



The Indirect Detection of Thermal Dark Matter

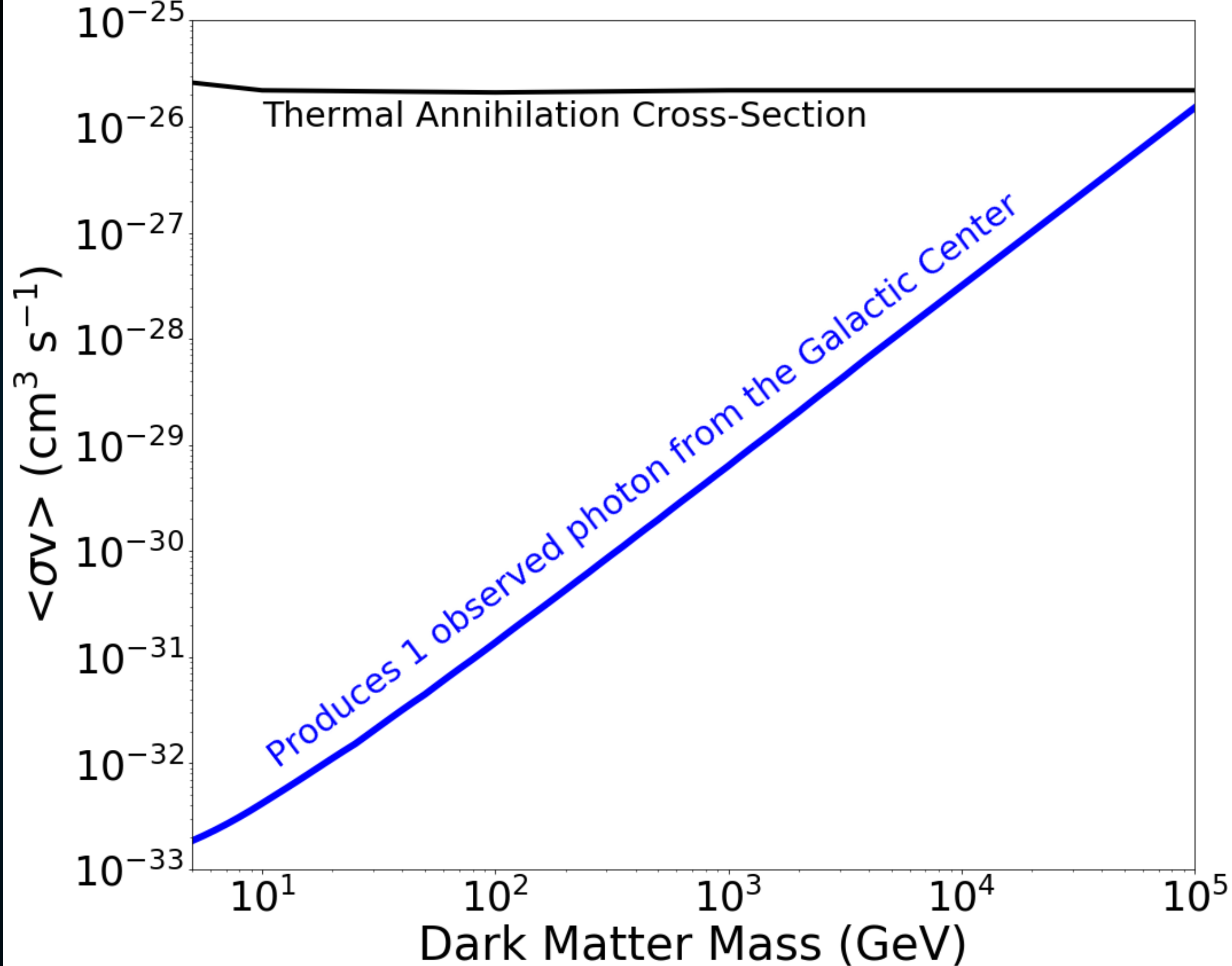
Tim Linden

Case Western Reserve University
Michelson Postdoctoral Prize Lectures
Physics Department Colloquium 10/4/18

