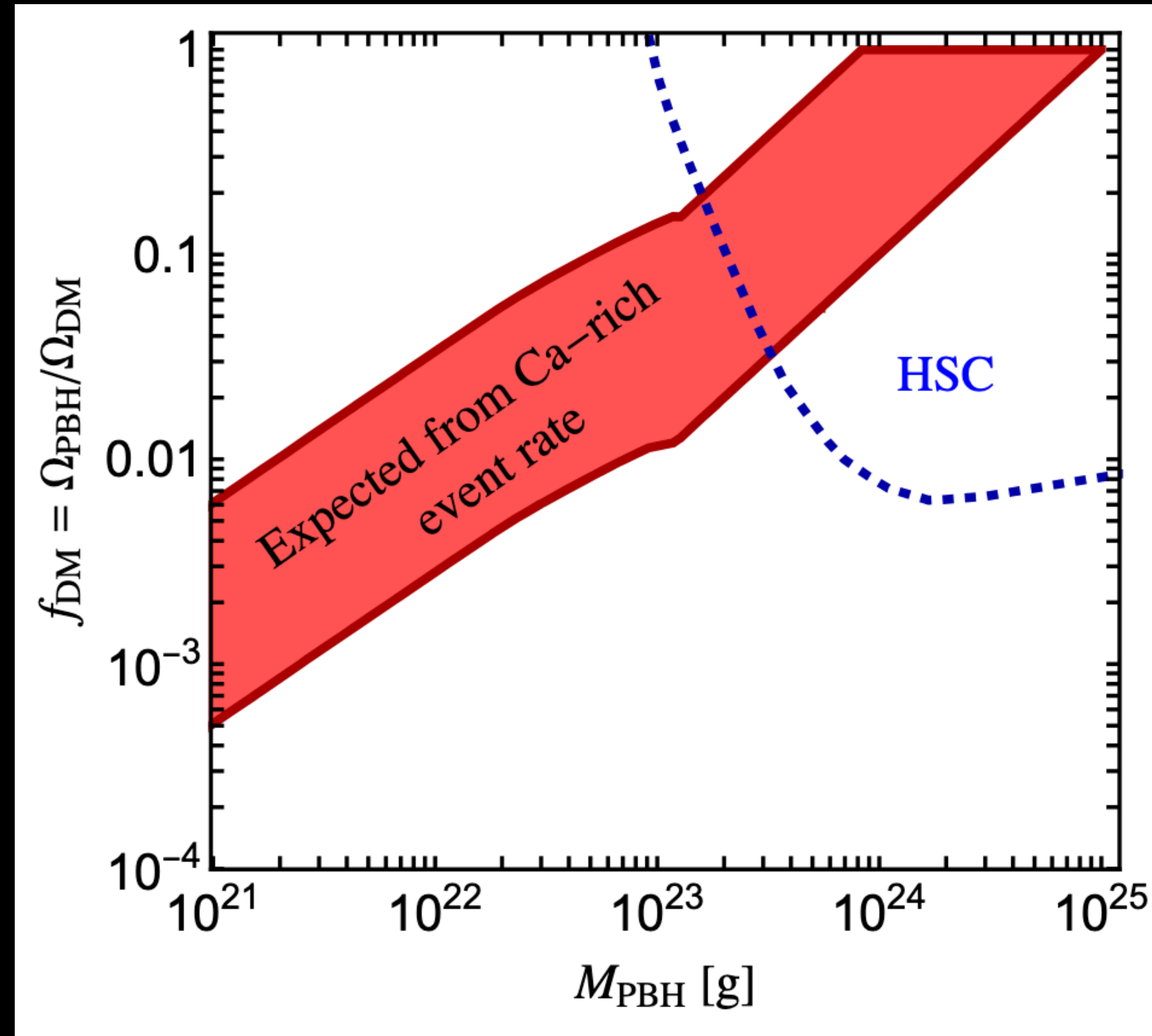
$$v_{\text{esc}} = 10.9 \left(\frac{M}{10^7 M_{\odot}} \right)^{1/3} \left[\frac{1+z}{9.5} \right]^{1/2} \frac{\text{km}}{\text{s}}$$





Fitting both the Number and Distribution of Ca-Rich Transients

- “Miracle” - Dark Matter must be relatively low-mass black holes (but can be a subdominant portion of the total dark matter density).
- Standard kinematic interaction rates and dark matter abundance.



