ev falos: Three Findings And Three Five! Puzzles



PSR B0656+14









EARLY LESSONS

- 1.) Pulsars are highly efficient e⁺e⁻ accelerators.
- 2.) TeV e⁺e⁻ are not confined in the source.



PSR B0656+14



• 3.) Regions near pulsar sources have unusually low diffusion coefficients.

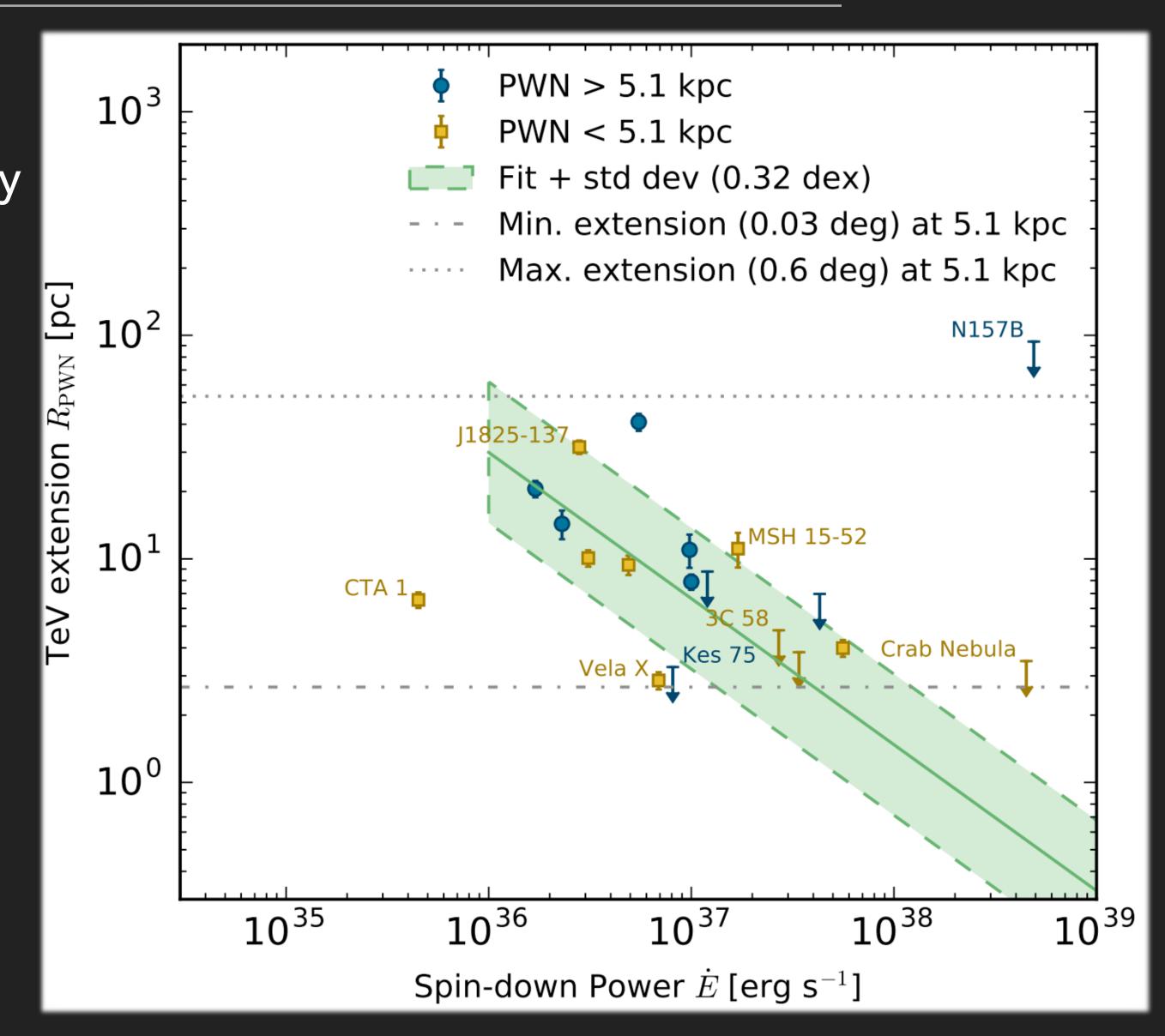
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TEV HALOS: A NEW SOURCE CLASS

TeV Halos are much larger than PWN, especially at low spin down power and large ages.

NOTE: The size of halos has the opposite time- dependence as the X-Ray PWN.

$$\begin{aligned} R_{\rm PWN} \simeq 1.5 \left(\frac{\dot{E}}{10^{35}\,{\rm erg/s}} \right)^{1/2} \times \\ \left(\frac{n_{\rm gas}}{1\,{\rm cm}^{-3}} \right)^{-1/2} \left(\frac{v}{100\,{\rm km/s}} \right)^{-3/2} {\rm pc} \end{aligned}$$

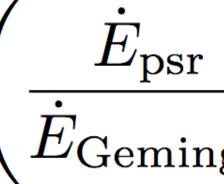


TEV HALOS: THE GEMINGA-CENTRIC MODEL

Make One Key Assumption:

ATNF Name	Dec. ($^{\circ}$)	Distance (kpc)	Age (kyr)	Spindown Lum. (erg s^{-1})	Spindown Flux (erg s ^{-1} kpc ^{-2})	2HWC
J0633+1746	17.77	0.25	342	3.2e34	4.1e34	2HWC J0631+169
B0656+14	14.23	0.29	111	3.8e34	3.6e34	2HWC J0700+143
B1951+32	32.87	3.00	107	3.7e36	3.3e34	
J1740+1000	10.00	1.23	114	2.3e35	1.2e34	
J1913+1011	10.18	4.61	169	2.9e36	1.1e34	2HWC J1912+099
J1831-0952	-9.86	3.68	128	1.1e36	6.4e33	2HWC J1831-098
J2032+4127	41.45	1.70	181	1.7e35	4.7e33	2HWC J2031+415
B1822-09	-9.58	0.30	232	4.6e33	4.1e33	
B1830-08	-8.45	4.50	147	5.8e35	2.3e33	
J1913+0904	9.07	3.00	147	1.6e35	1.4e33	
B0540+23	23.48	1.56	253	4.1e34	1.4e33	

The following correlation is consistent with the data.



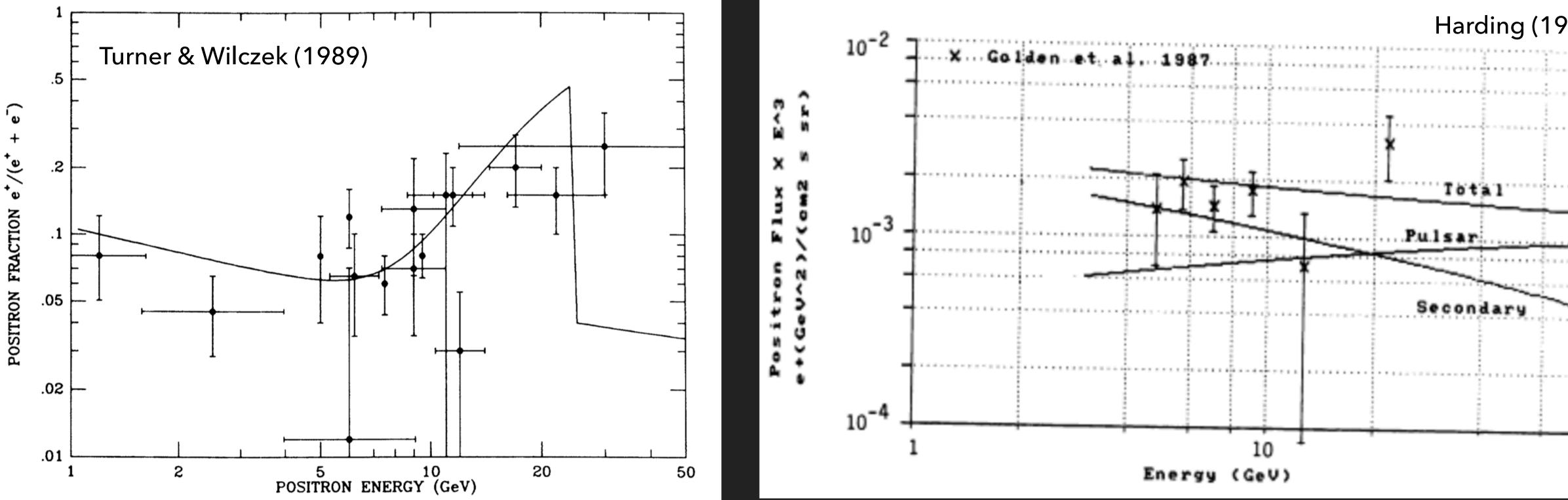


 $\phi_{\rm TeV \ halo} = \left(\frac{\dot{E}_{\rm psr}}{\dot{E}_{\rm Geminga}}\right) \left(\frac{d_{\rm Geminga}^2}{d_{\rm psr}^2}\right) \phi_{\rm Geminga}$

IMPLICATION 1: THE POSITRON EXCESS

EARLY LESSONS

• 1.) Pulsars are highly efficient e⁺e⁻ accelerators.



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IMPLICATION 1: THE POSITRON EXCESS

• What were the uncertainties in pulsar scenarios of the positron excess?

• I: The e⁺e⁻ production efficiency?

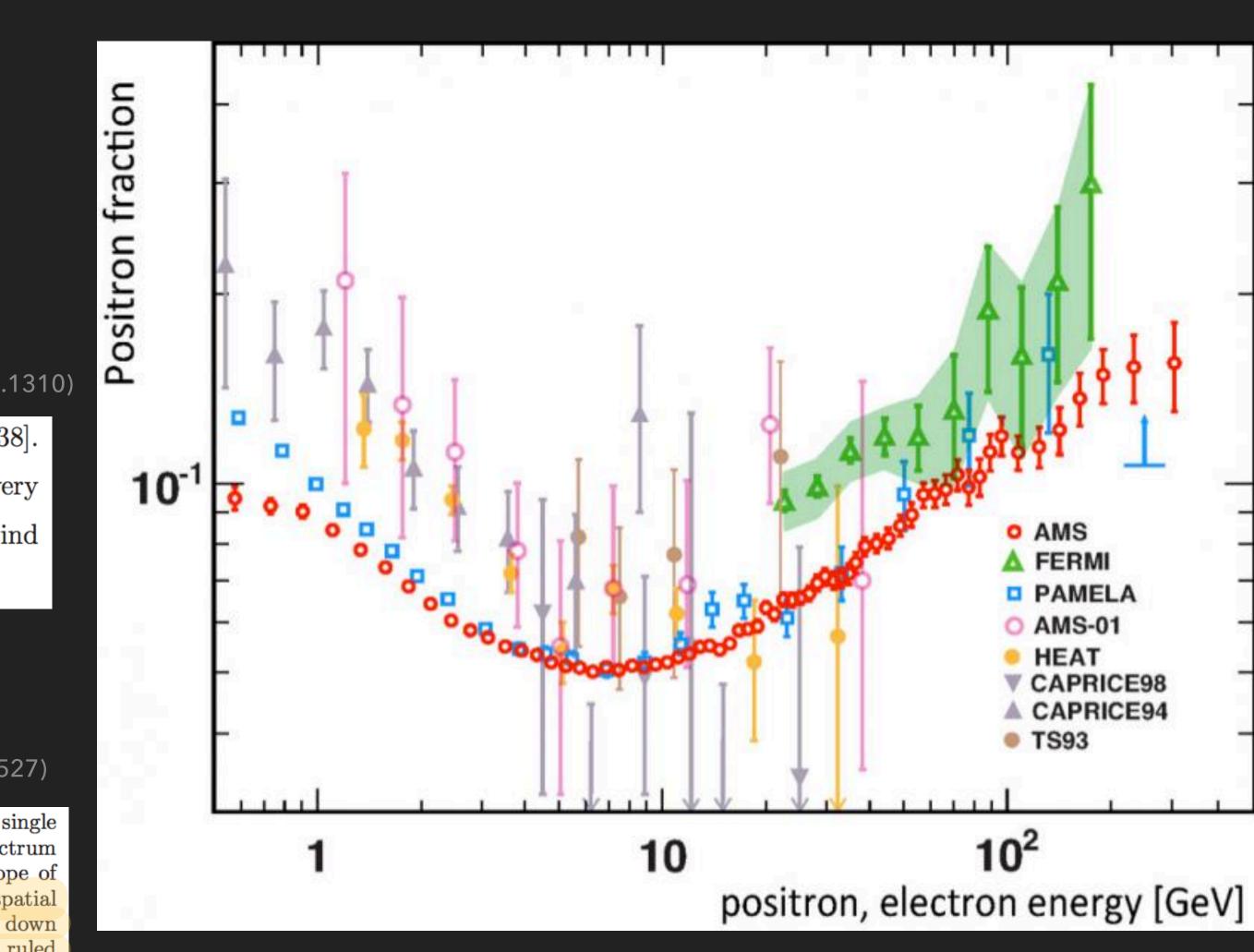
Profumo (0812.4457); Malyshev et al. (0903.1310)

%. A quantitative discussion of plausible values for $f_{e^{\pm}}$ was recently given in Ref. [38]. We shall not review their discussion here, but Ref. [38] argues (see in particular their very informative App. B and C) that in the context of a standard model for the pulsar wind nebulae, a reasonable range for $f_{e^{\pm}}$ falls between 1% and 30%.

• II: The e⁺e⁻ spectrum.

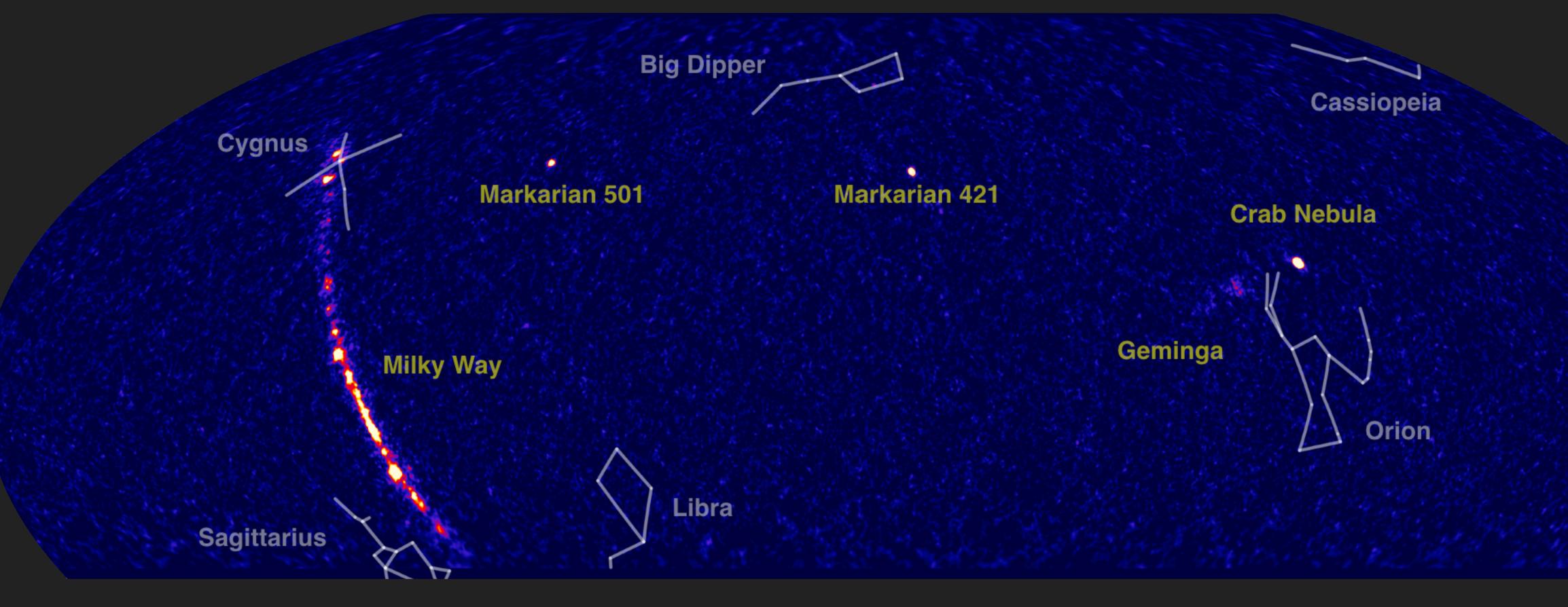
Hooper et al. (0810.1527)

part of their energy adiabatically because of the expansion of the wind. The energy spectrum injected by a single pulsar depends on the environmental parameters of the pulsar, but some attempts to calculate the average spectrum injected by a population of mature pulsars suggest that the spectrum may be relatively hard, having a slope of $\sim 1.5-1.6$ [18]. This spectrum, however, results from a complex interplay of individual pulsar spectra, of the spatial and age distributions of pulsars in the Galaxy, and on the assumption that the chief channel for pulsar spin down is magnetic dipole radiation. Due to the related uncertainties, variations from this injection spectra cannot be ruled out. Typically, one concentrates the attention on pulsars of age $\sim 10^5$ years because younger pulsars are likely to still



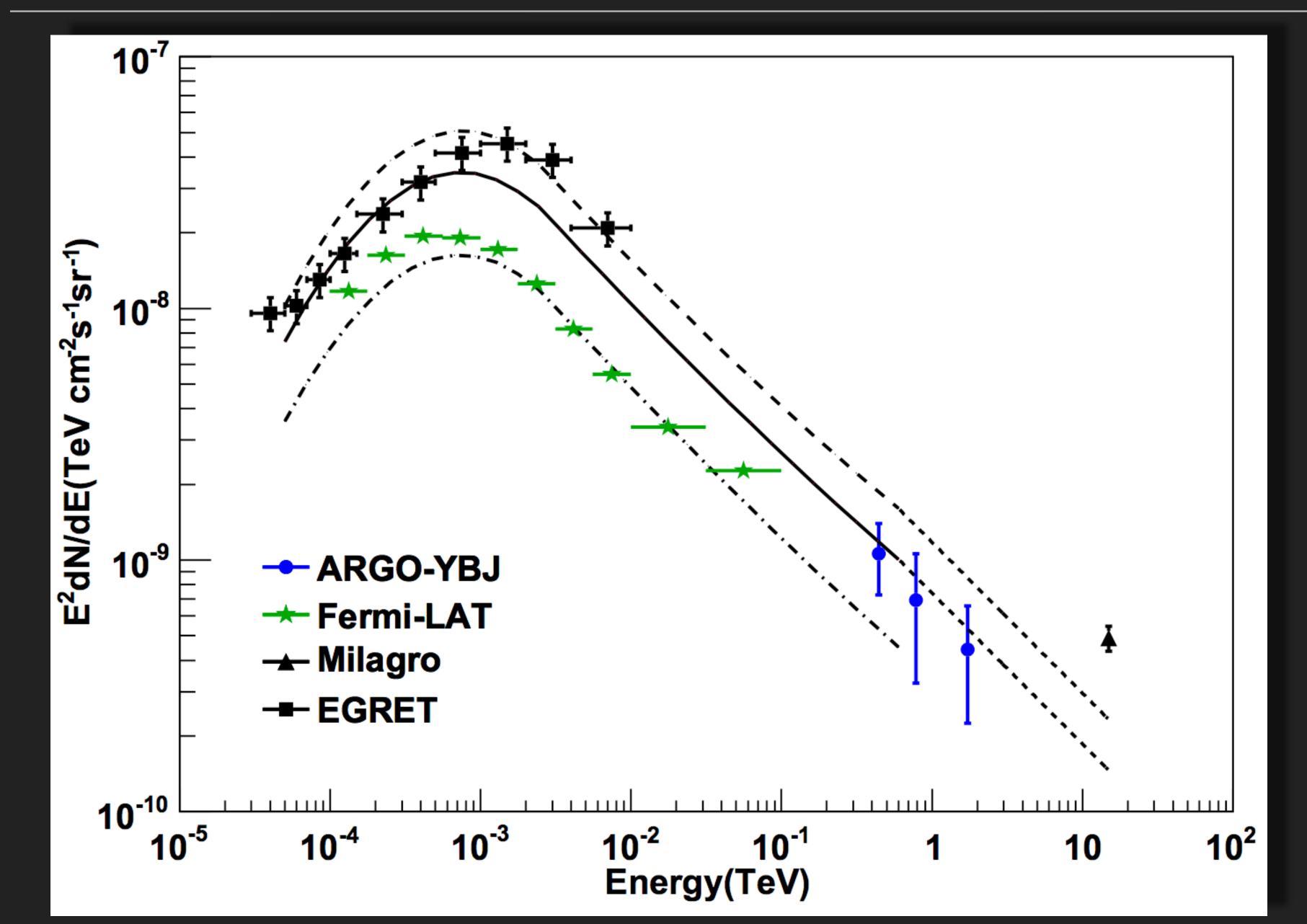
EARLY LESSONS

- 2.) TeV e⁺e⁻ are not confined in the source.



• 1.) Pulsars are highly efficient e⁺e⁻ accelerators.