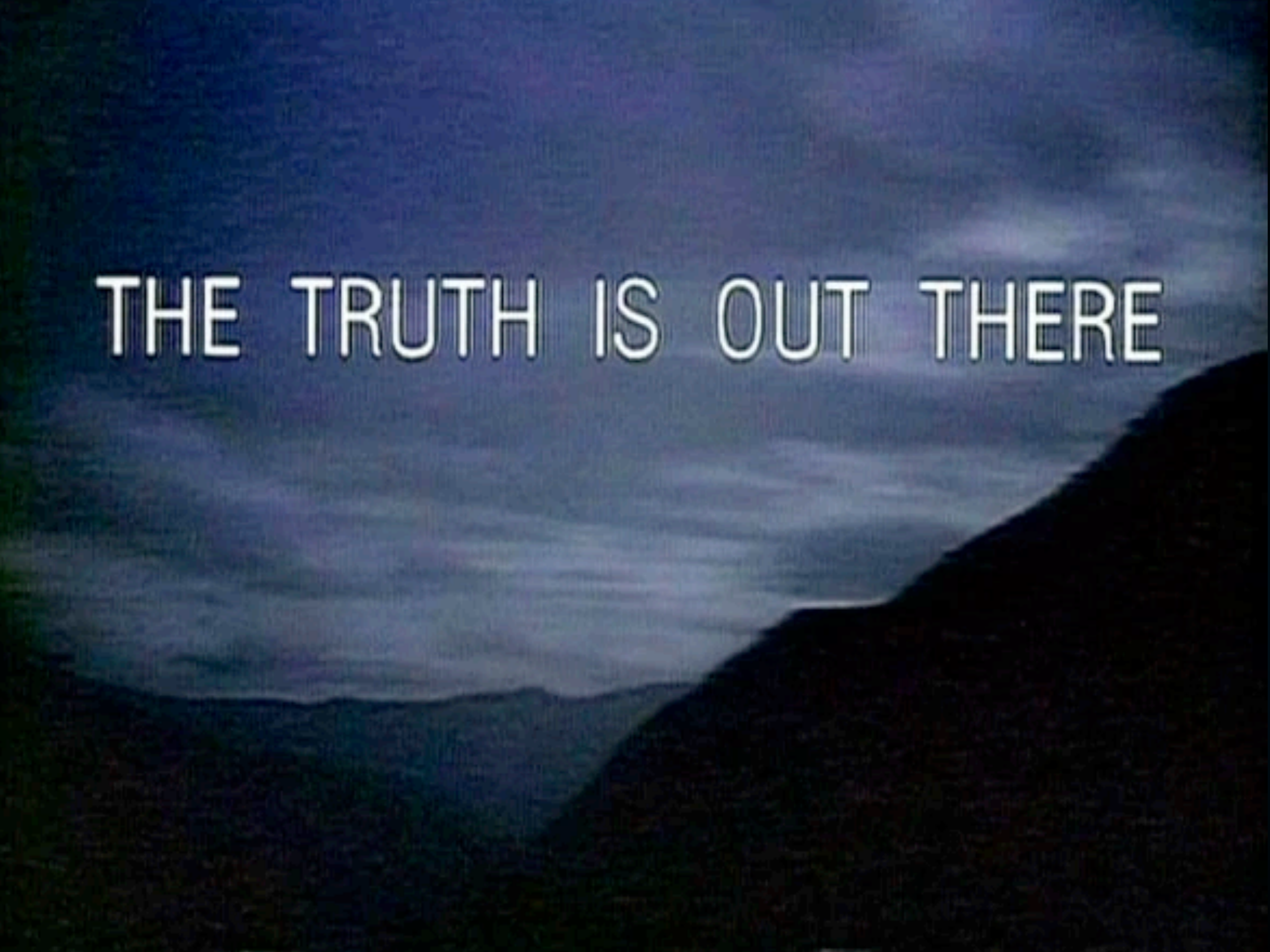


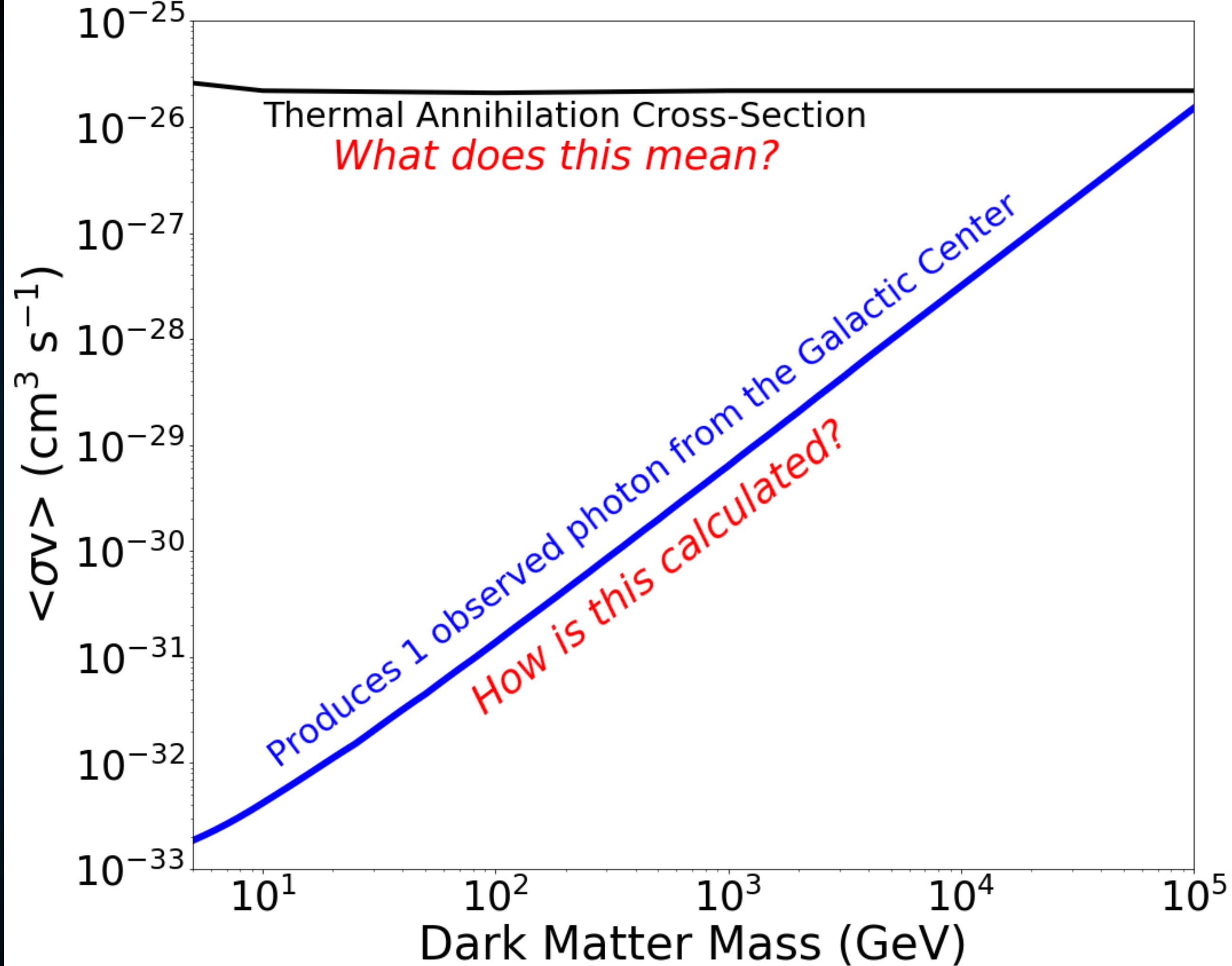
THE OHIO STATE UNIVERSITY