Observational Follow-Ups are Motivated

- Searches for Ca-Rich SNe are a key science component for upcoming LSST analyses.
- JWST follow-ups of these sources can potentially detect nearby dwarf galaxies.

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How to Do Science in the High-Risk High-Reward Regime

- **“What we need to do is combine to form a larger collaboration that can push the type of dedicated experiments that are capable of setting strong limits on dark matter in celestial bodies.”**










































How to Do Science in the High-Risk High-Reward Regime

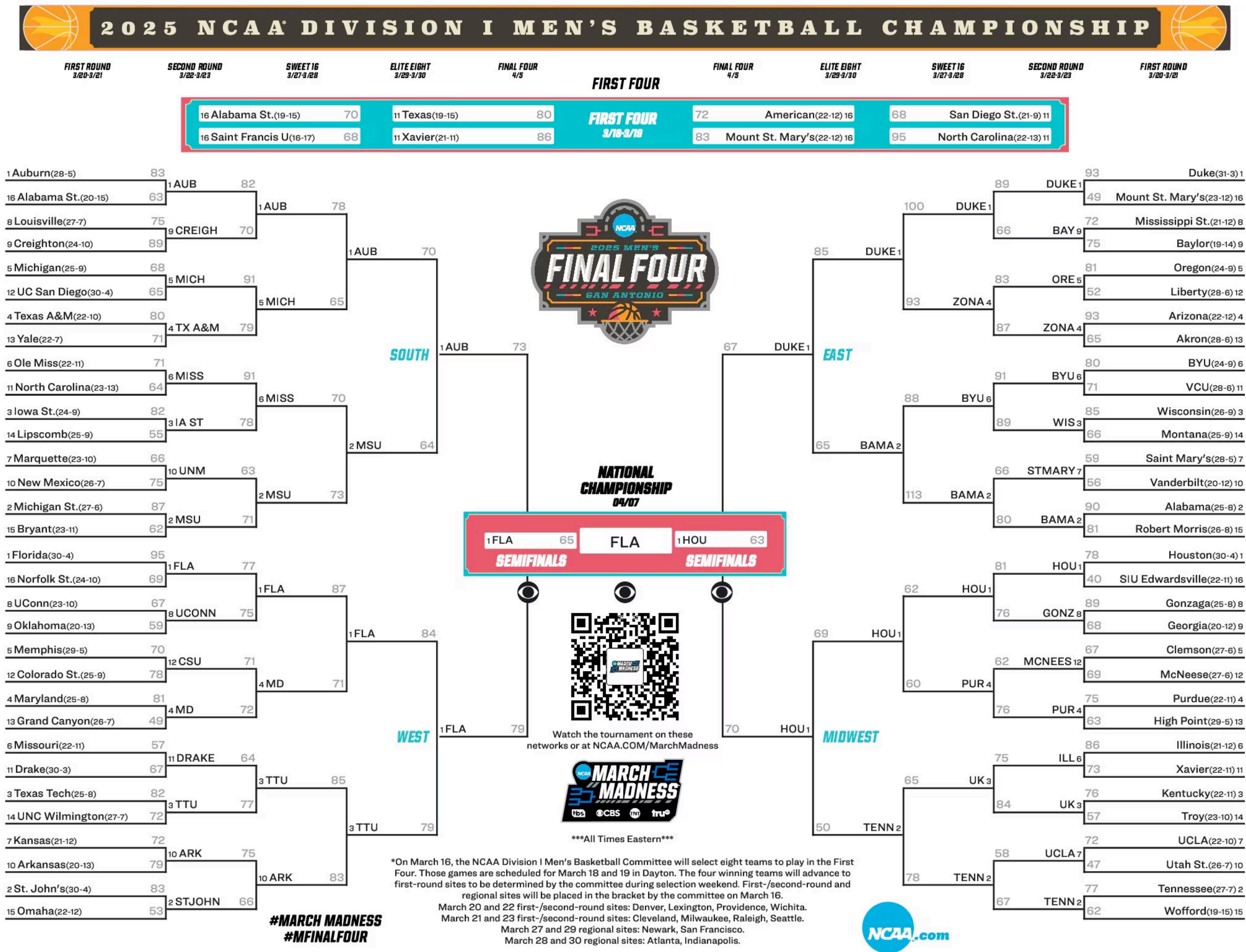
- “In order to get funding, we need to couple our astrophysical searches with a dedicated experimental program.”