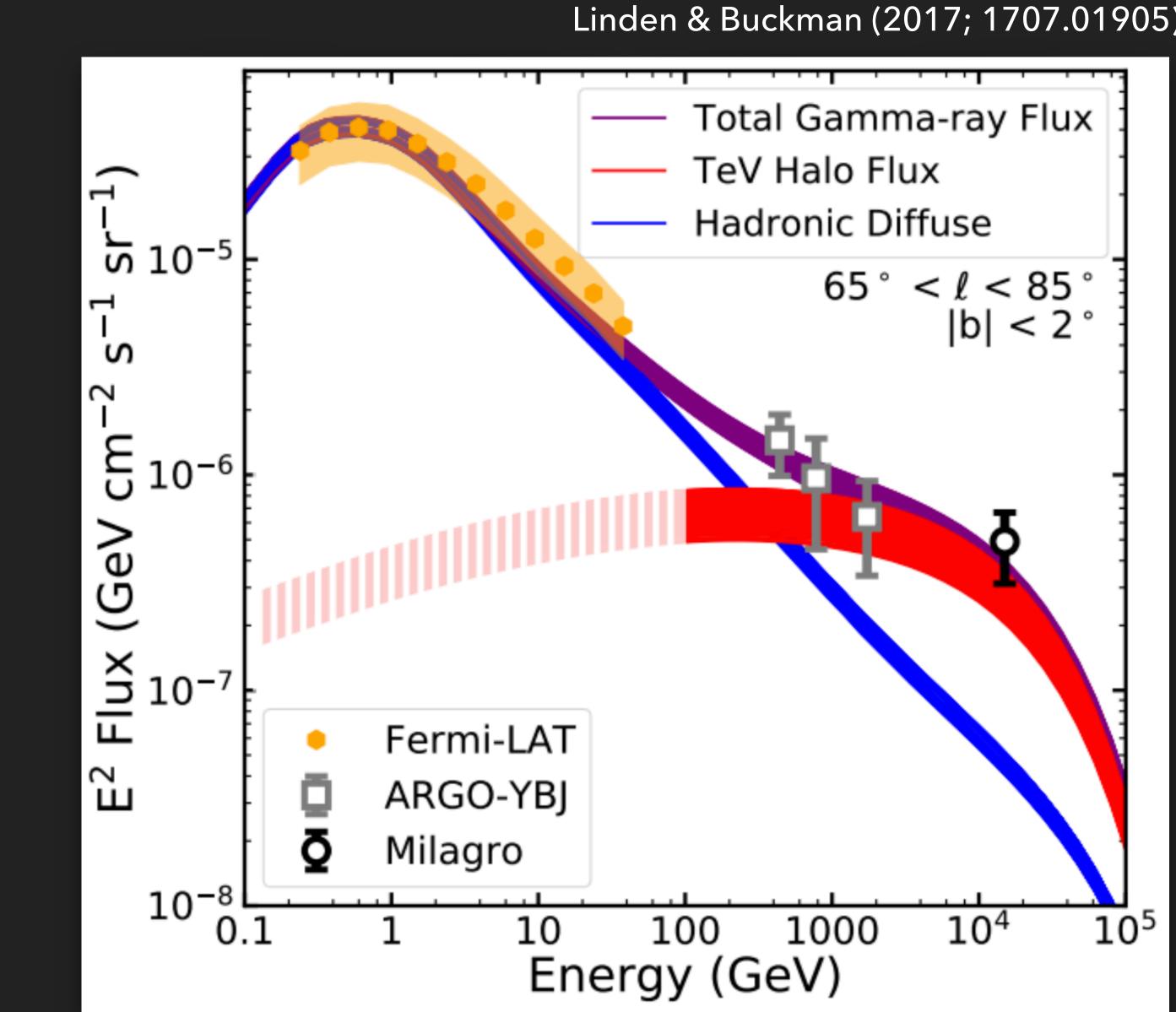




- TeV halos naturally explain the spectrum and intensity of this emission.
- Multiple halos observed with E^{-2.0} spectra.

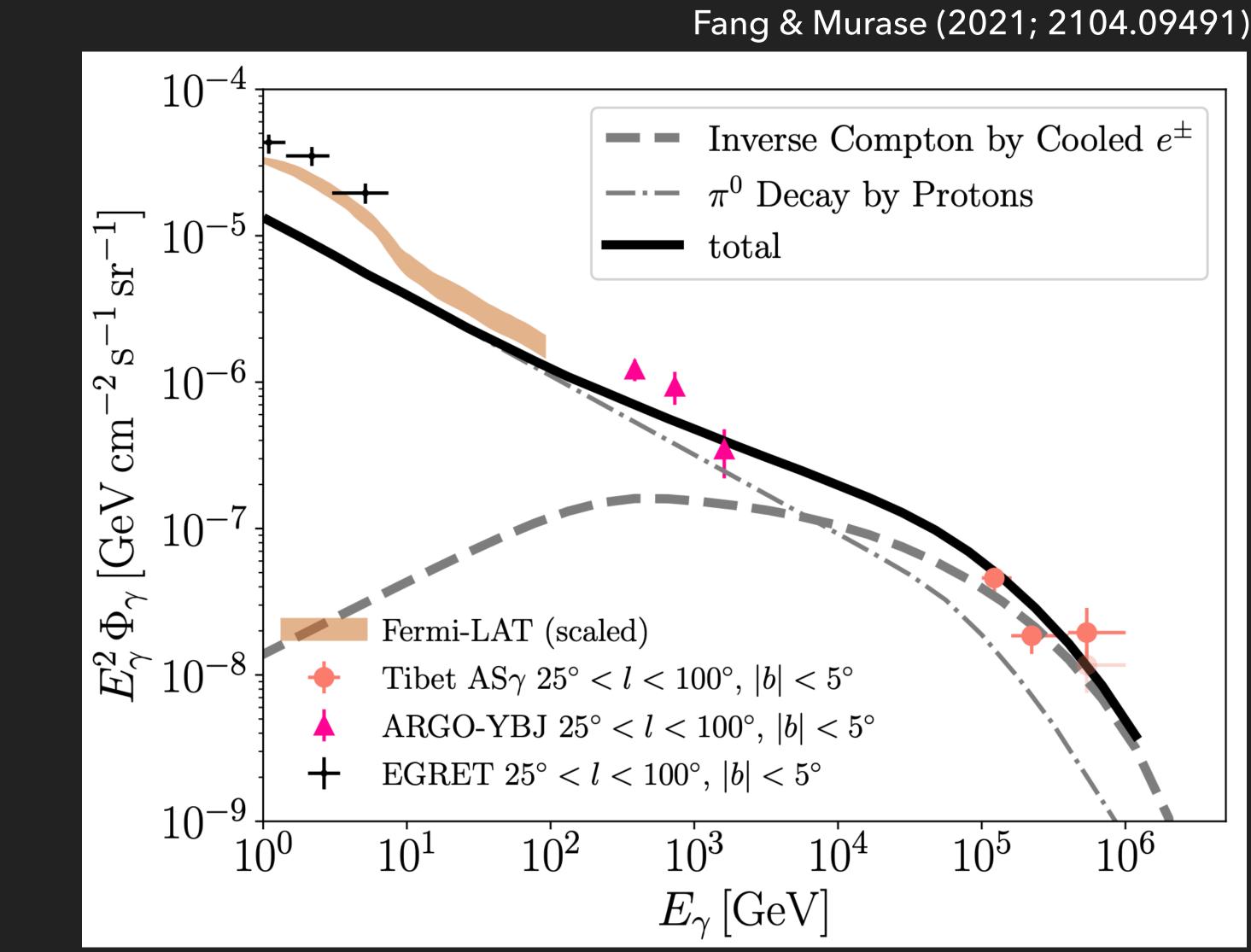
- Note "Halo" is not needed
 - Pulsar efficiency ~10%
 - Power must escape PWN





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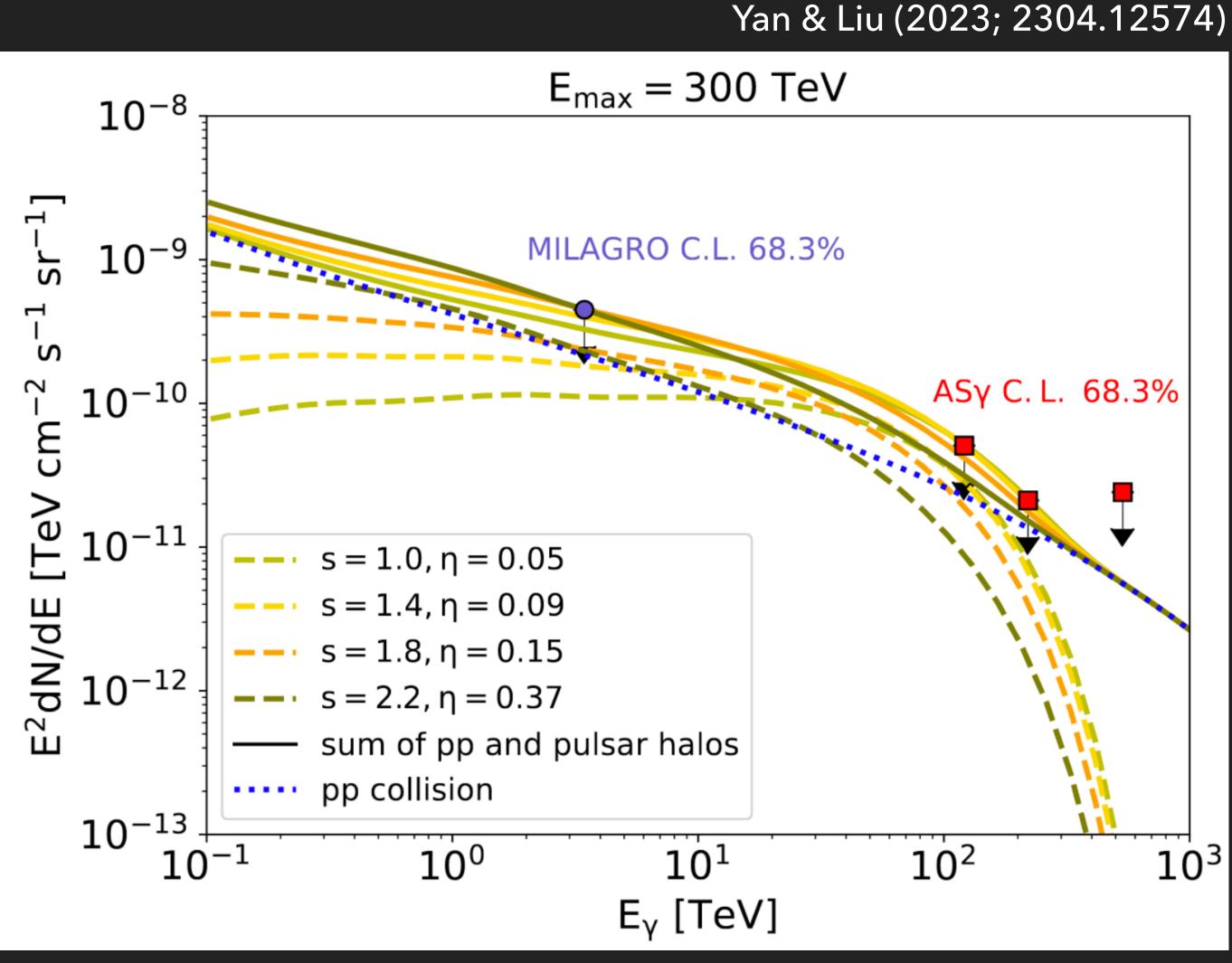
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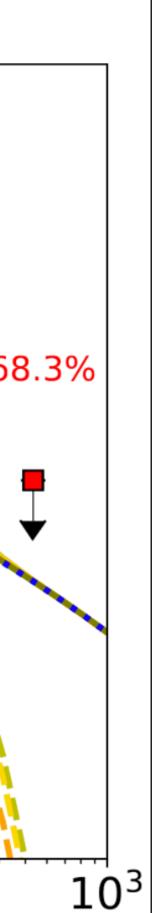
Tibet ASγ data

- TeV halos naturally explain the spectrum and intensity of this emission.
- Multiple halos observed with E^{-2.0} spectra.

- Note "Halo" is not needed
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LHAASO Data



EARLY LESSONS

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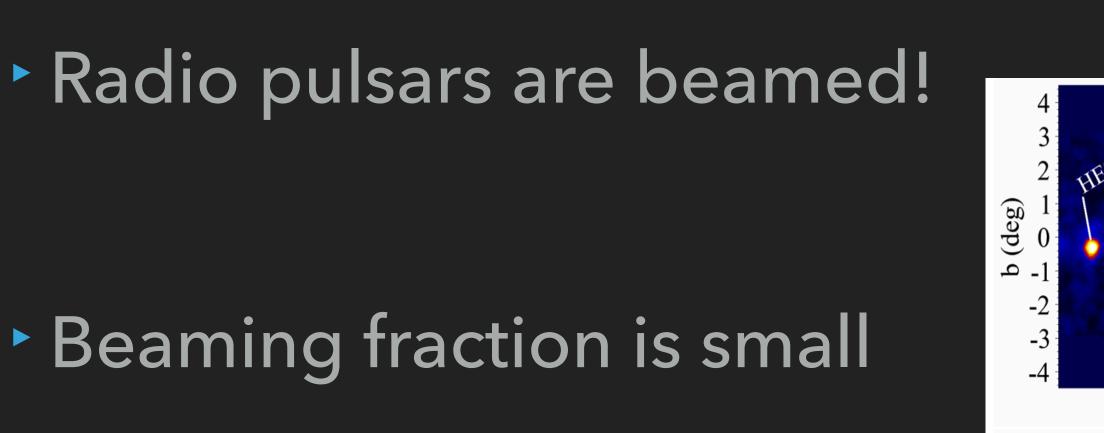
PSR B0656+14



• 3.) Regions near pulsar sources have unusually low diffusion coefficients.

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IMPLICATION 3: MOST TEV SOURCES ARE POWERED BY PULSARS

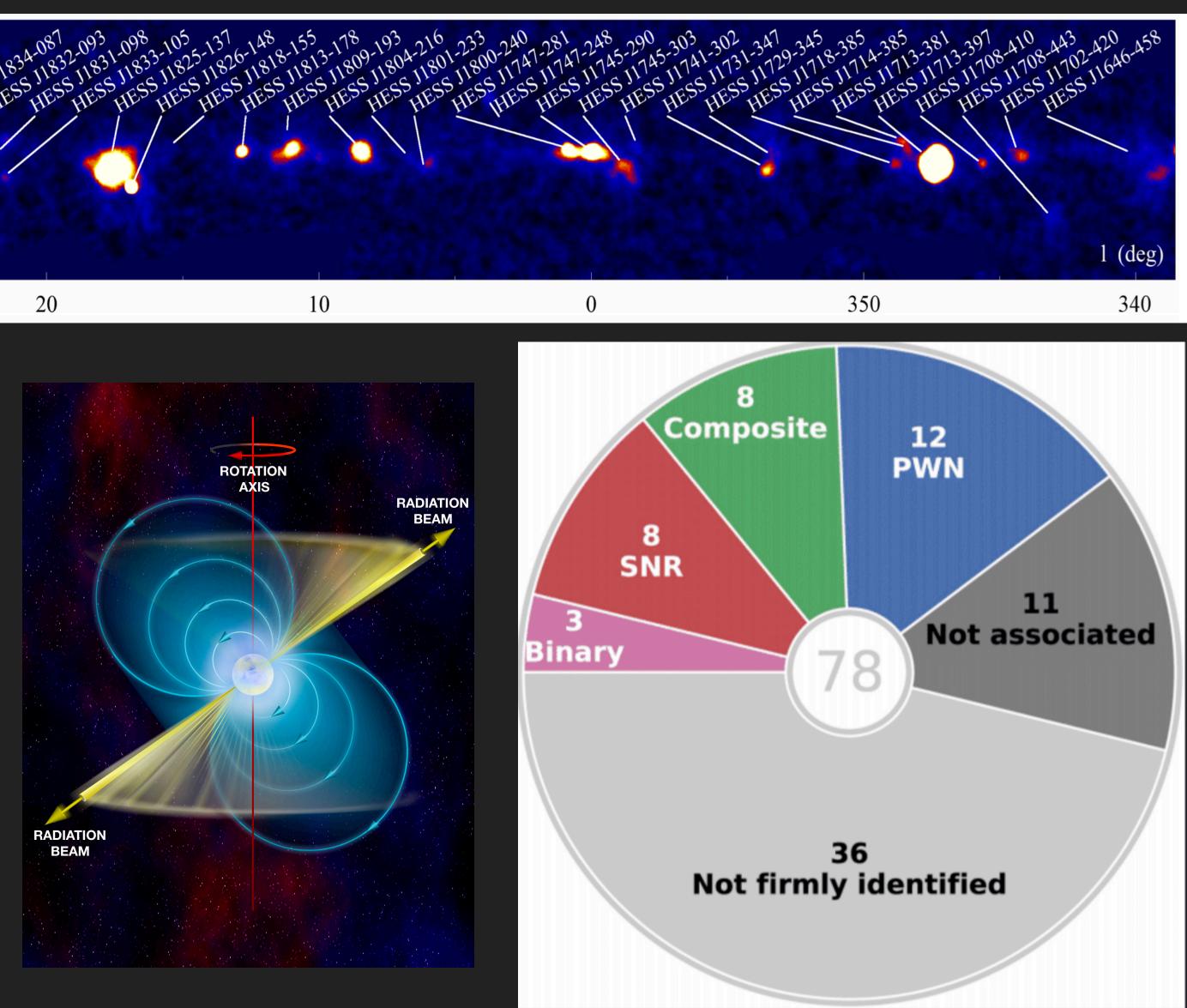


Tauris & Manchester (1998)

$$f = \left[1.1 \left(\log_{10} \left(\frac{\tau}{100 \text{ Myr}}\right)\right)^2 + 15\right]\%$$

This varies between 15-30%.

Most pulsars are unseen in radio!



Lessons from HAWC PWNe observations: the diffusion constant is not a constant; Pulsars remain the likeliest sources of the anomalous positron fraction; Cosmic rays are trapped for long periods of time in pockets of inefficient diffusion

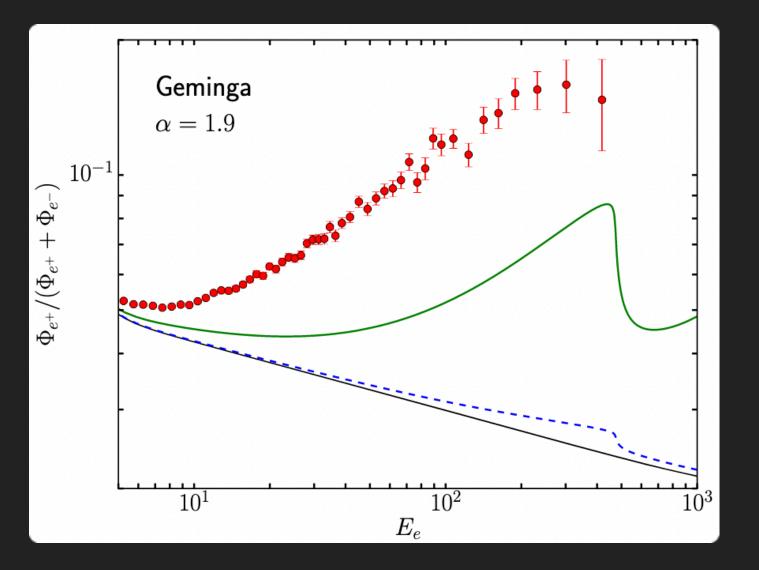
¹Department of Physics, University of California, 1156 High St. Santa Cruz, CA 95060, United States of America

Stefano Profumo,^{1,2,*} Javier Reynoso-Cordova,^{2,3,†} Nicholas Kaaz,^{1,‡} and Maya Silverman^{1,§}

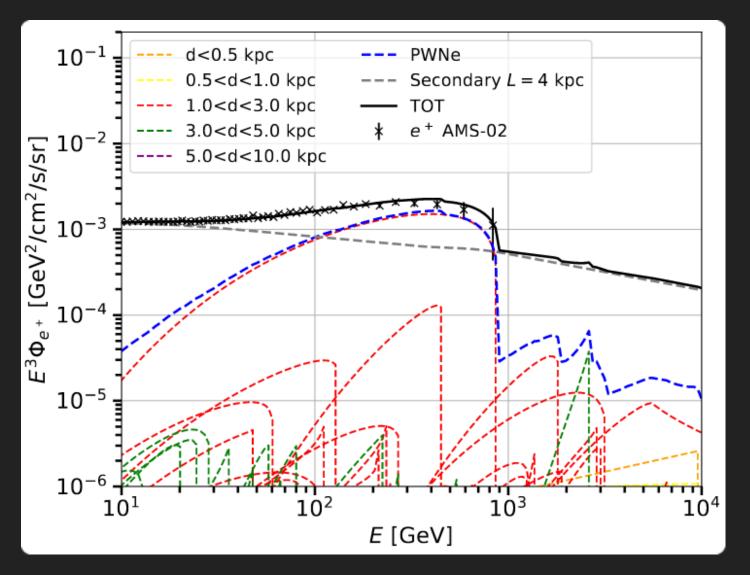
²Santa Cruz Institute for Particle Physics, 1156 High St. Santa Cruz, CA 95060, United States of America ³ Departamento de Física, DCI, Campus León, Universidad de Guanajuato, 37150, León, Guanajuato, México

Recent TeV observations of nearby pulsars with the HAWC telescope have been interpreted as evidence that diffusion of high-energy electrons and positrons within pulsar wind nebulae is highly inefficient compared to the rest of the interstellar medium. If the diffusion coefficient well outside the nebula is close to the value inferred for the region inside the nebula, high-energy electrons and positrons produced by the two observed pulsars could not contribute significantly to the local measured cosmic-ray flux. The HAWC collaboration thus concluded that, under the assumption of isotropic and homogeneous diffusion, the two pulsars are ruled out as sources of the anomalous high-energy positron flux. Here, we argue that since the diffusion coefficient is likely *not* spatially homogeneous, the assumption leading to such conclusion is flawed. We solve the diffusion equation with a radially dependent diffusion coefficient, and show that the pulsars observed by HAWC produce potentially perfect matches to the observed high-energy positron fluxes. We also study the implications of inefficient diffusion within pulsar wind nebulae on Galactic scales, and show that cosmic rays are likely to have very long residence times in regions of inefficient diffusion. We describe how this prediction can be tested with studies of the diffuse Galactic emission.