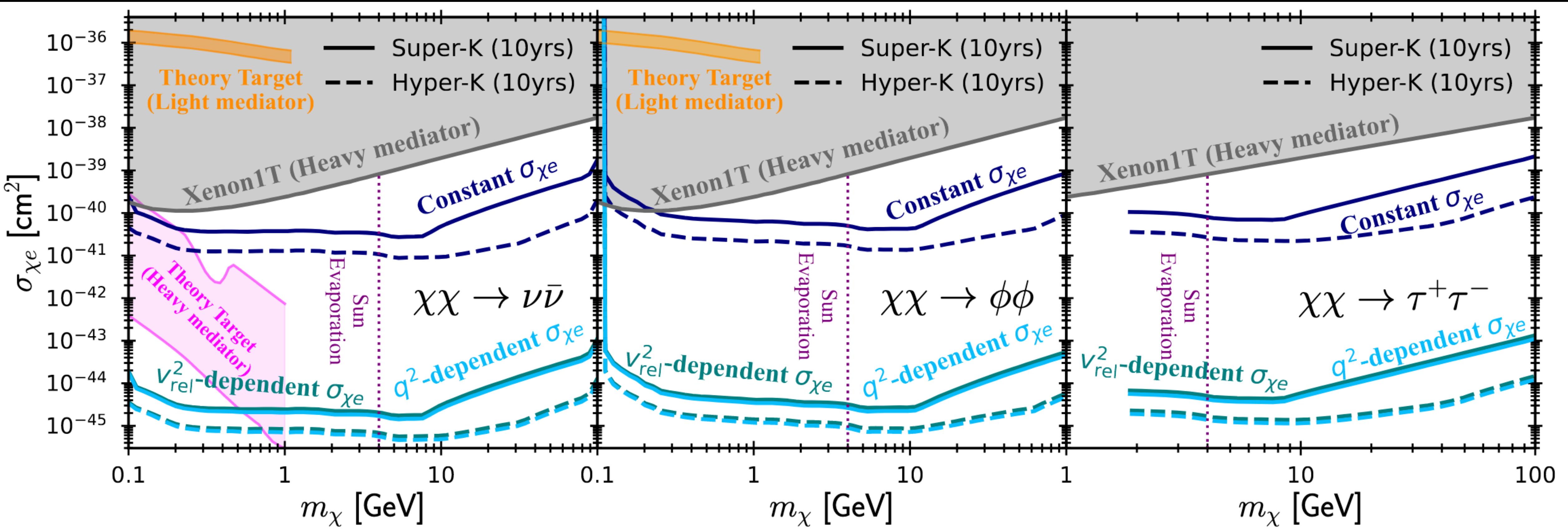
Conclusions

DD	GeV Scale	Sun		     	Heating Signatures (Thermal Spectra)
	Asymmetric	Jupiter		 	Thermal Transport (Seismology)
	Self-Interacting	S-Cluster			Thermal Transport (Affect on Neutrino Flux)
	Light Mediator	G-Objects		  	Thermal Transport (Cooling Curves)
				 	Escaping Signals (DM Neutrinos)
no DD	<GeV Mass			   	Escaping Signals (Cosmic Rays/ γ -Rays)
	>TeV Mass	GC Brown Dwarfs		 	Structural Changes (Longer Lifespan of CBs)
	Spin-Dependent	GC White Dwarfs		   	Structural Changes (Destruction of CBs)
	Inelastic			   	Structural Changes (Spectral Changes)
	Velocity-Dependent	Dark Stars		   	Structural Changes (Population Level)

Conclusions

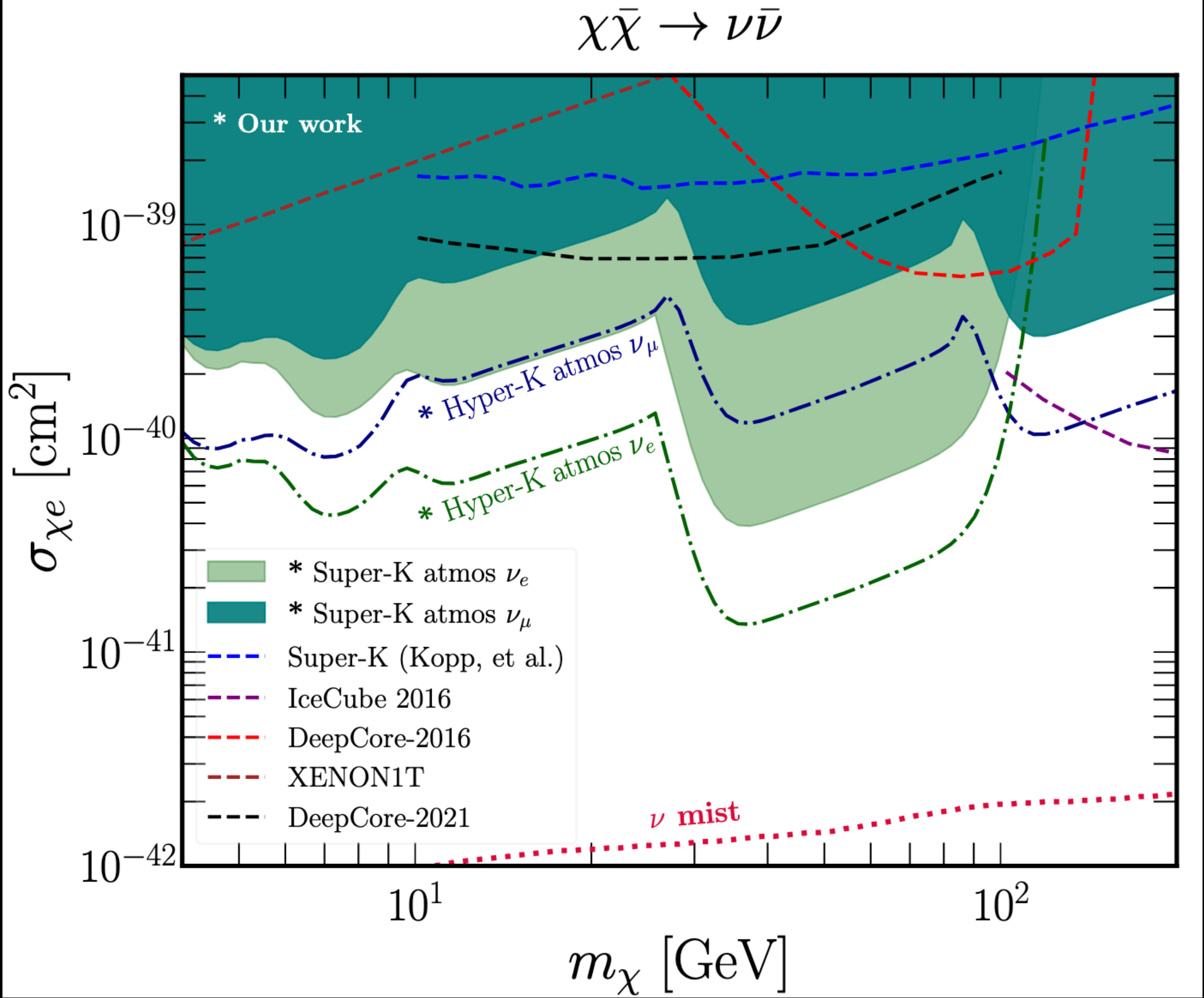


Super-Kamiokande



- When the cross-sections are velocity or momentum dependent, the high velocity of electrons makes constraints much stronger, probing the theoretical targets for leptophilic DM.