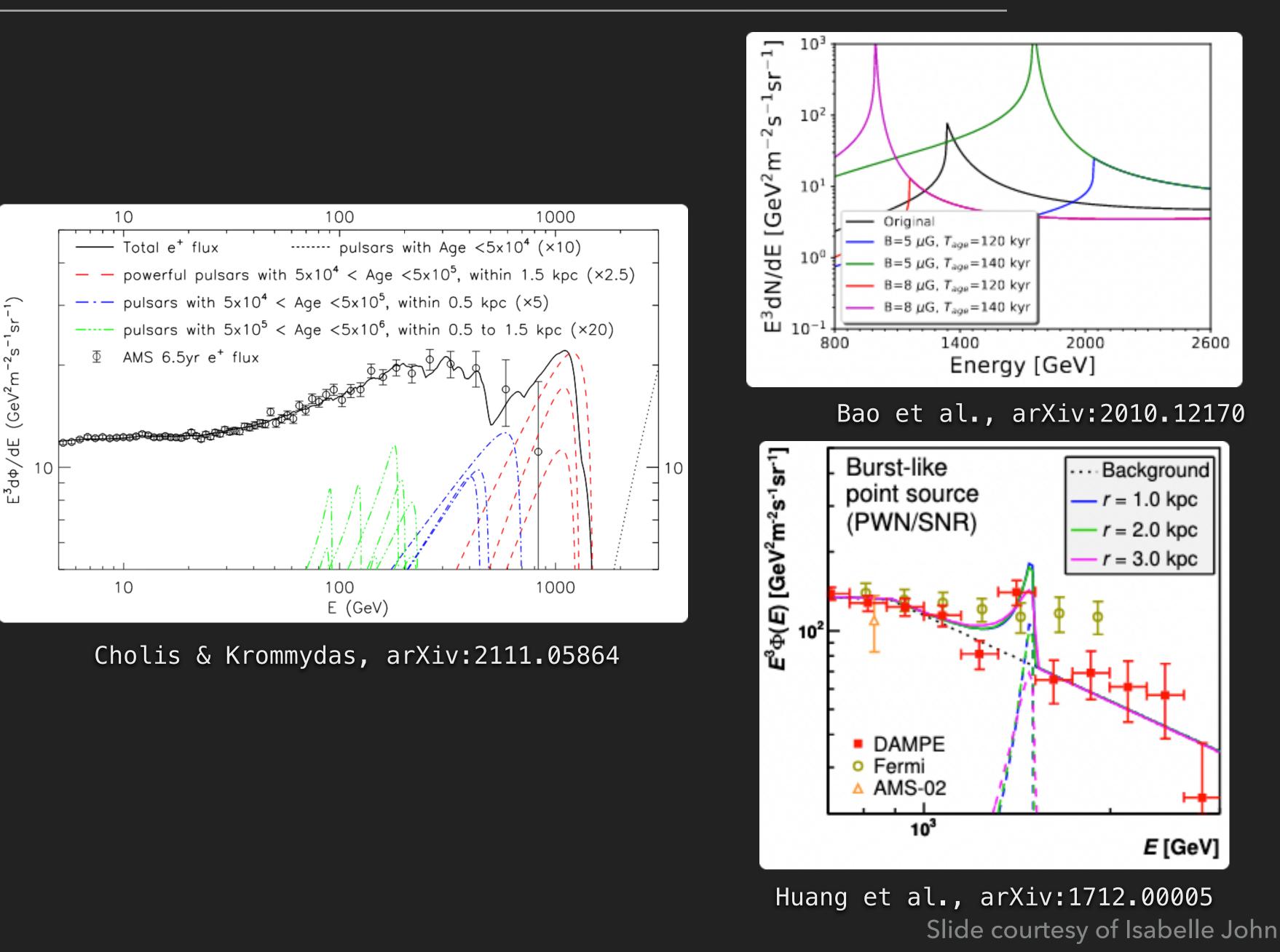
PACS numbers:



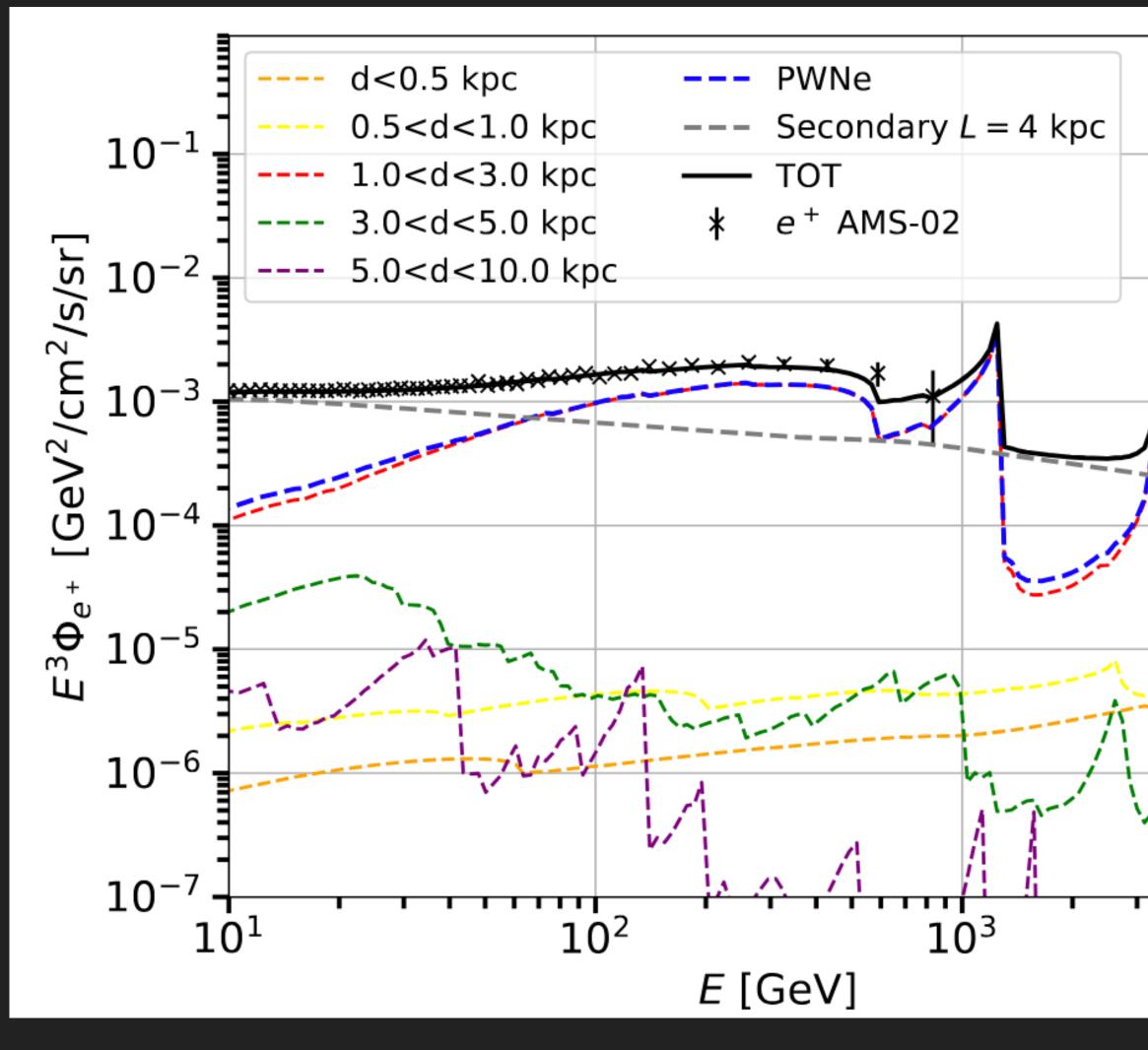
Hooper et al., arXiv:1702.08436



Orusa et al., arXiv:2107.06300

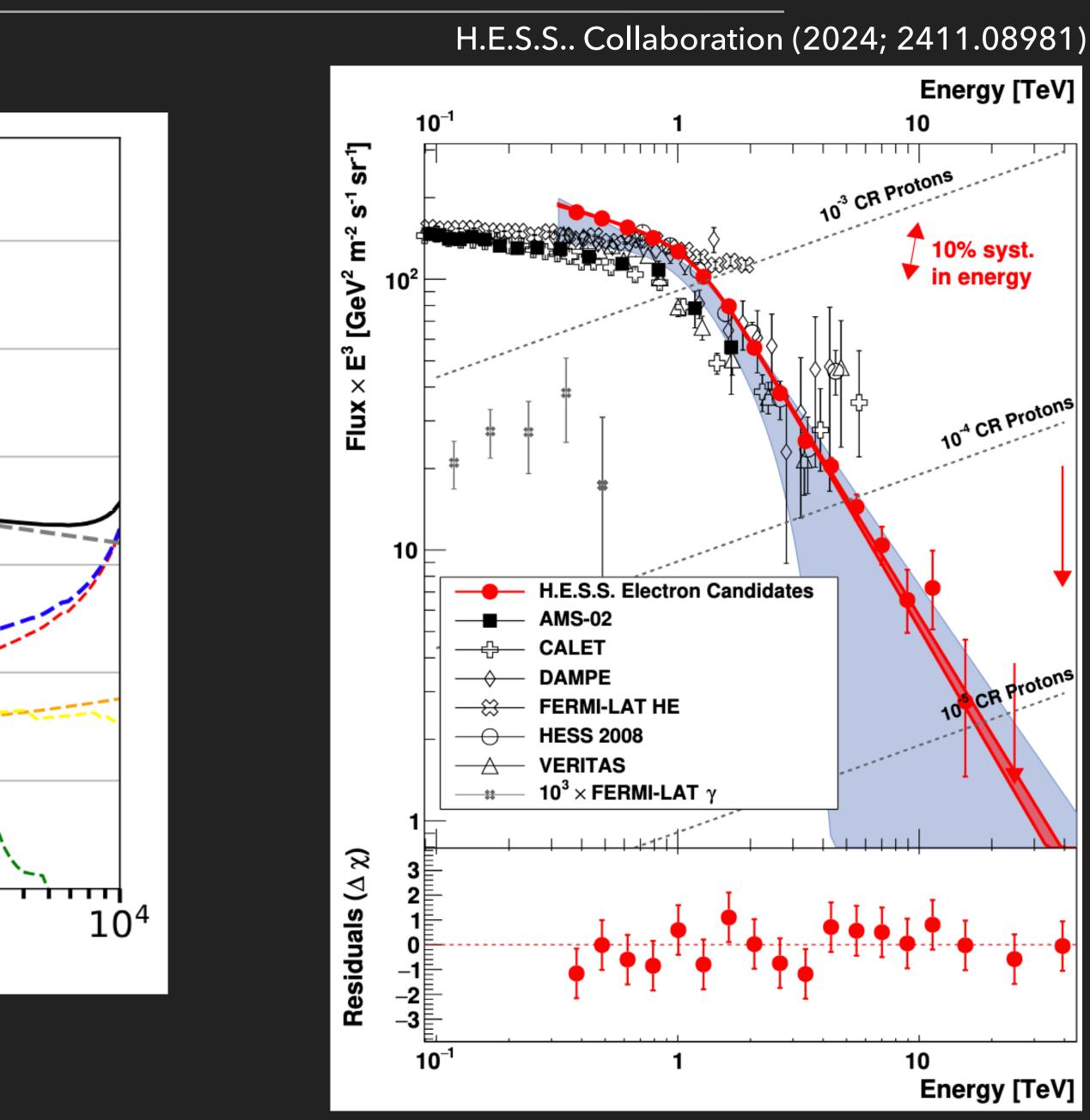


• Debates on the Number of Pulsars

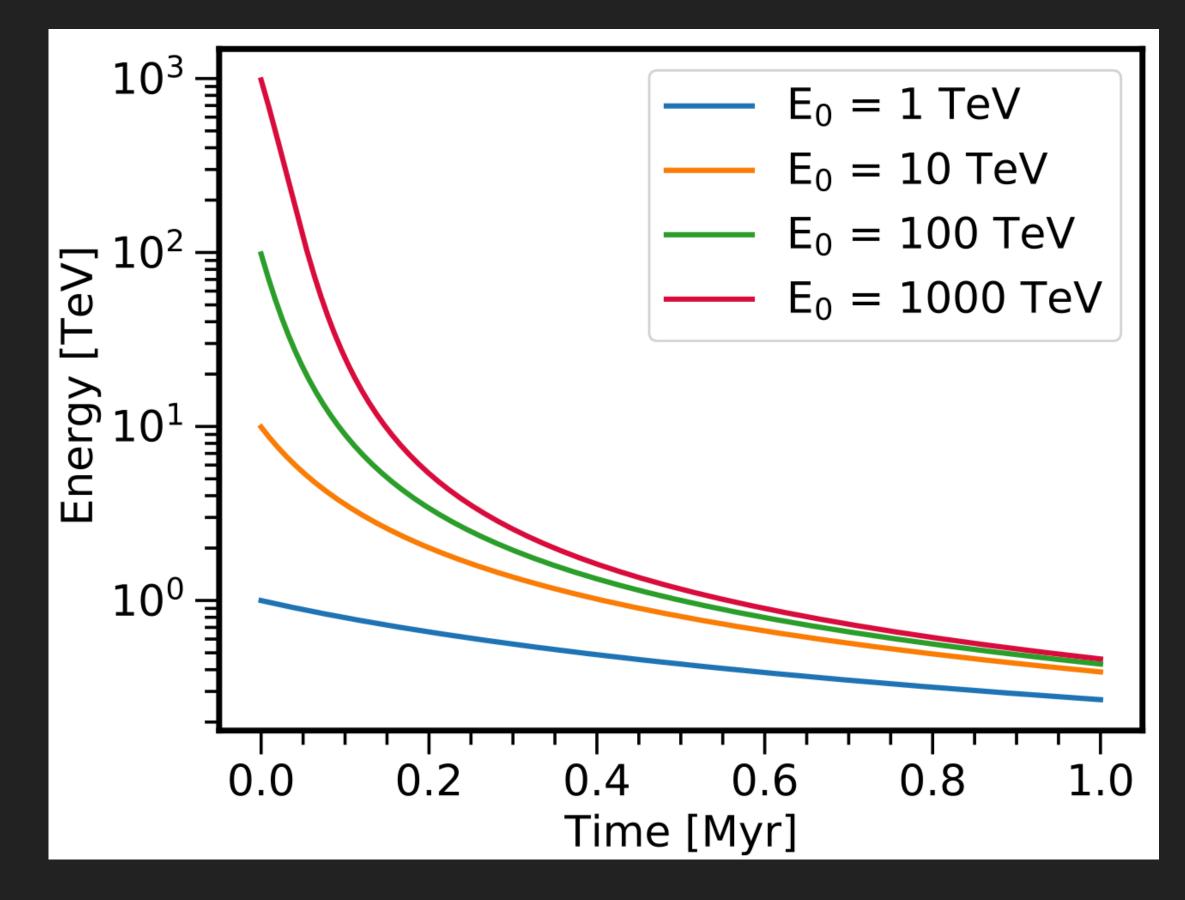


Orusa et al. (2021; 2107.06300)



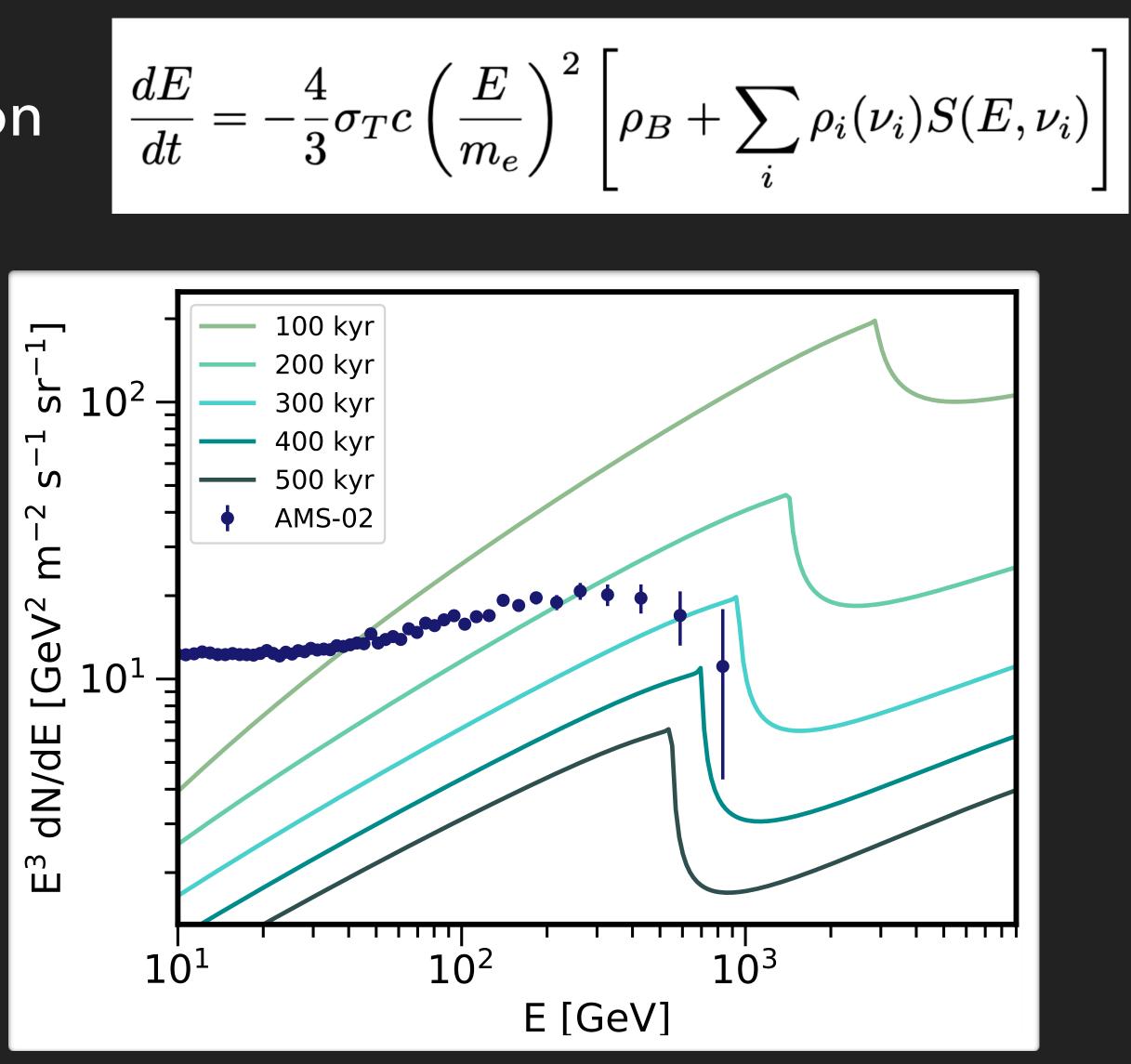


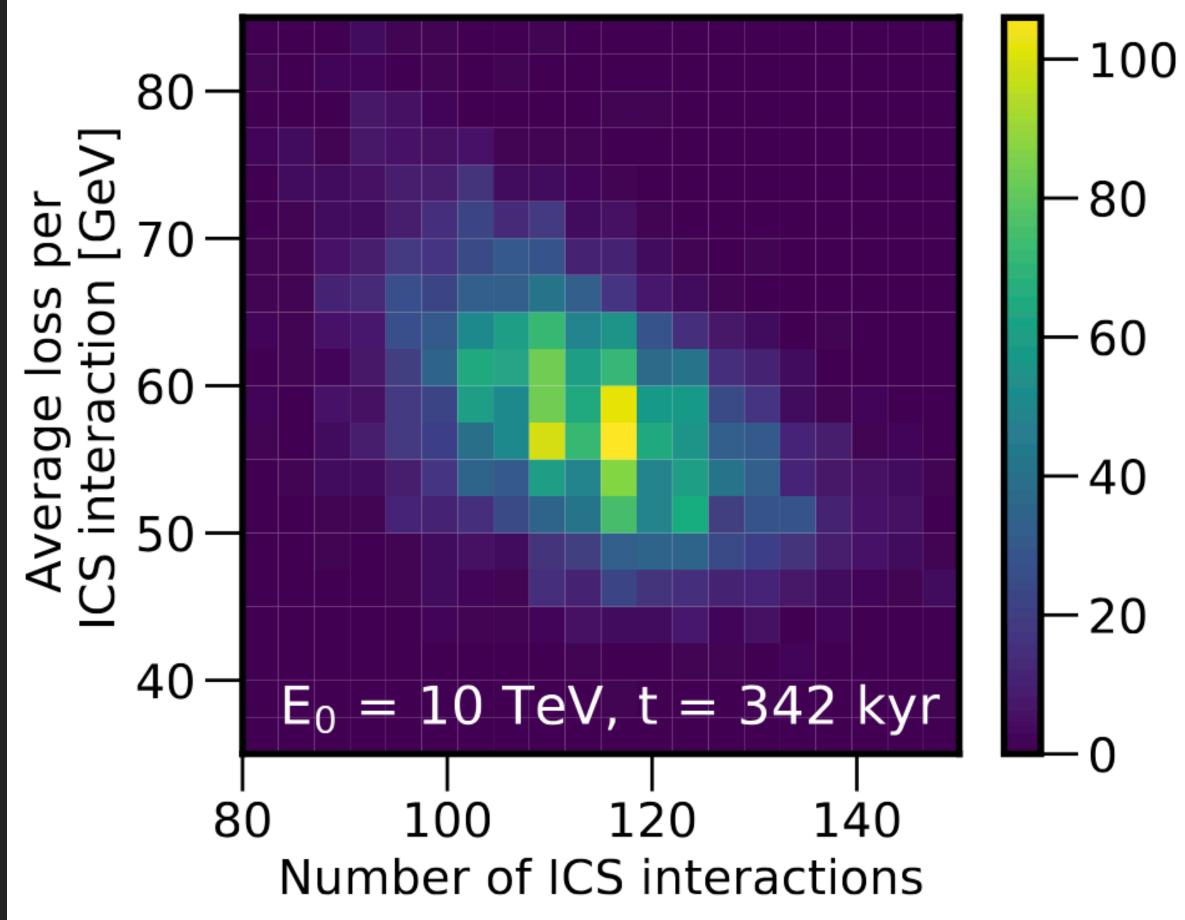
For pulsars, there is a key error: Studies generally use a continuous approximation for electron energy losses:



John & Linden (2022; 2206.04699)

$$\frac{dE}{dt} = -\frac{4}{3}\sigma_T c \left(\frac{E}{m_e}\right)^2 \left[\rho_B + \sum_i \rho_i(\nu_i) S(I_i)\right]$$

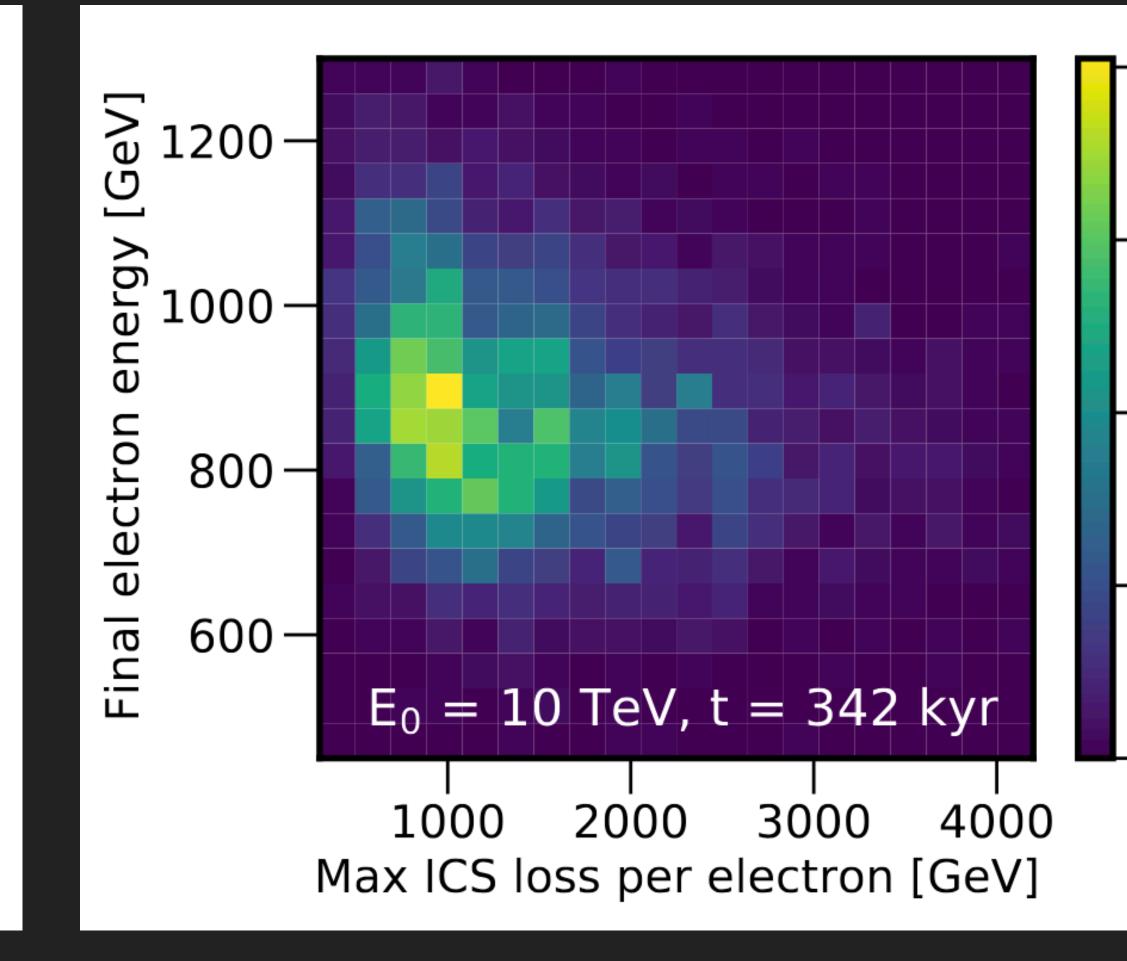




 But ICS interactions are very rare and stochastic. The energy after a given time is not determined by the initial energy.

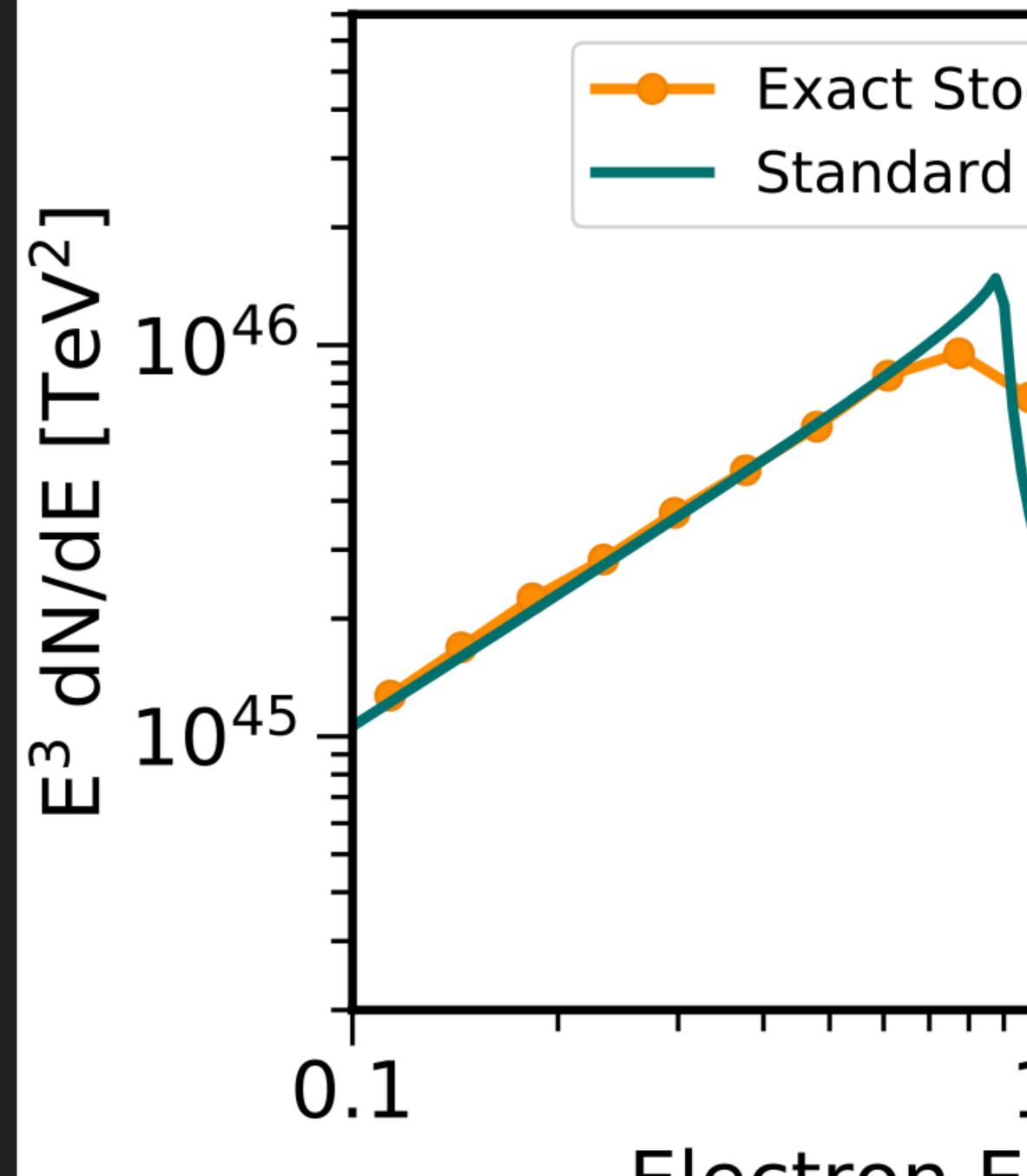
John & Linden (2022; 2206.04699)

Number of counts









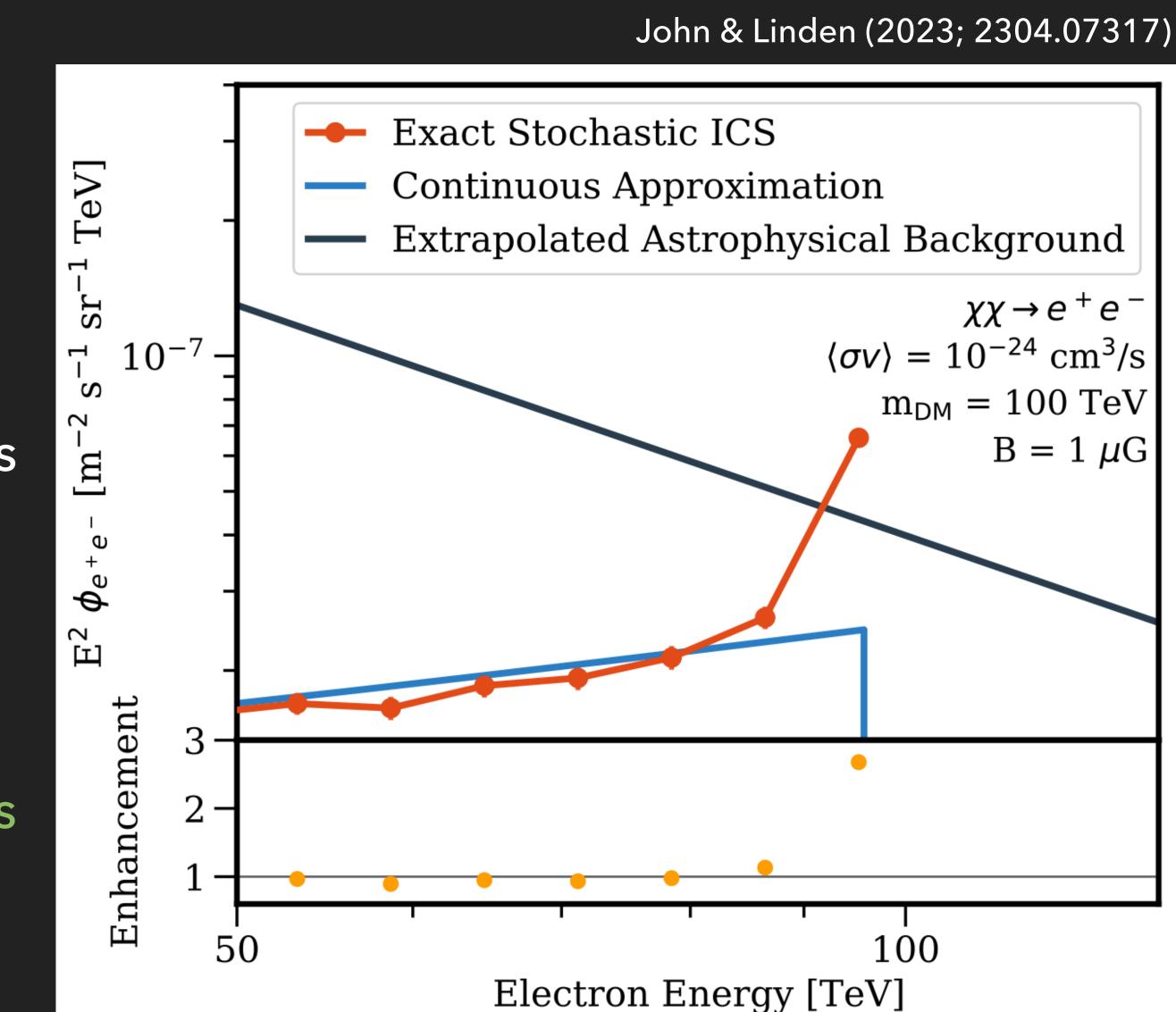
Exact Stochastic ICS Standard Analytic Approximation 10Electron Energy [TeV]

PUZZLE II: DARK MATTER VS. PULSARS IN THE POSITRON DATA

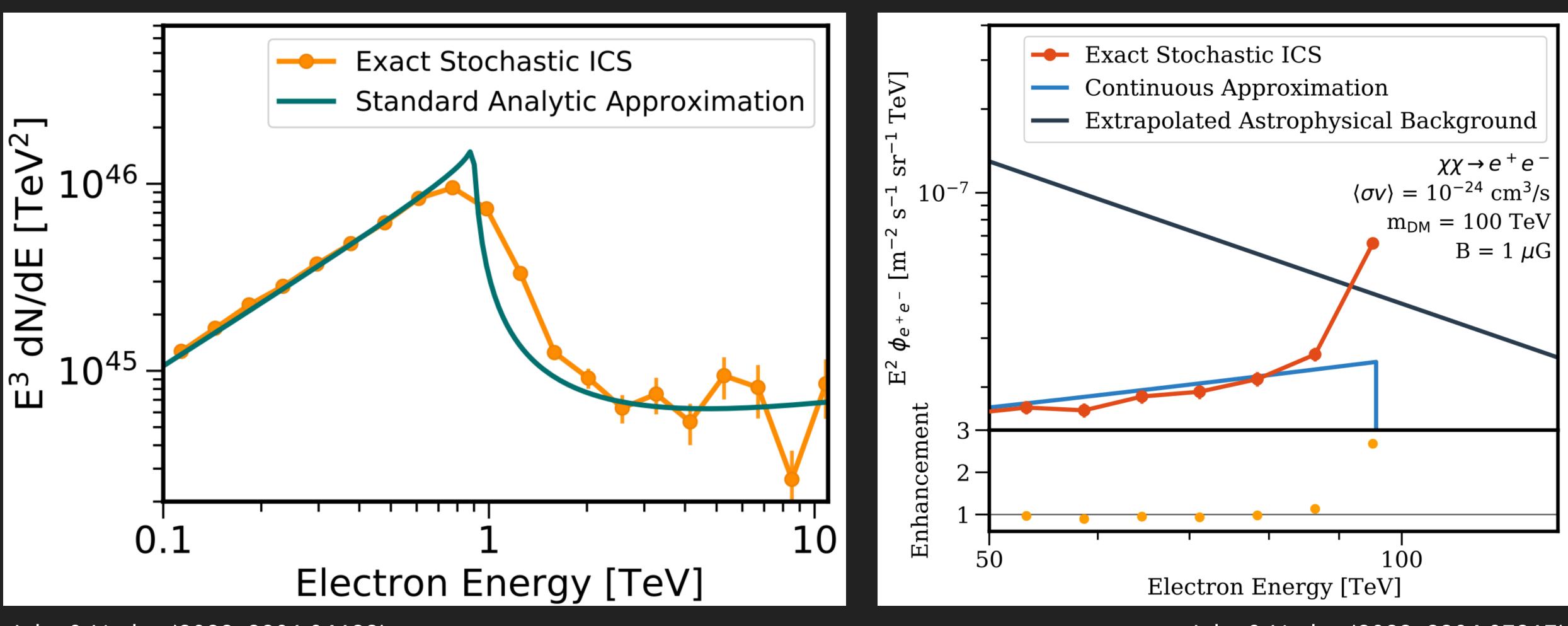
For dark matter, the spectral cutoff is not produced by ICS cooling, but from the dark matter mass.

The stochasticity of cooling instead means that some particles don't cool at all, enhancing the peak.

Correctly accounting for ICS energy losses makes it possible to differentiate dark matter and pulsars via their positron spectrum.



PUZZLE II: DARK MATTER VS. PULSARS IN THE POSITRON DATA



John & Linden (2022; 2206.04699)

ONLY DARK MATTER CAN PRODUCE SHARP SPECTRA IN THE POSITRON DATA!

John & Linden (2023; 2304.07317)





PUZZLE III: COMPLEX HALOS



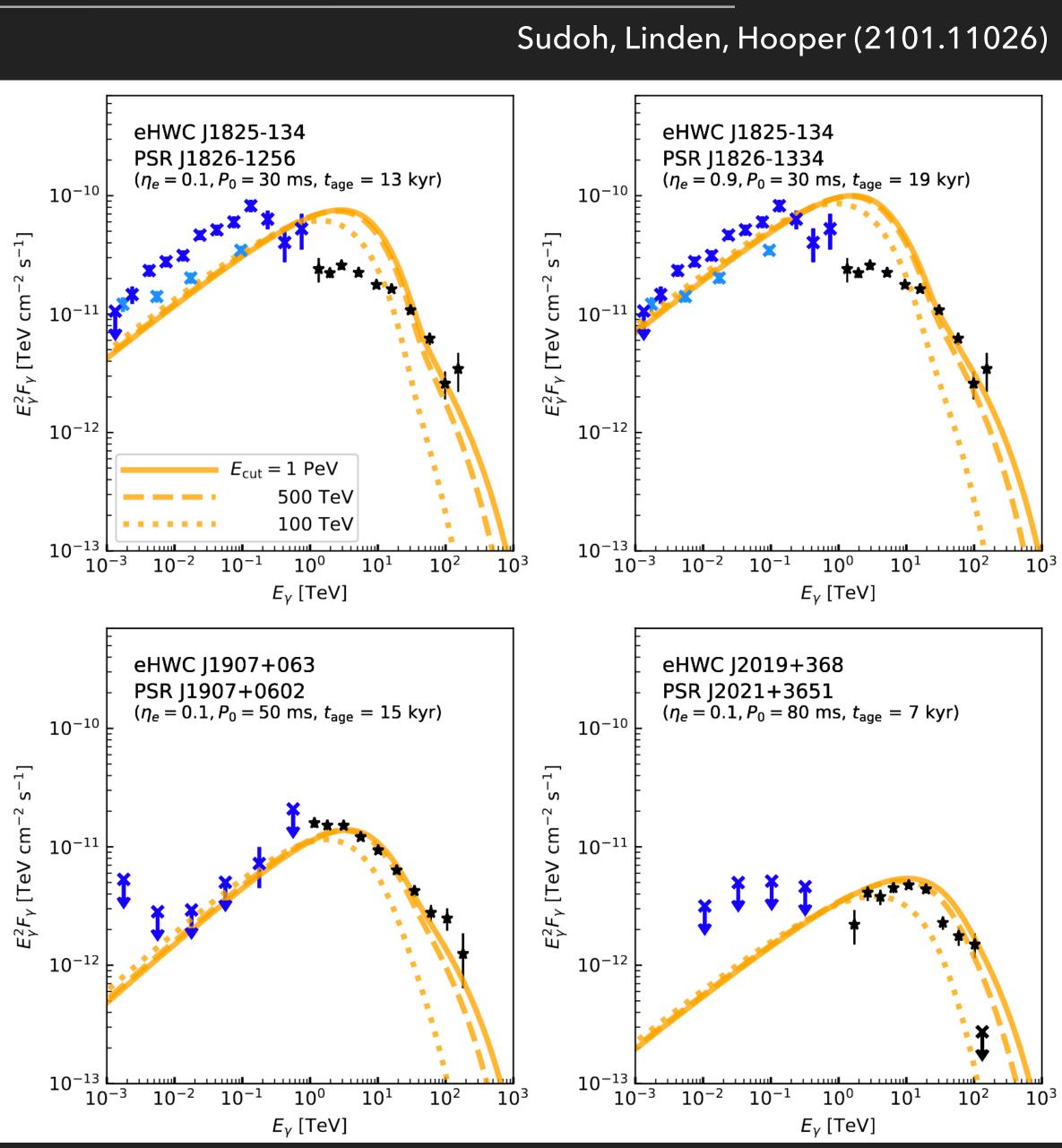


PUZZLE IIIA: COMPOSITE OBJECTS

Most of the highest energy HAWC sources have positions consistent with pulsars.

Ages only 7-20 kyr.

Interplay between PWN and Halo is of critical importance.

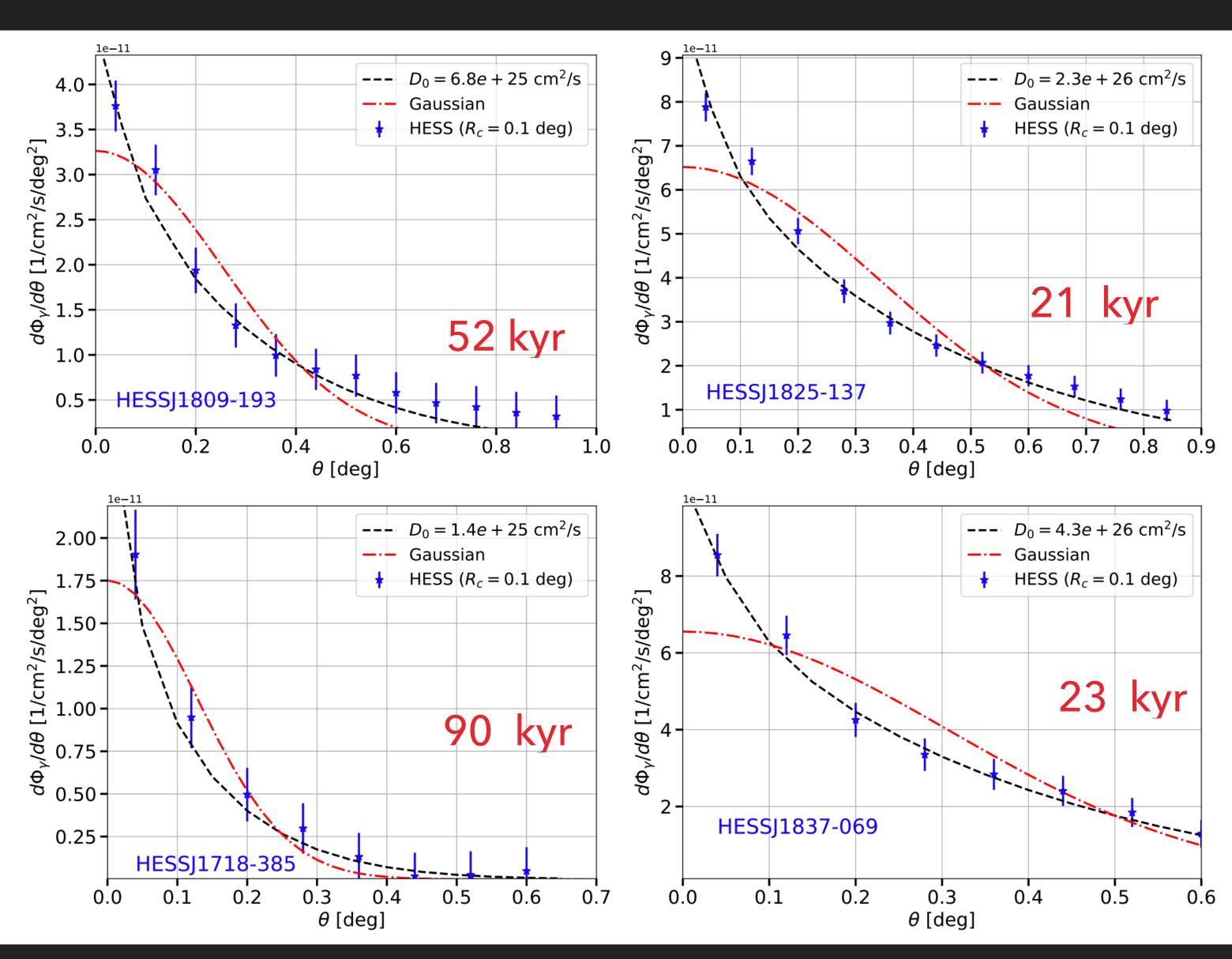


PUZZLE IIIA: COMPOSITE OBJECTS

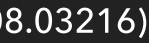
Most of the highest energy HAWC sources have positions consistent with pulsars.

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Di Mauro, Manconi, Donato (2019; 1908.03216)



PUZZLE IIIB: MIRAGE HALOS

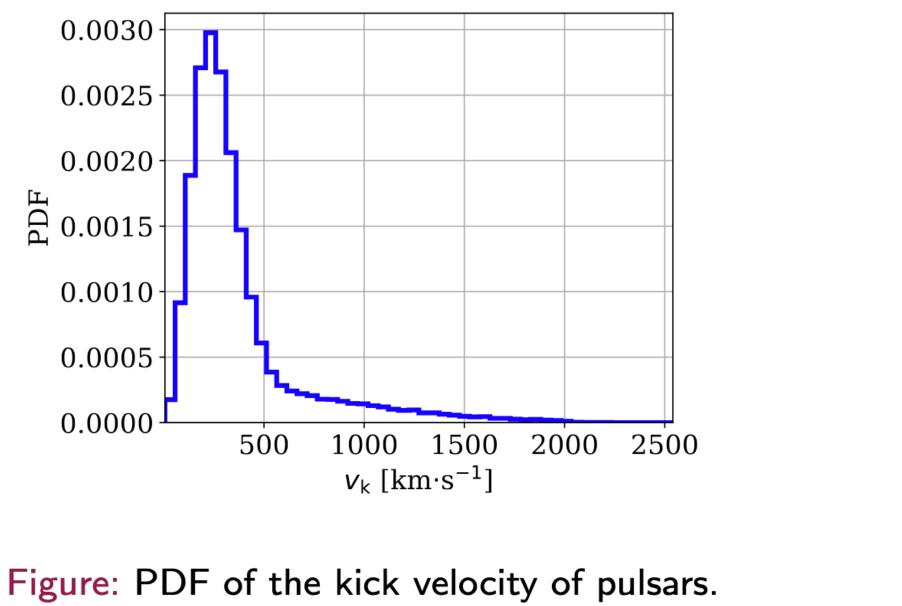
Property of the pulsars: Kick velocity

Kick velocity distribution

Taken from Faucher-Giguère et al. (2006), modulus of all components:

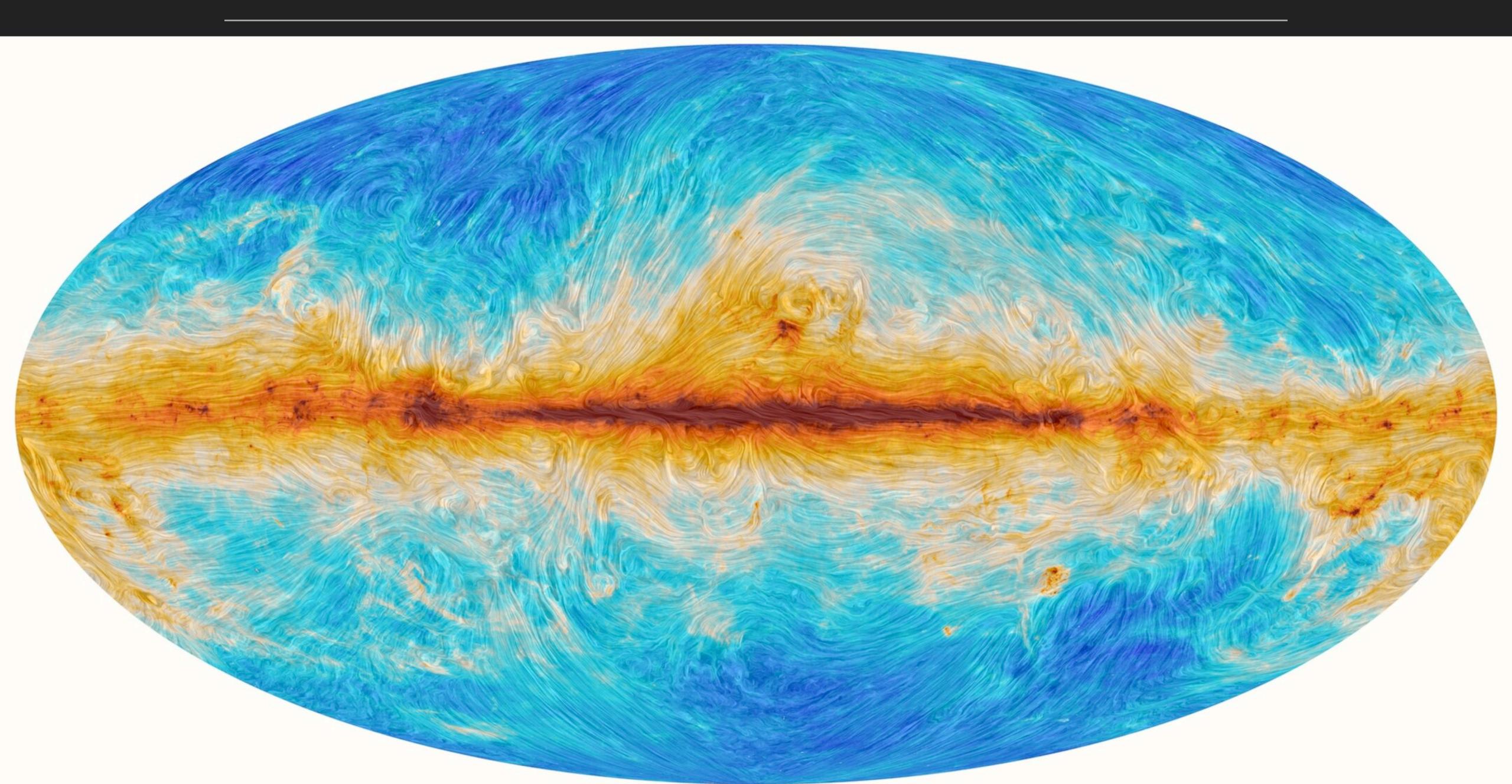
$$f(v_k^{x,y,z}) = w \mathcal{N}(v_k, \sigma = 160 \text{ km/s}) + (1 - w) \mathcal{N}(v_k, \sigma = 780 \text{ km/s})$$
 (1)

with w = 0.90.

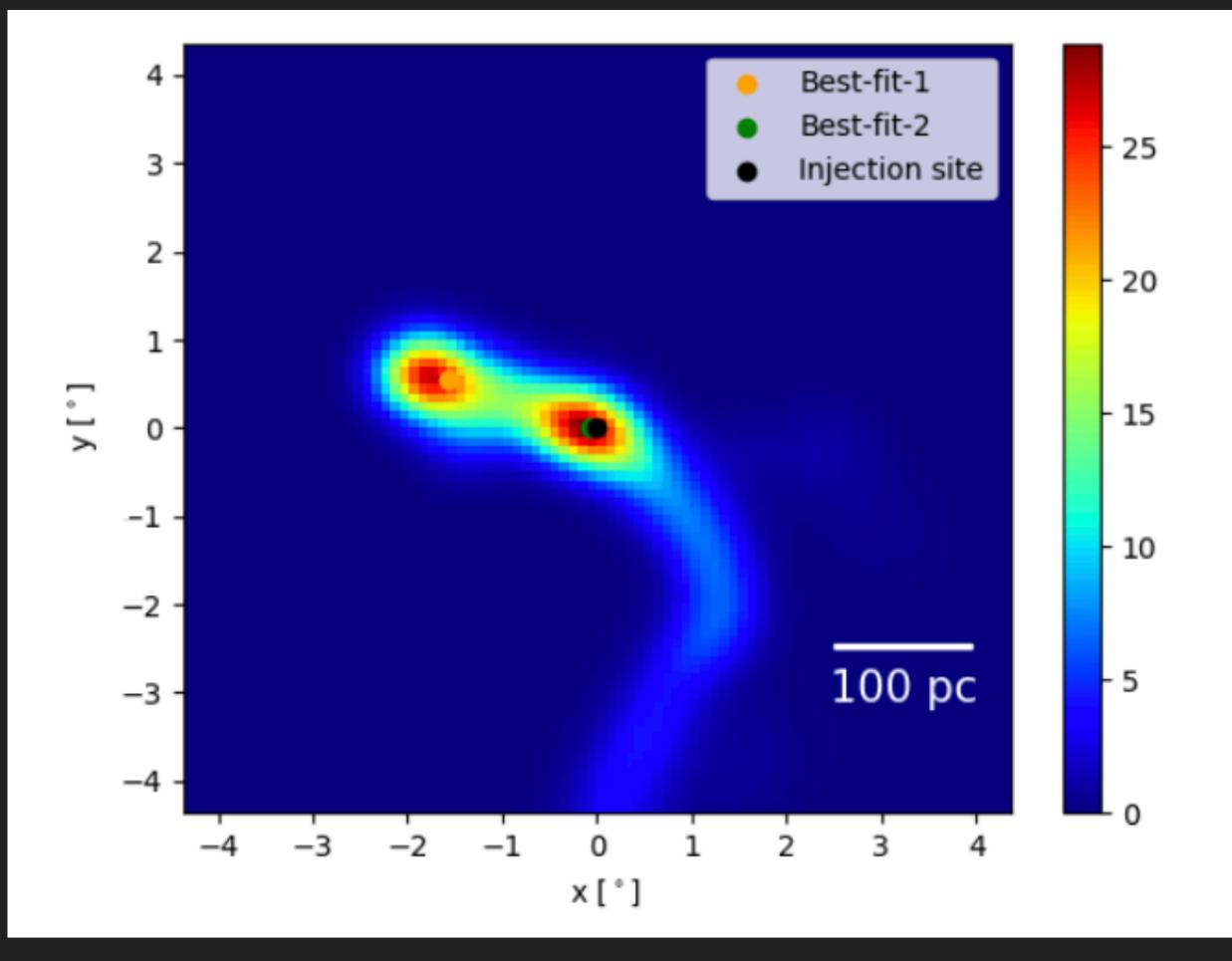




PUZZLE IIIB: MIRAGE HALOS

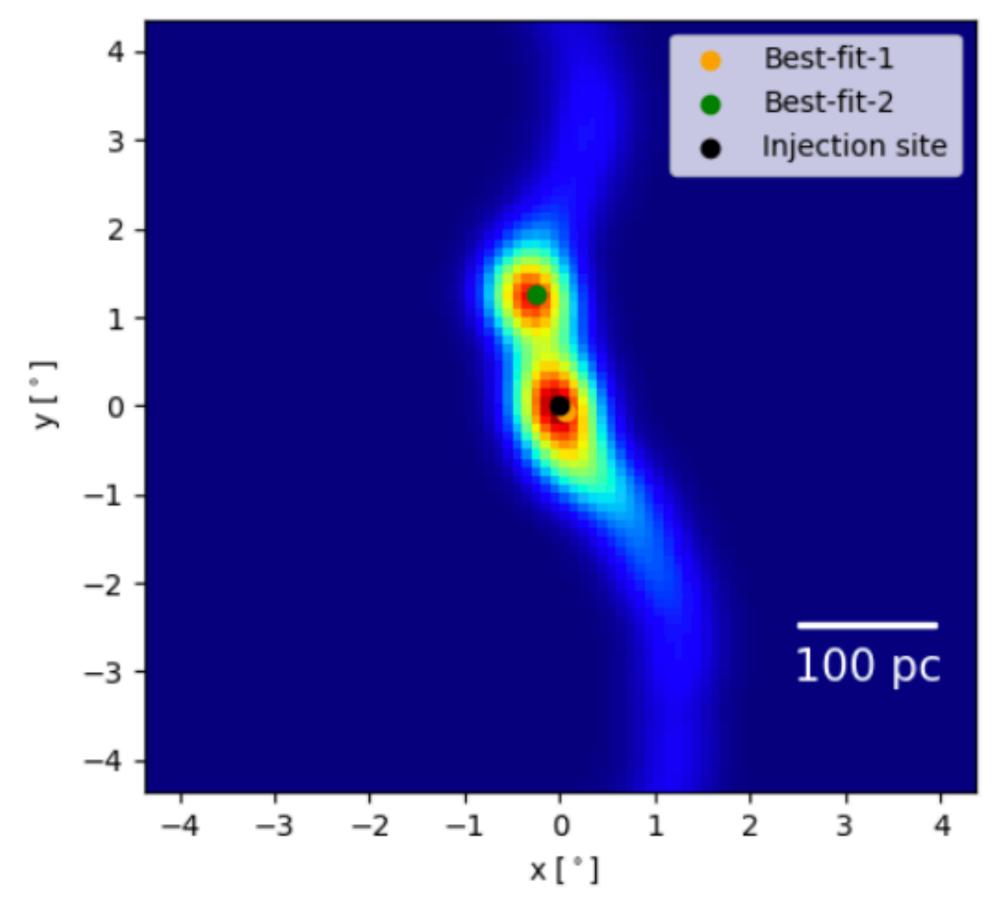


PUZZLE IIIB: MIRAGE HALOS

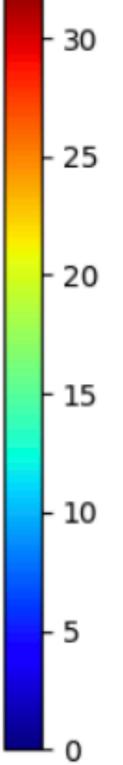


 "Mirage" TeV halos - Many more systems may be difficult to detect or analyze, because they break the modeling assumption of spherical symmetry.

Bao et al. (2407.02478)







PUZZLE IIIC: TEV HALOS POWERED BY MILLISECOND PULSARS?

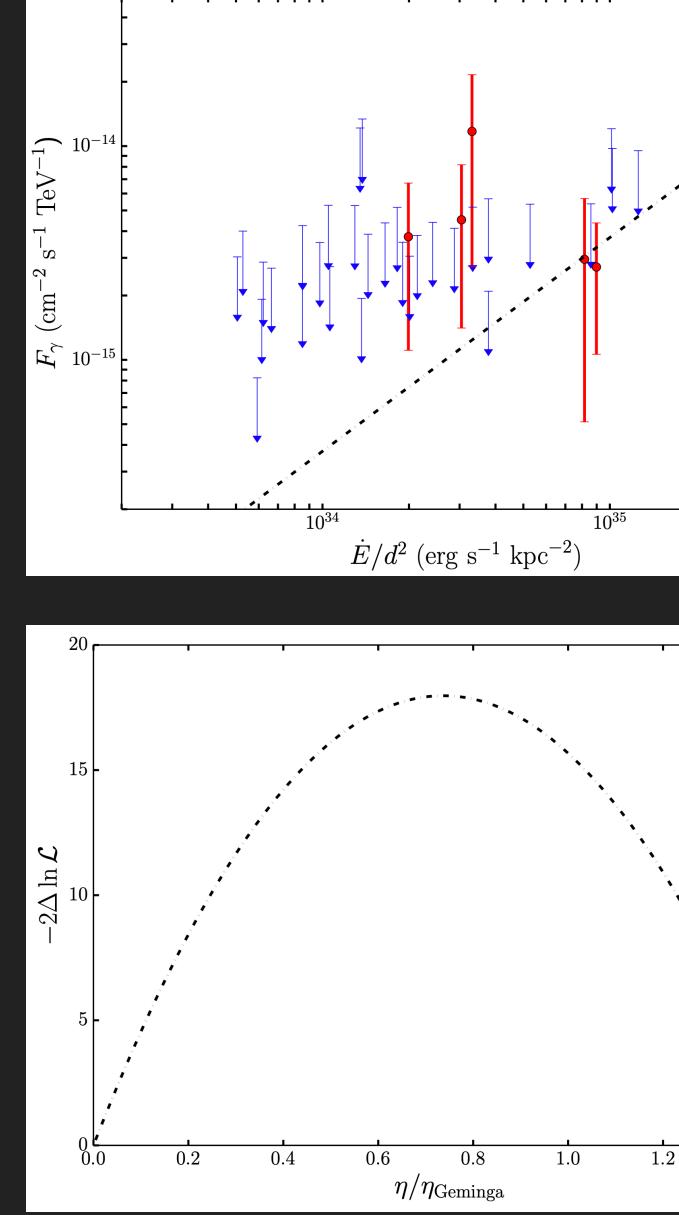
Do MSPs Have TeV Halos?

- Tentative: 4.24σ Poisson evidence from a HAWC stacking analysis (~ 2.3σ from blank sky test).
- Possible MSP Detection by LHAASO
- Important theoretical implications:
 - Cosmic-Ray confinement near pulsars?
 - Cosmic-Ray diffusion at high latitudes
 - PWN/Magnetospheric acceleration models.