A Note Regarding 2503.07713



ν limits from Super-Kamiokande on dark matter-electron scattering in the Sun

Dhashin Krishna ^{1,*}, Rinchen Sherpa ^{1,†}, Akash Kumar Saha ^{1,‡},
^{1,‡} Tarak Nath Maity ^{2,§}, Ranjan Laha ^{1,¶} and Nirmal Raj ^{1,**}

¹ Centre for High Energy Physics, Indian Institute of Science, C. V. Raman Avenue, Bengaluru 560012, India
² Physics, The University of Sydney, ARC Centre of Excellence for Dark Matter Particle Physics, NSW 2006,
(Dated: March 12, 2025)

Dark matter scattering on electrons in the Sun may gravitationally capture and self-annihilate into neutrinos and anti-neutrinos, or other final states that in turn decay to them. The most stringent limits on the fluxes of atmospheric electron-type and muon-type neutrinos by Super-Kamiokande of order 10^{-39} cm² over a mass range of 4–200 GeV. These outdo direct limits from IceCube and previously set limits at IceCube. We also derive limits on dark matter-electron scattering cross sections from atmospheric neutrino observations restricted to the same mass range.

A Note Regarding 2503.07713

Gould, 1987a (Astrophys.J. 321 (1987) 560)

$$R(w \rightarrow v) = \frac{4\mu_+^4}{\pi^{\frac{1}{2}}} N \sigma \frac{v}{w} \int_0^\infty dx \int_{-\infty}^\infty dy \kappa^3 (x+y) e^{-\kappa^2 u^2} \theta(v - |y|) \theta(x-w) \quad (\text{A11})$$

$$= \frac{2}{\pi^{\frac{1}{2}}} \frac{\mu_+^2}{\mu} N \sigma \frac{v}{w} [\chi(-\alpha_-, \alpha_+) + \chi(-\beta_-, \beta_+) e^{-\frac{M}{2T}(v^2 - w^2)}]. \quad (\text{A15})$$

Garani & Palomares-Ruiz (1702.02768)

$$\begin{aligned} R_i(w \rightarrow v) &= \int n_i(r) \frac{d\sigma_i}{dv} |\boldsymbol{w} - \boldsymbol{u}| f_i(\boldsymbol{u}, r) d^3\boldsymbol{u} \\ &= \frac{2}{\sqrt{\pi}} \frac{n_i(r)}{u_i^3(r)} \int_0^\infty du u^2 \int_{-1}^1 d\cos\theta \frac{d\sigma_i}{dv} |\boldsymbol{w} - \boldsymbol{u}| e^{-u^2/u_i^2(r)}, \end{aligned} \quad (\text{A.1})$$

A Note Regarding 2503.07713

$$C = \left[\left(\frac{8}{3\pi} \right)^{\frac{1}{2}} \sigma n_W \bar{v} \right] \left[\frac{M_B}{m} \right] \left[\frac{3v_{\text{esc}}^2}{2\bar{v}^2} \langle \hat{\phi} \rangle \right] [\xi_\eta(\infty)] \left\langle \frac{\hat{\phi}}{\langle \hat{\phi} \rangle} \left(1 - \frac{1 - e^{-A^2}}{A^2} \right) \frac{\xi_1(A)}{\xi_\eta(\infty)} \right\rangle, \quad (2.31)$$

Moreover, the distribution of energy loss is uniform over this interval. On the other hand, scattering from velocity w to a velocity less than v , requires an energy loss of *at least*

$$\frac{\Delta E}{E} \geq \frac{w^2 - v^2}{w^2} = \frac{u^2}{w^2}. \quad (2.11)$$

Combining expressions (2.9) and (2.11) gives the probability that a given scattering will leave the WIMP with less than escape energy,