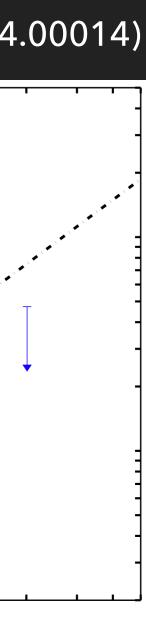
LHAASO Collaboration (2023; 2305.17030)

1LHAASO J0216+4237u 0.33 ATNF PSR J0218+4232 4FGL J0218.1+4232 0.33

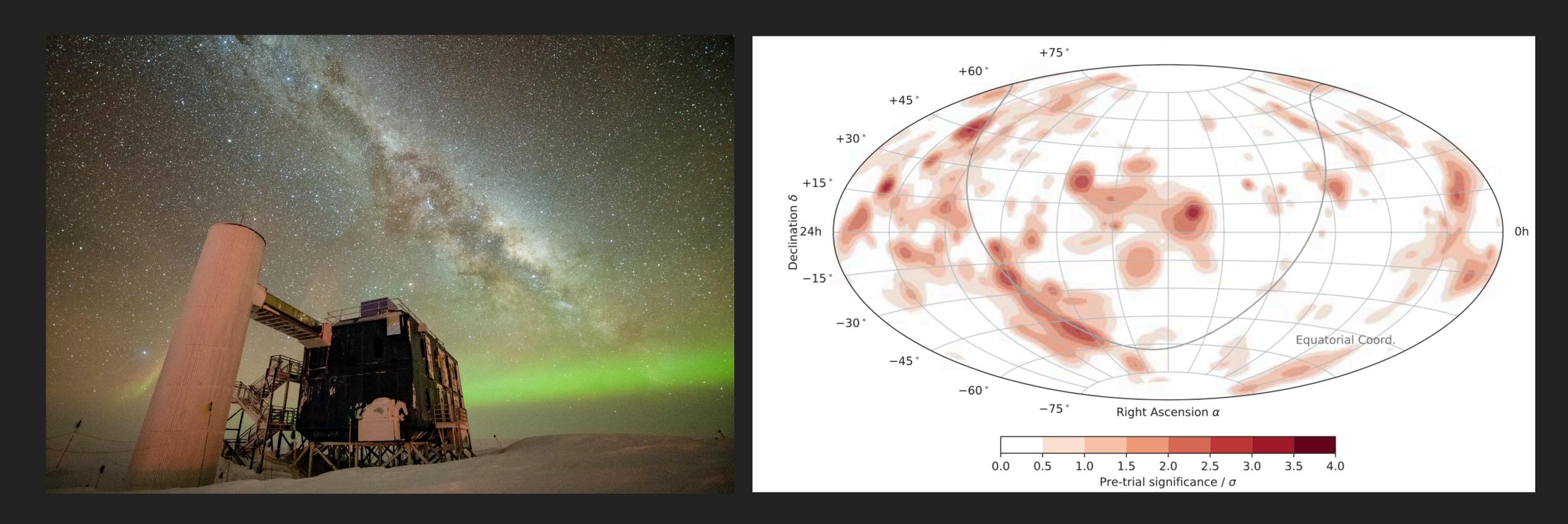
Hooper, TL (2021; 2104.00014)



 $\dot{E} = 2.44 \times 10^{35} \text{ erg s}^{-1}, \tau_c = 476000.0 \text{ kyr}, d = 3.15 \text{ kpc}$ PSR J0218+4232;MSP;

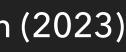






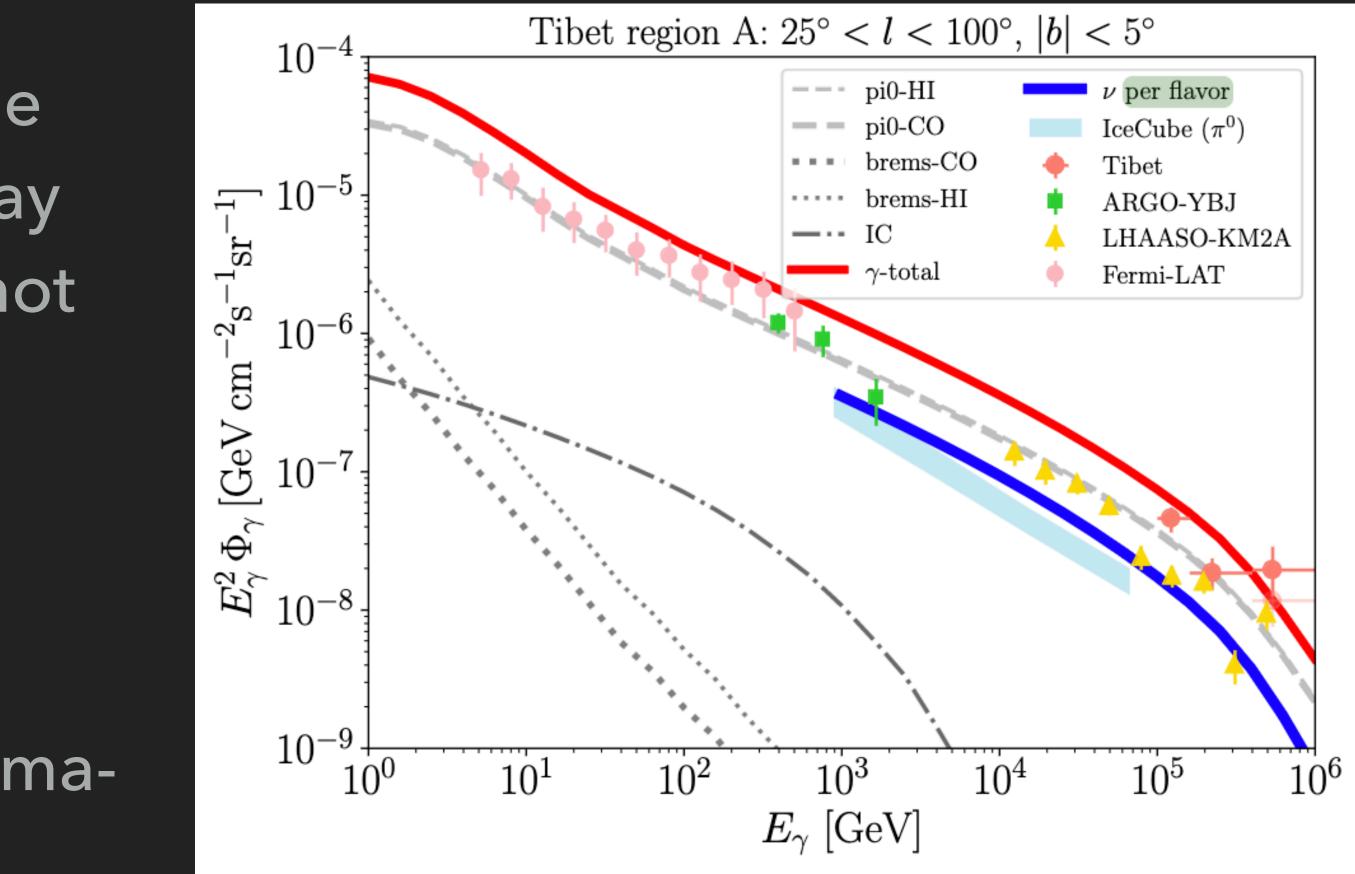
IceCube detection of a galactic neutrino flux – with a normalization that is ~4x brighter than expectations from the Fermi-LAT extrapolation.

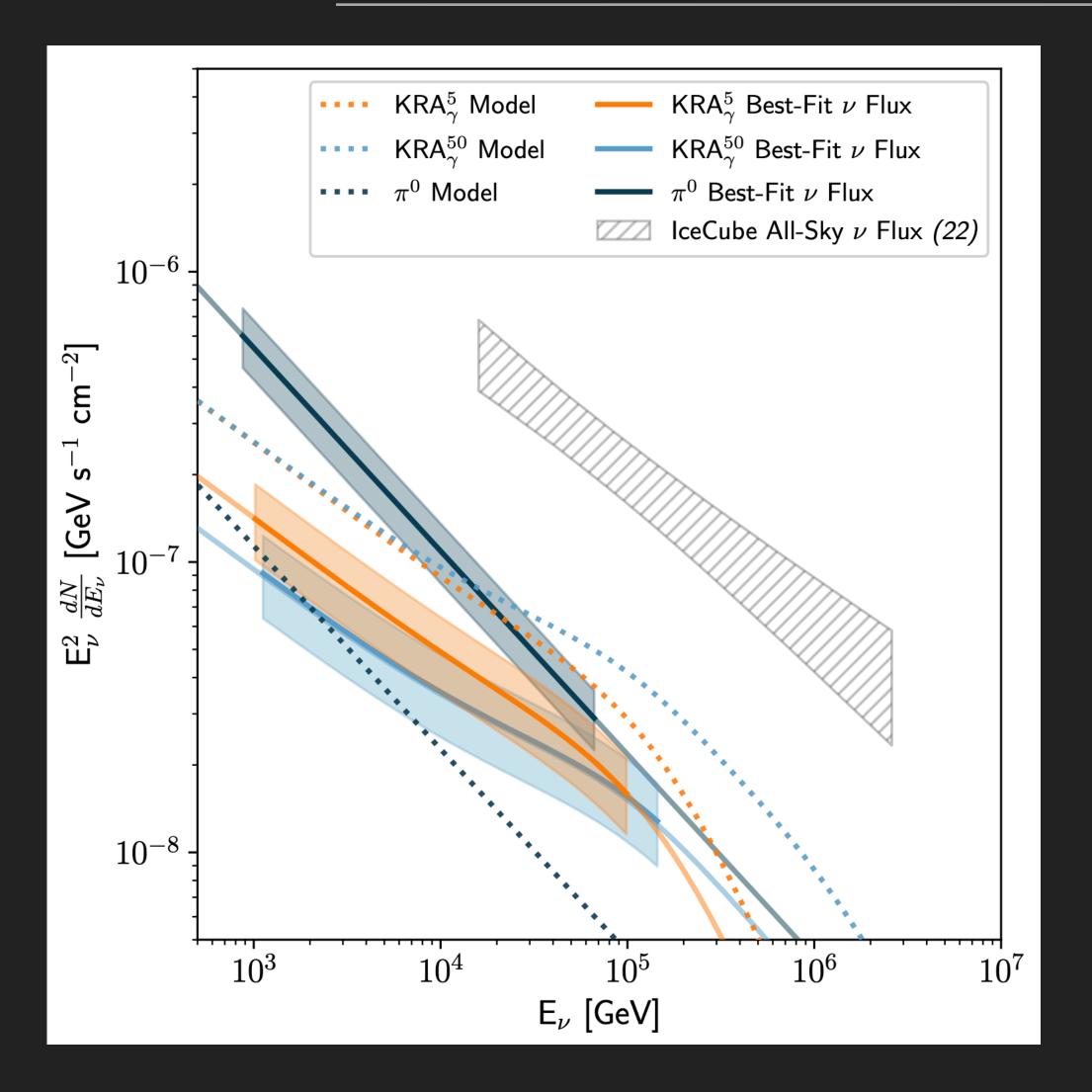
IceCube Collaboration (2023)



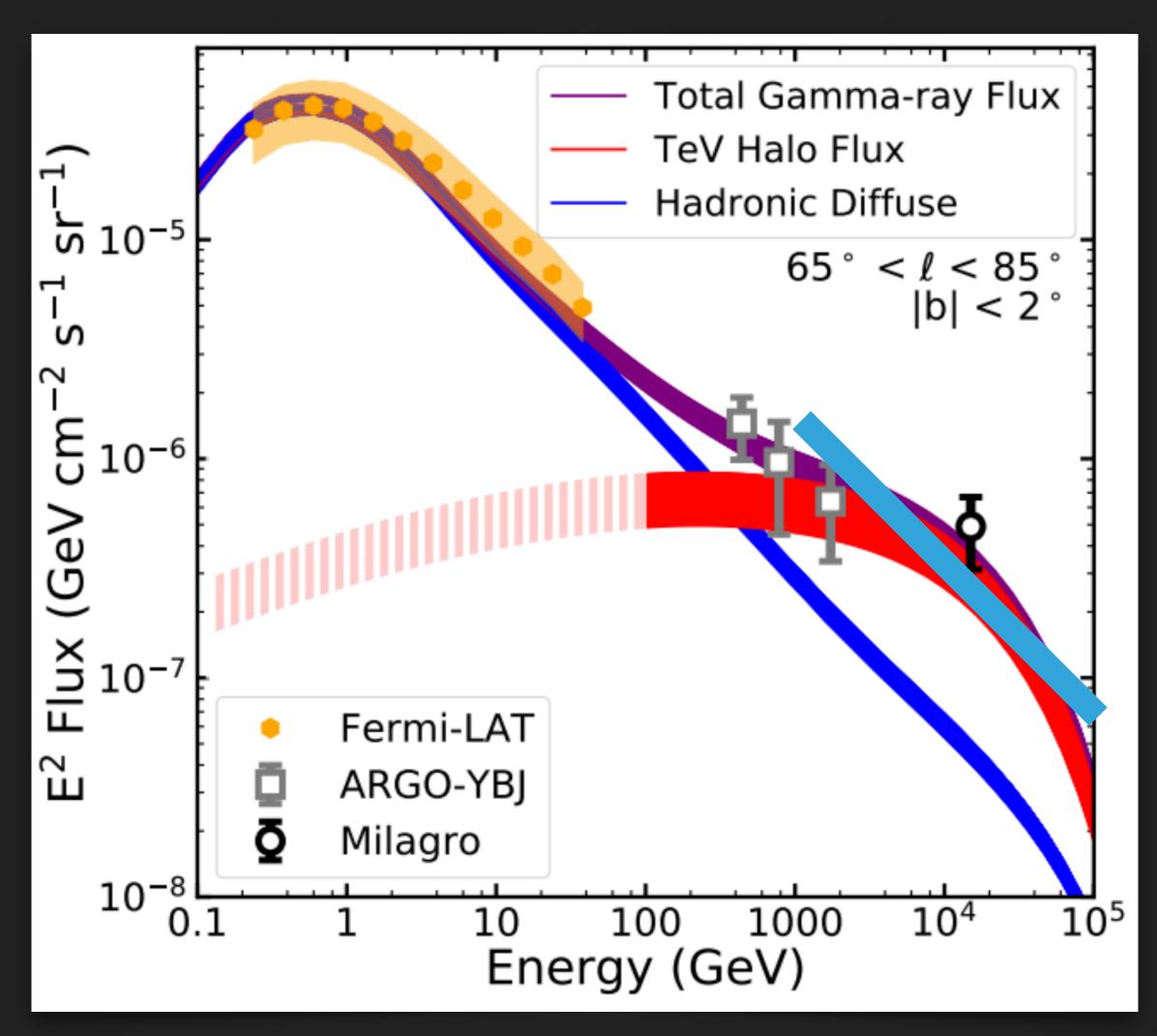
 If the IceCube neutrino flux from the galaxy is higher, then the gamma-ray flux from hadronic processes (i.e., not halos) could also be higher.

 In Fang et al. this is capable of producing the diffuse galactic gammaray emission Fang et al. (2023; 2306.17275)

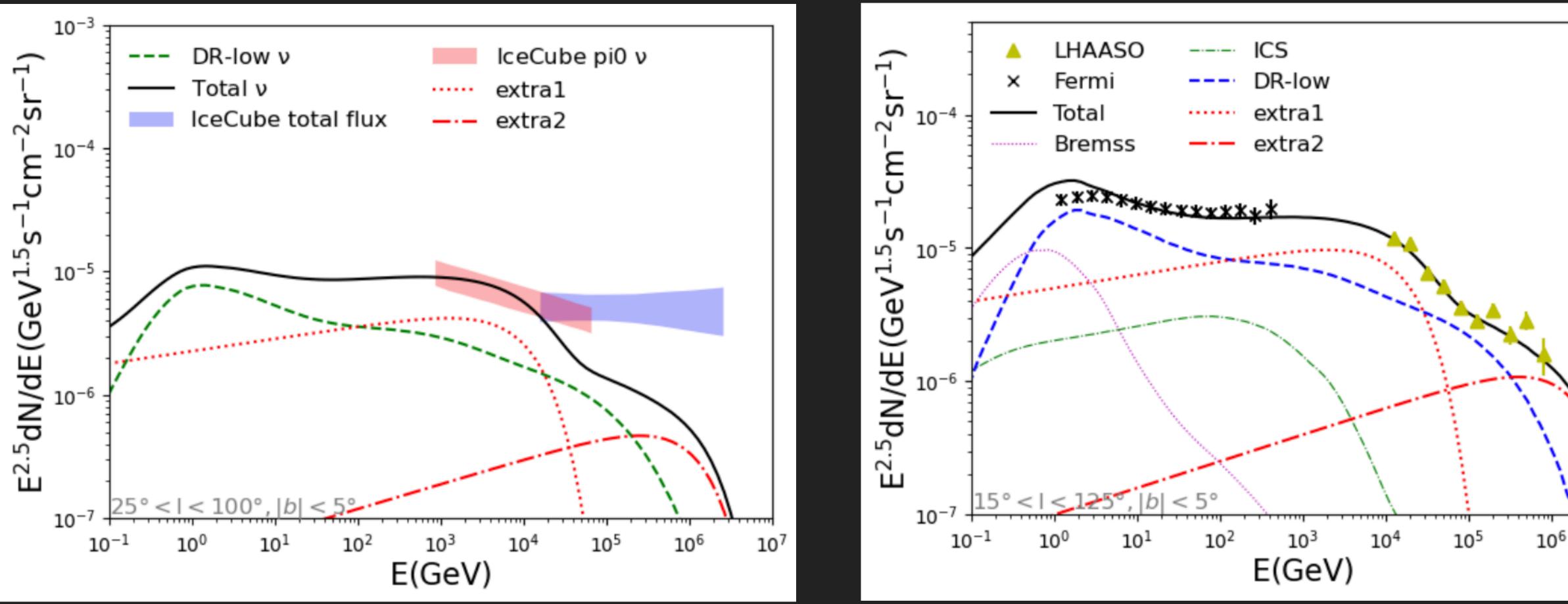




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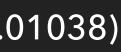
IceCube detection of a galactic neutrino flux – with a normalization that



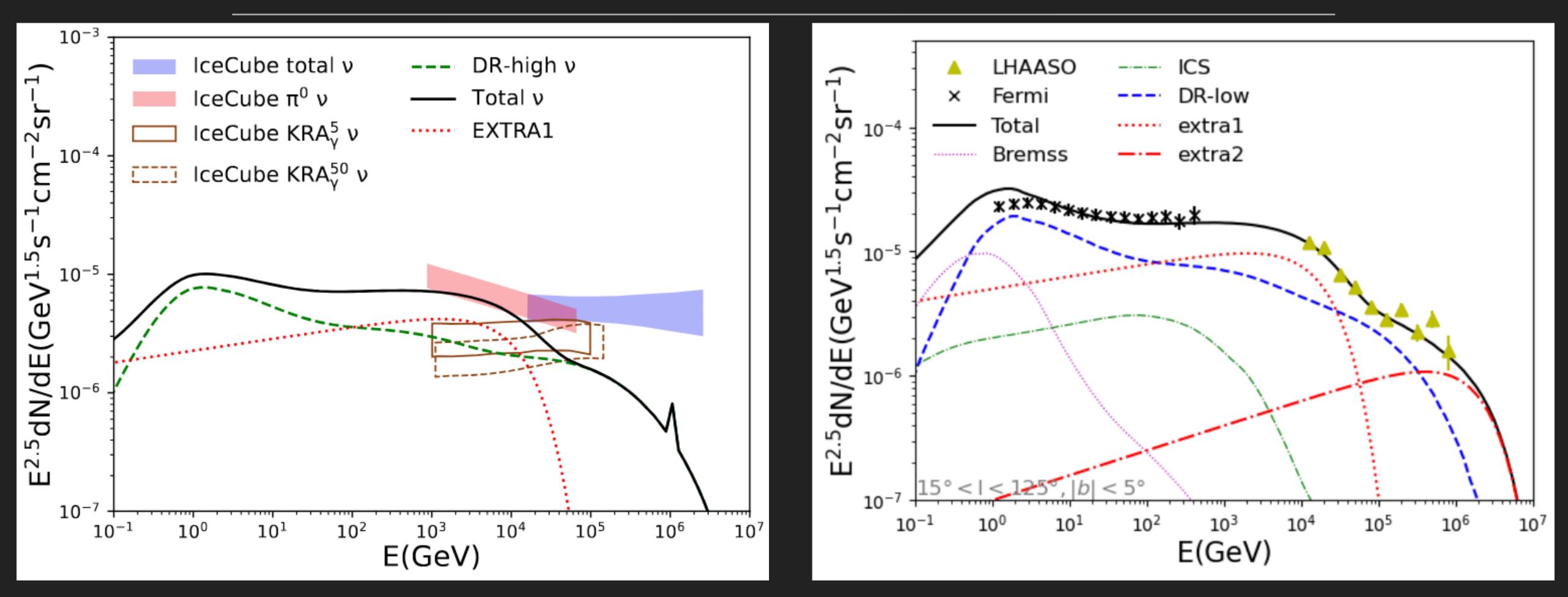
Models that explain the IceCube neutrino flux still require an additional gamma-ray component (here: "Extra1 and Extra2").

In this model it is hadronic.

Shao et al. (2023; 2307.01038)



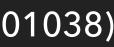




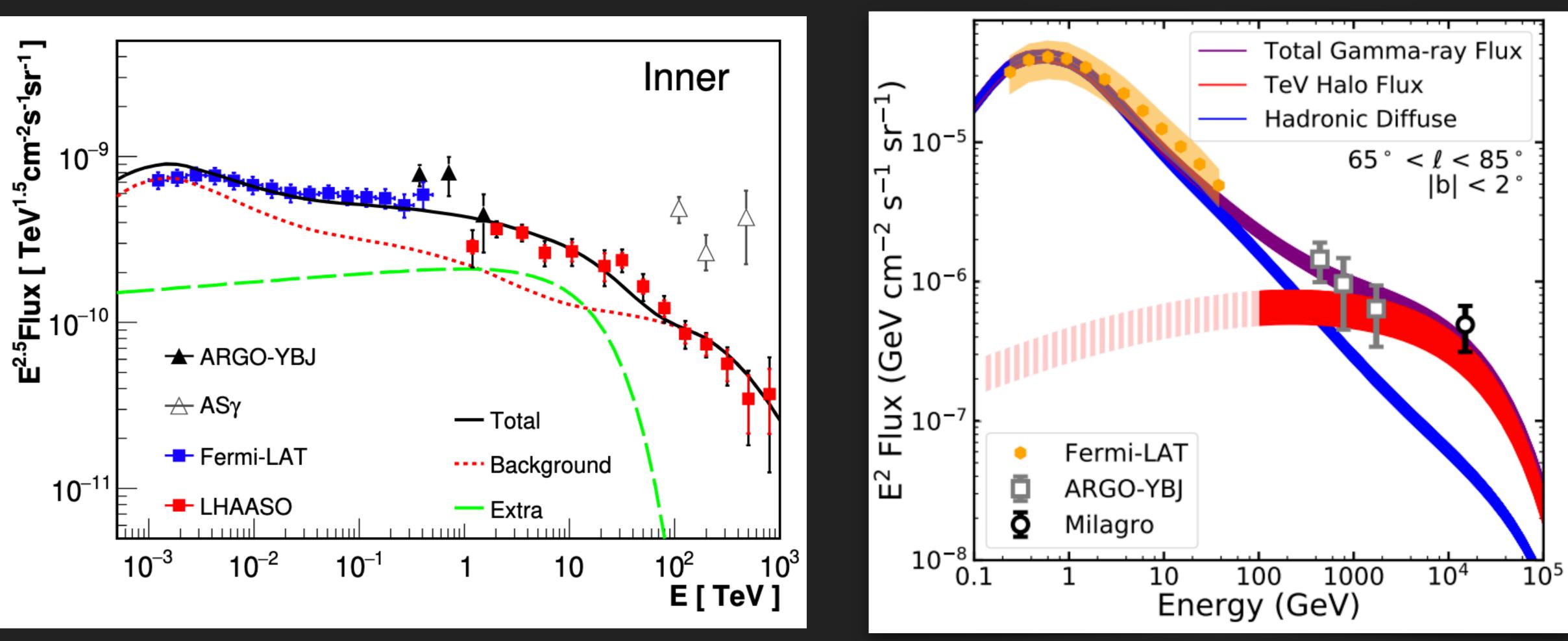
Models that explain the IceCube neutrino flux still require an additional gamma-ray component (here: "Extra1 and Extra2").