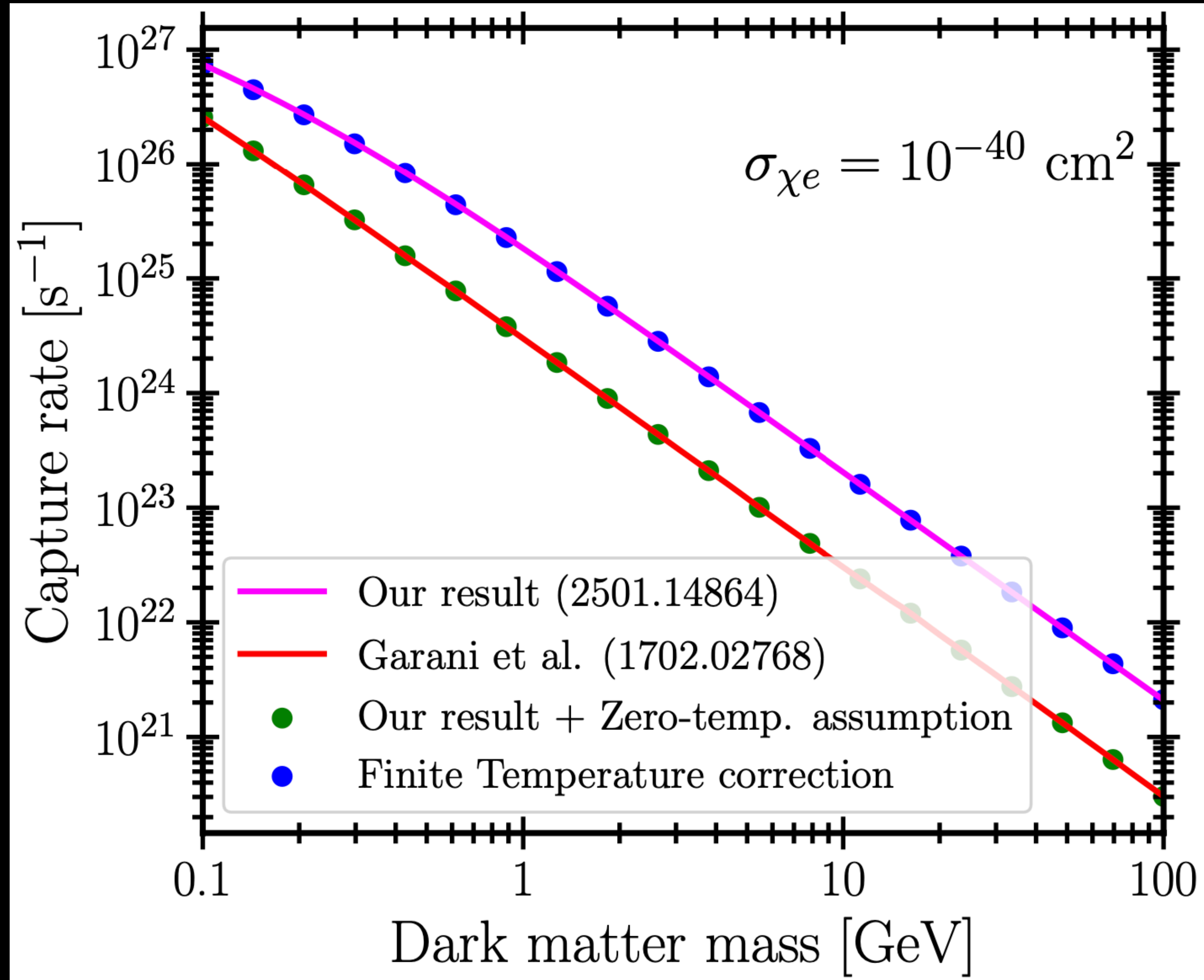
$$\frac{\mu_+^2}{\mu} \cdot \left(\frac{\mu}{\mu_+^2} - \frac{u^2}{w^2} \right) \theta \left(\frac{\mu}{\mu_+^2} - \frac{u^2}{w^2} \right). \quad (2.12)$$

The rate of scattering from w to less than v is just the product of the total rate of scattering, $\sigma n w$, with the conditional probability (2.12). This result may be written,

$$\Omega_v^-(w) = \frac{\sigma n}{w} \left(v^2 - \frac{\mu_-^2}{\mu} u^2 \right) \theta \left(v^2 - \frac{\mu_-^2}{\mu} u^2 \right). \quad (2.13)$$

A Note Regarding 2503.07713

- **Incorrectly adding a zero-temperature kinematic cutoff significantly suppresses the leptophilic dark matter capture rate in the Sun (by a factor of ~7).**
- **Correcting this error leads to stronger limits in many studies.**



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$$A - M \geq -1 \quad \text{Constraints}$$