In this model it is likely leptonic.

Shao et al. (2023; 2307.01038)



PUZZLE IV: HADRONIC VS. LEPTONIC DIFFUSE TEV EMISSION LHAASO Collaboration (2411.16021)



EXTRA component as well, with a very similar spectrum to Geminga.

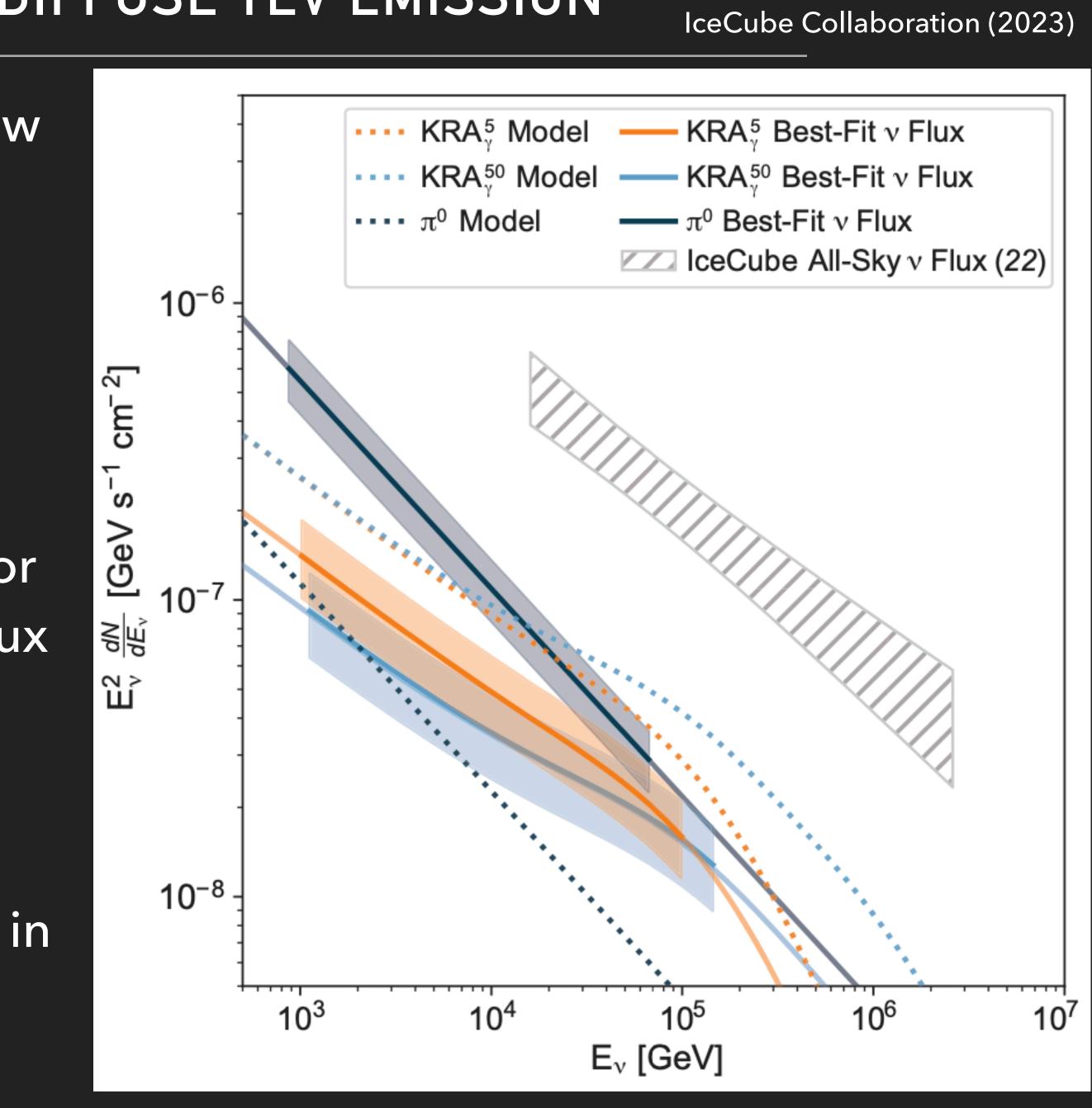
• LHAASO collaboration recently reported a diffuse spectrum requiring an



IceCube neutrino flux is unknown at low energies (nearly order of magnitude uncertainties from models that fit the data to within 1σ .

On top of this, there is an intrinsic factor of 2 uncertainty in even the IceCube flux measurement.

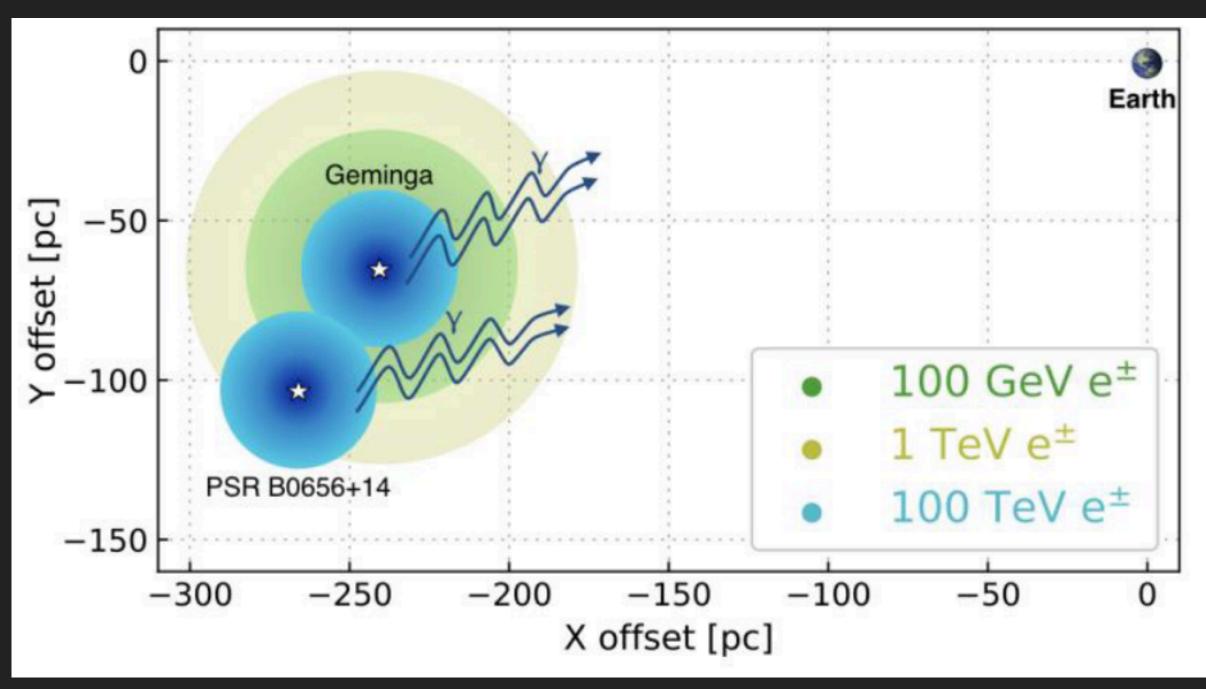
There is also a factor of ~2 uncertainty in the TeV halo flux owing to the "Geminga-like" assumption

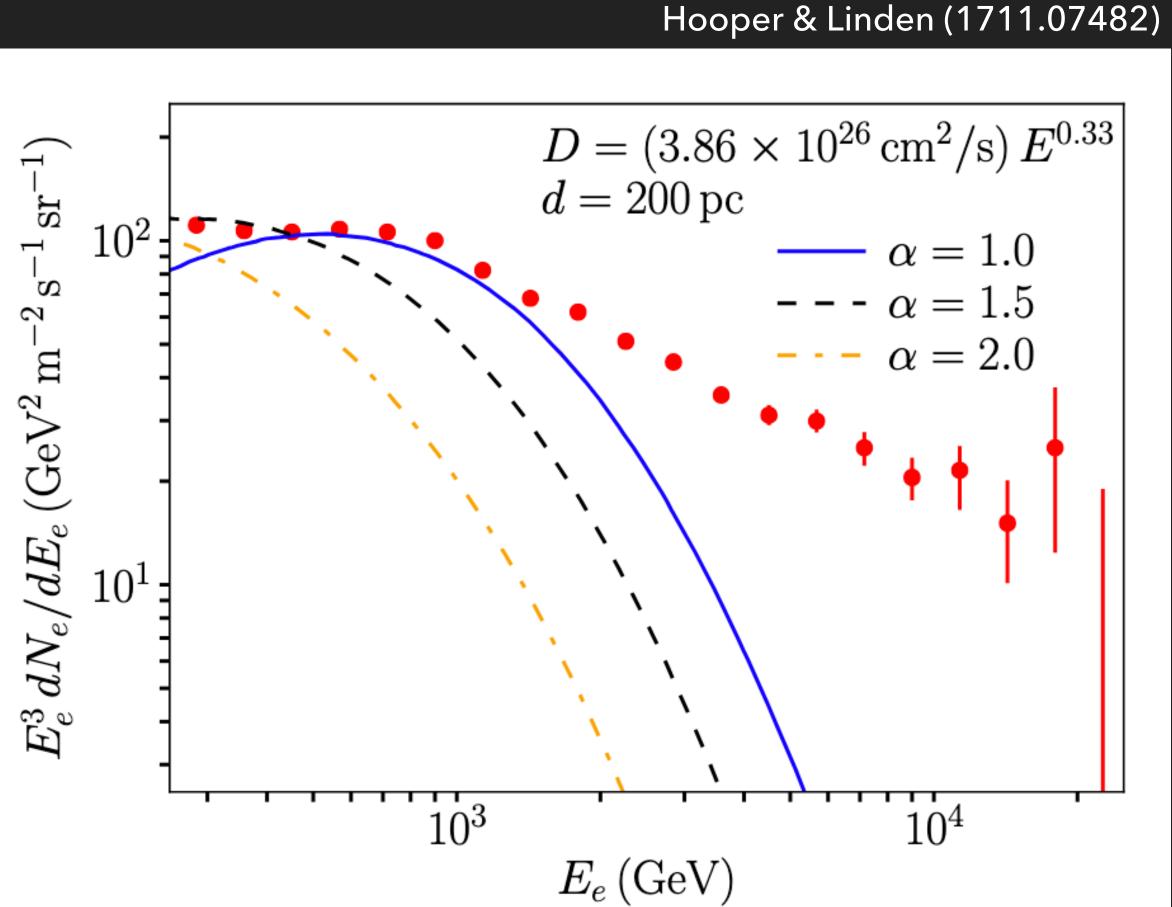


PUZZLE V: WHAT ARE TEV HALOS?

Failed Model 1: One zone models

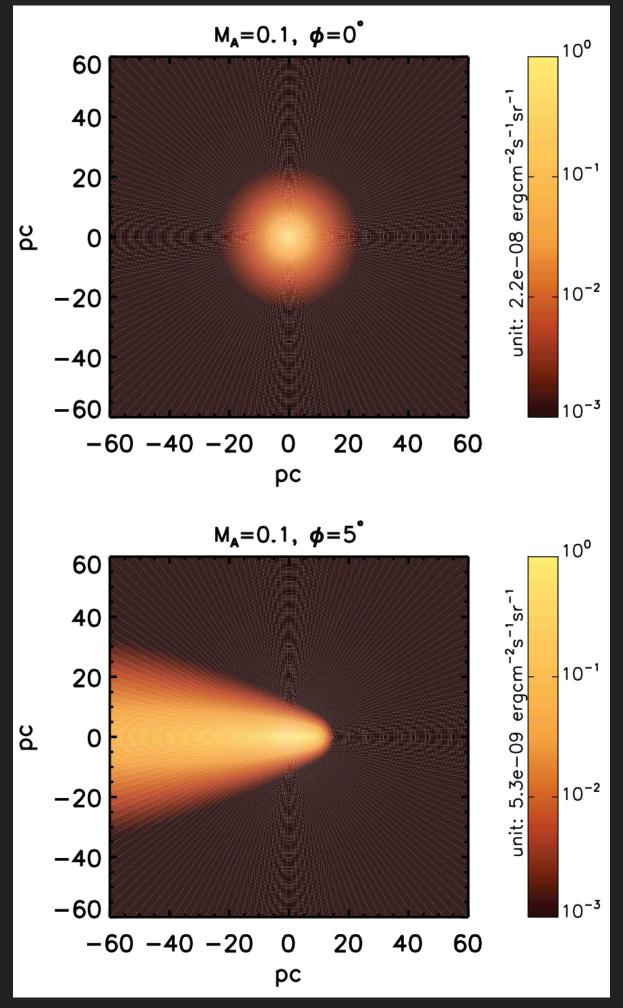
HAWC Collaboration (1711.06223)

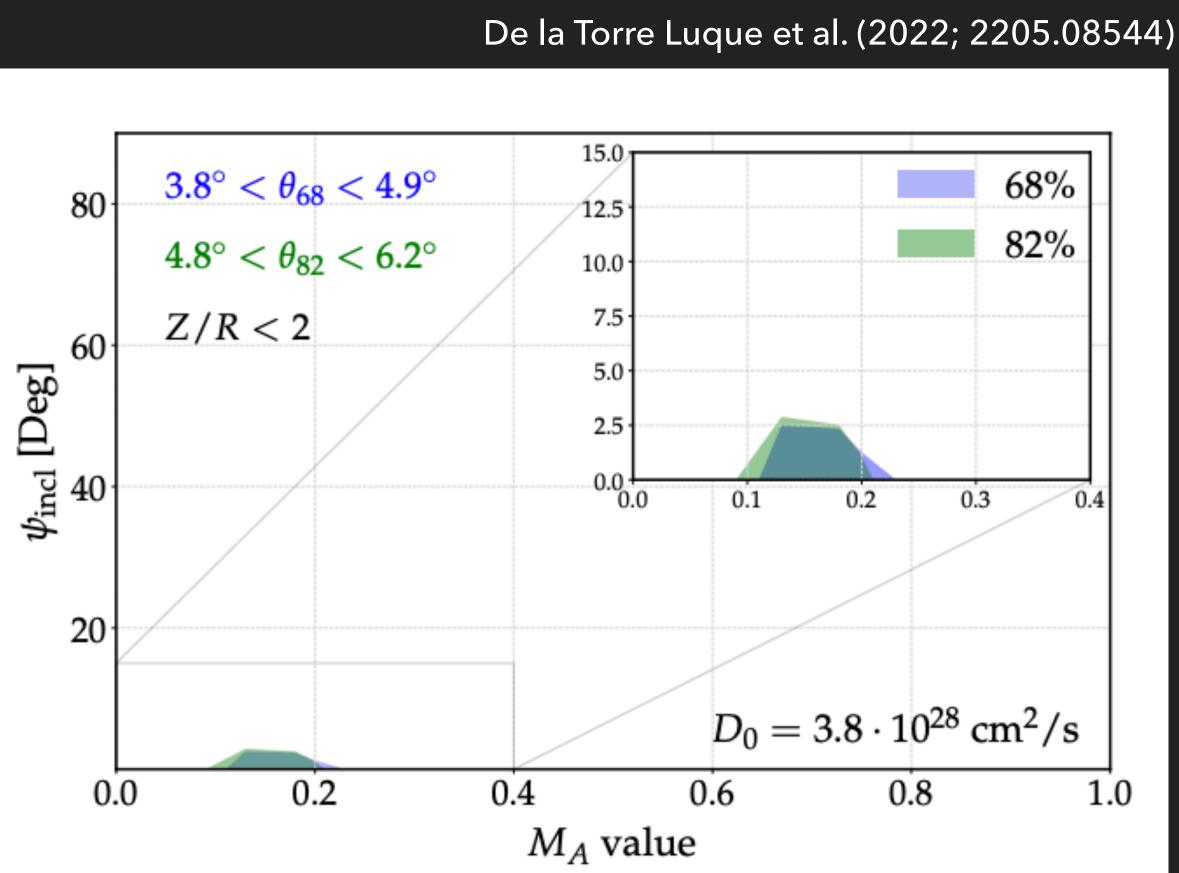




Failed Model 2: Magnetic fields anomalously pointed towards the Earth

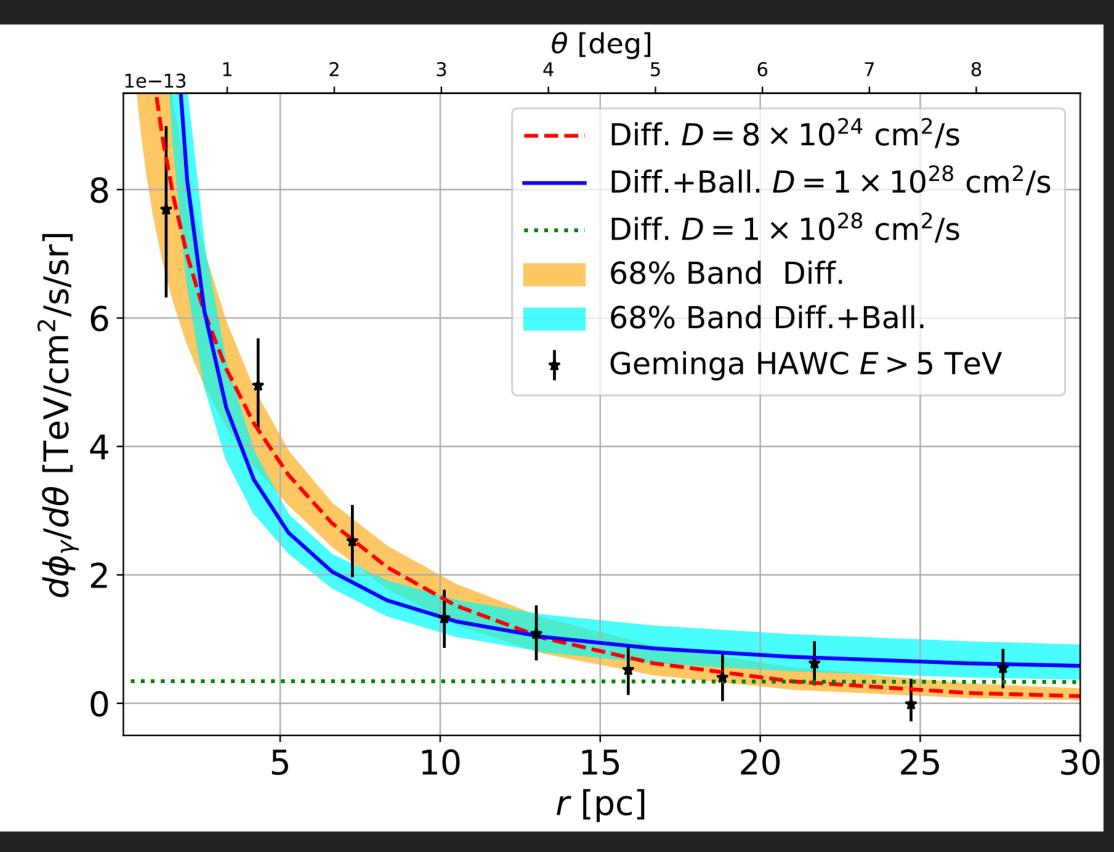
Liu, Yan, Zhang (2019; 1904.11536)



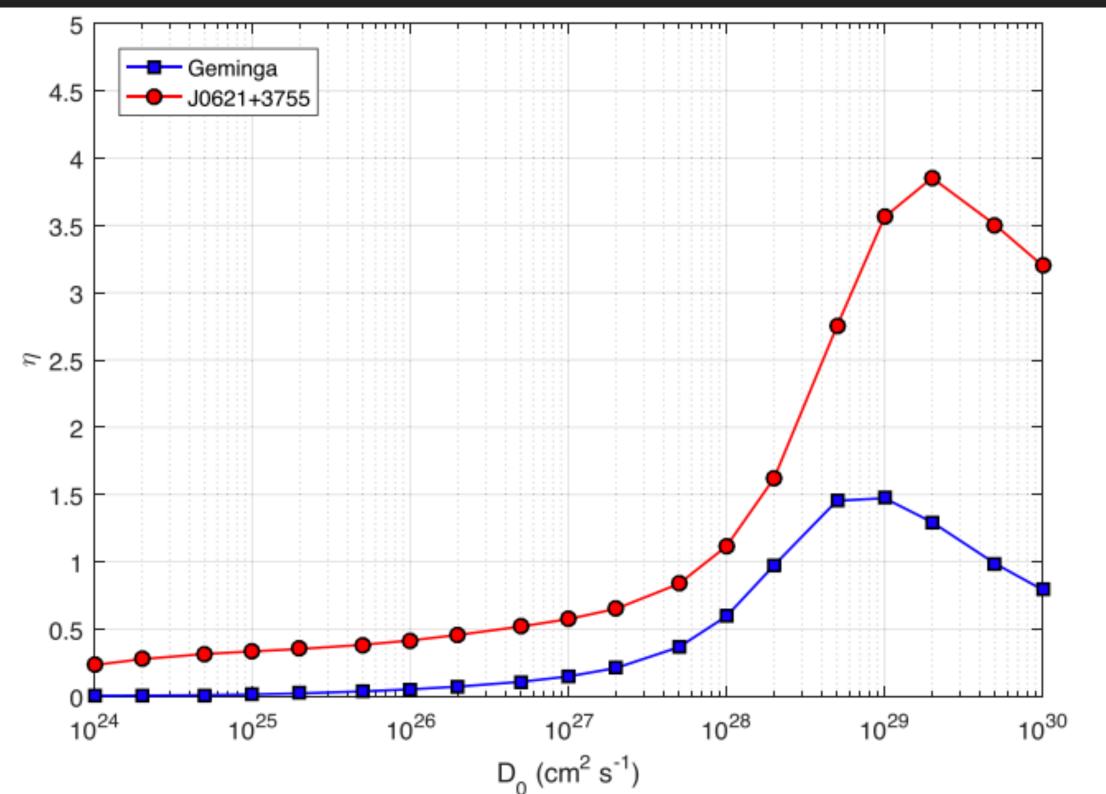


Failed Model 3: Rectilinear propagation during gamma-ray production.

Recchia et al. (2021; 2106.02275)



Bao et al. (2021; 2107.07395)



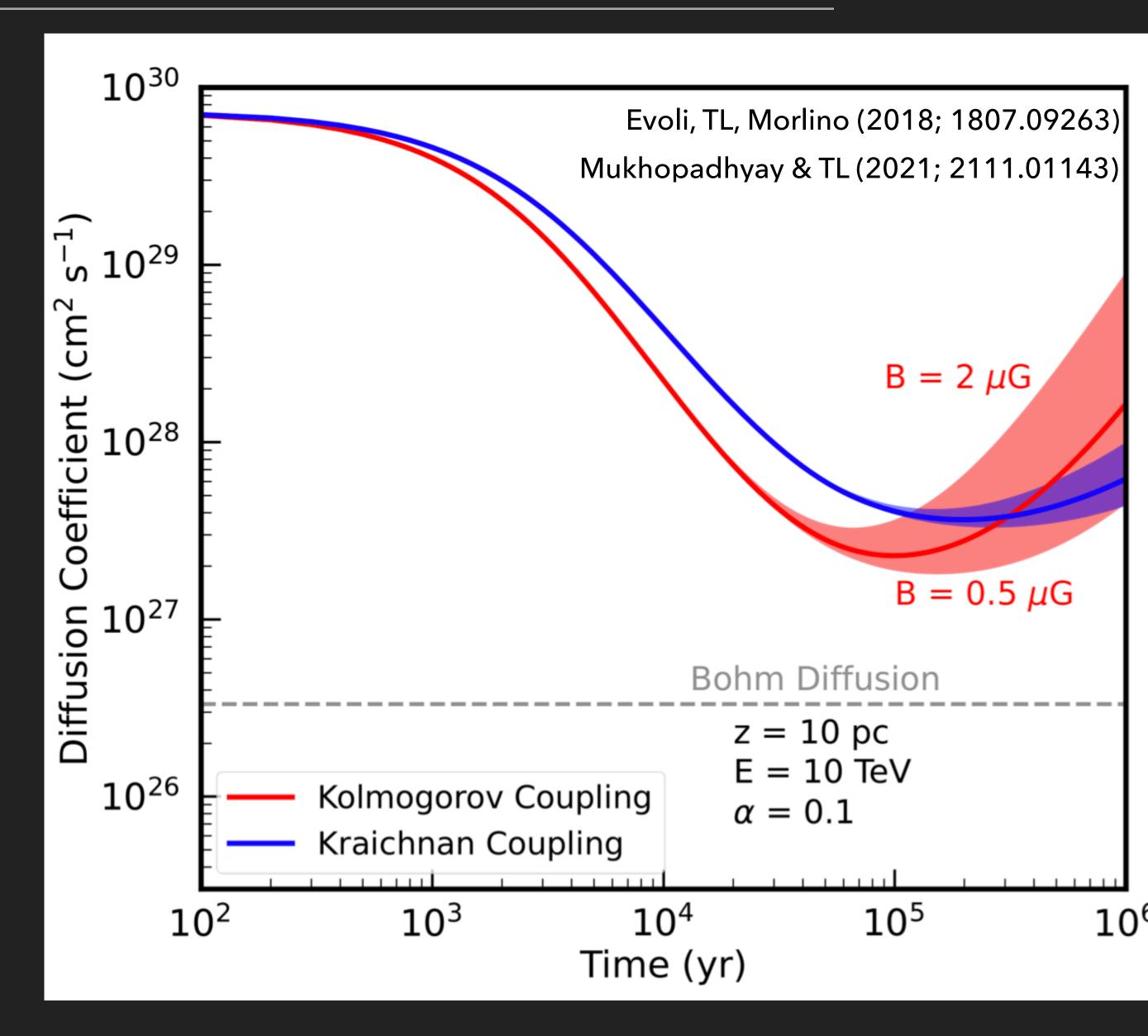


 Self-confinement models (and most other models for inhibited diffusion) - require the high energy of a very young pulsar.

 Probing the diffusion around the youngest systems is critical for understanding TeV halo dynamics.

$$\frac{\partial \mathcal{W}}{\partial t} + v_A \frac{\partial \mathcal{W}}{\partial z} = (\Gamma_{\rm CR} - \Gamma_{\rm D}) \mathcal{W}(k, z, t)$$

$$\Gamma_{
m CR}(k) = rac{2\pi}{3} rac{c|v_A|}{k\mathcal{W}(k) U_0} \left[p^4 rac{\partial f}{\partial z}
ight]_{p_{
m res}}$$



$$\frac{\partial W}{\partial t} + v_A \frac{\partial W}{\partial z} = (\Gamma_C t)$$

$$\Gamma_{CR}(k) = \frac{2\pi}{3} \frac{c |v_{\alpha}|}{k W(k)} \left(\frac{B_0^2}{8\pi}\right)^{-1} \left[p^4 \frac{\partial f}{\partial z}\right]_{p_{\text{res}}}$$

$$D(p,t) = \frac{4}{3\pi} \frac{cr_L(p)}{k_{\rm res}W(z,k_{\rm res}))}$$

 $_{CR} + \Gamma_{NLD}) W(k, z, t)$

$$\Gamma_{NLD}(k) = c_k v_{\alpha} \begin{cases} k^{3/2} W^{1/2} & \text{Kolmogor}\\ k^2 W & \text{Kraichna} \end{cases}$$



Many uncertainties in these models:

- Role of Supernova Remnant
- Disruption by molecular gas or magnetic fields
- Pulsar Proper Motion
- ▶ 1D vs. 3D diffusion
- non-Resonant Terms
- Halos in close proximity

Possible origin of the slow-diffusion region around Geminga

Kun Fang¹* Xiao-Jun Bi^{1,2}† Peng-Fei Yin¹‡

¹ Key Laboratory of Particle Astrophysics, Institute of High Energy Physics, Chinese Academy of Sciences, Beijing 100049, China ² School of Physical Sciences, University of Chinese Academy of Sciences, Beijing 100049, China

23 July 2019

ABSTRACT

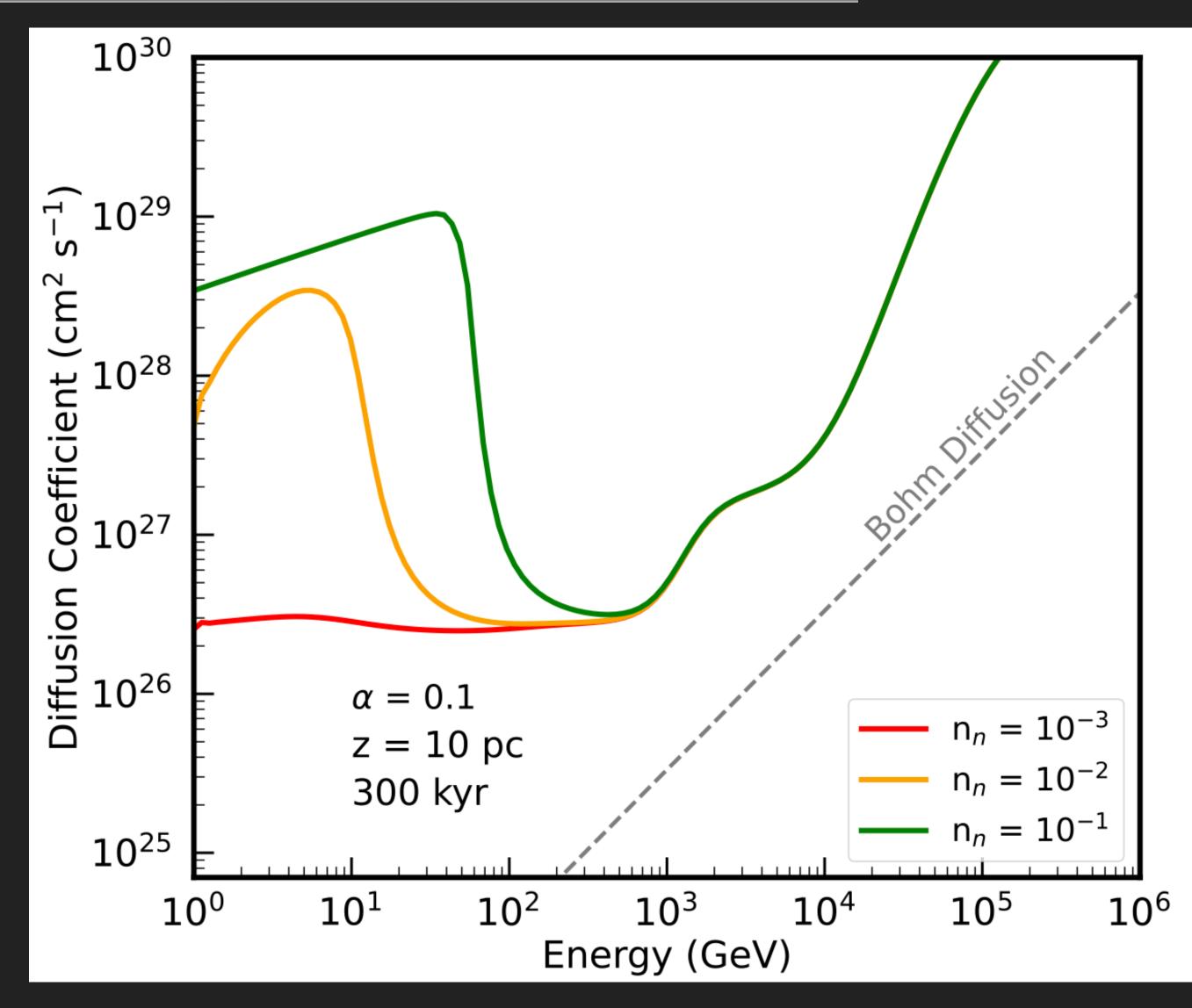
Geminga pulsar is surrounded by a multi-TeV γ -ray halo radiated by the high energy electrons and positrons accelerated by the central pulsar wind nebula (PWN). The angular profile of the γ -ray emission reported by HAWC indicates an anomalously slow diffusion for the cosmic-ray electrons and positrons in the halo region around Geminga. In the paper we study the possible mechanism for the origin of the slow diffusion. At first, we consider the self-generated Alfvén waves due to the streaming instability of the electrons and positrons released by Geminga. However, even considering a very optimistic scenario for the wave growth, we find this mechanism DOES NOT work to account for the extremely slow diffusion at the present day if taking the proper motion of Geminga pulsar into account. The reason is straightforward as the PWN is too weak to generate enough high energy electrons and positrons to stimulate strong turbulence at the late time. We then propose an assumption that the strong turbulence is generated by the shock wave of the parent supernova remnant (SNR) of Geminga. Geminga may still be inside the SNR, and we find that the SNR can provide enough energy to generate the slow-diffusion circumstance. The TeV halos around PSR B0656+14, Vela X, and PSR J1826-1334 may also be explained under this assumption.

Key words: cosmic rays – ISM: individual objects: Geminga nebula – ISM: supernova remnants – turbulence

Many uncertainties in these models:

- Role of Supernova Remnant
- Disruption by molecular gas or magnetic fields
- Pulsar Proper Motion
- ▶ 1D vs. 3D diffusion
- non-Resonant Terms
- Halos in close proximity

Evoli, TL, Morlino (2018; 1807.09263) Mukhopadhyay & TL (2021; 2111.01143)

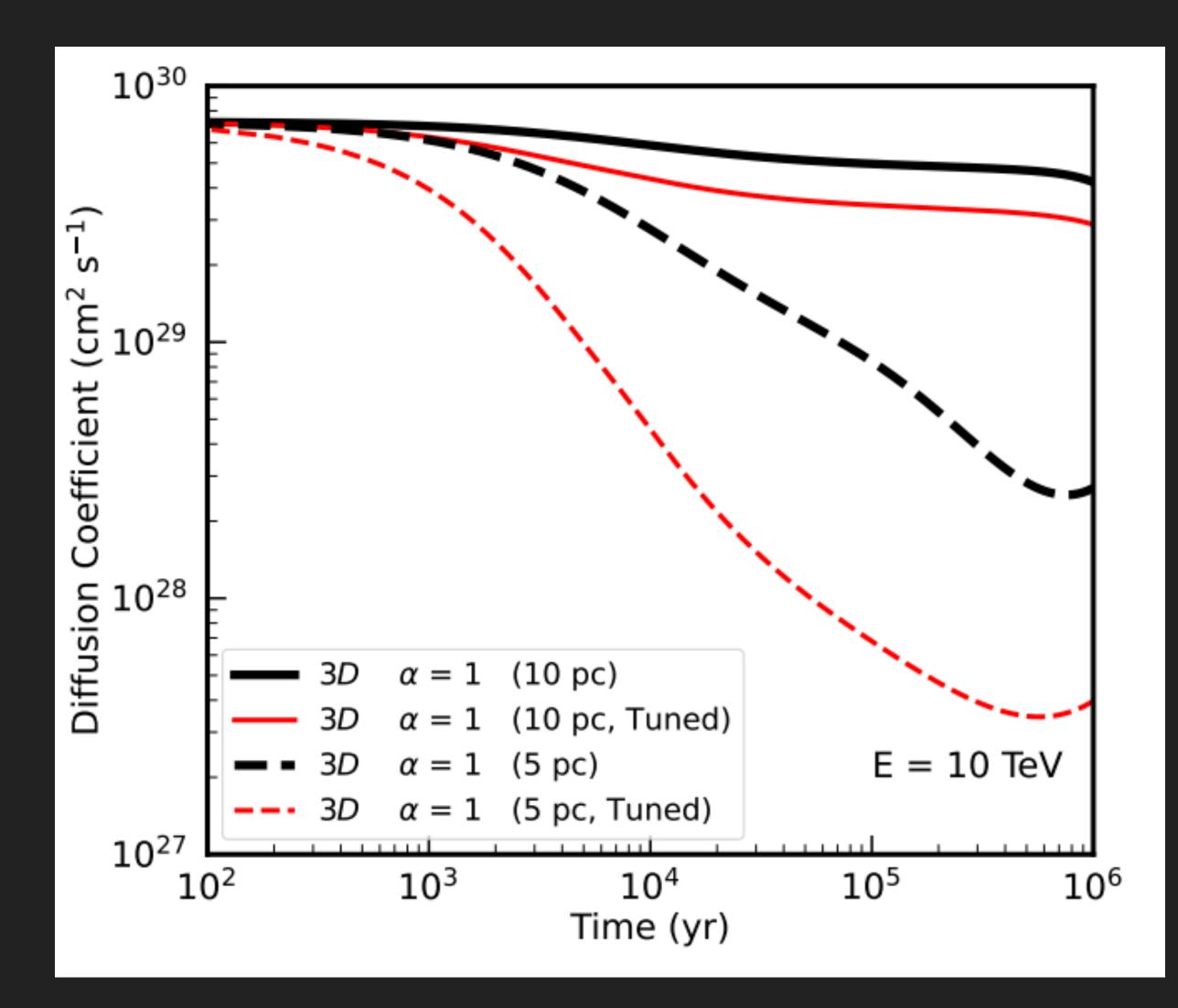




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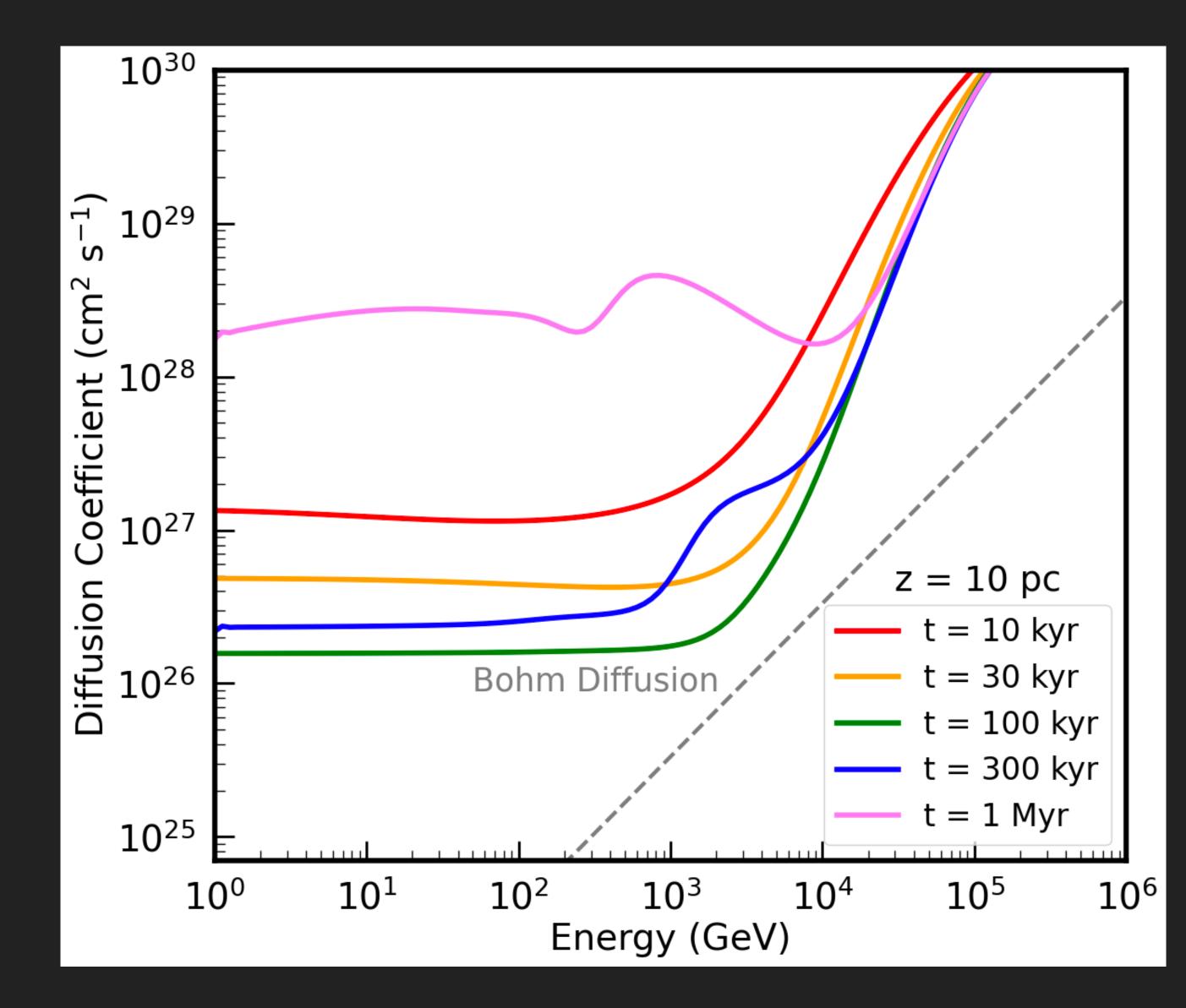
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- Several Predictions of these Models:
 - Relatively flat low-energy diffusion coefficient.

- Highly energy dependent diffusion coefficient at high energies.
 - 100 TeV halo detections challenge this interpretation!

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CONCLUSIONS

recycled pulsars).

emission.

- The next-generation lessons are harder:
 - Understand the diversity of sources.
 - Understand fundamentals of halo diffusion.
 - Understand interplay between leptonic and hadronic sources.

TeV halos are a common feature around middle-aged (and possibly young and

The early lessons were easy – TeV halos prove that pulsars produce the positron flux, and clearly provide a significant fraction of the TeV sources and diffuse TeV