CENTER FOR COSMOLOGY AND
ASTROPARTICLE PHYSICS

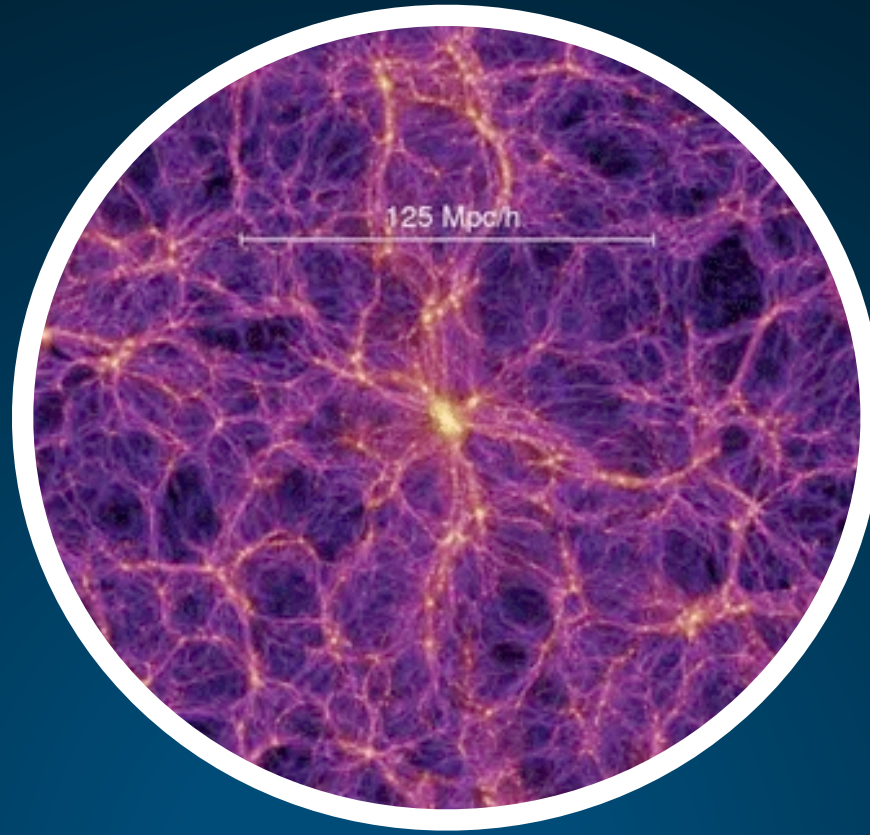
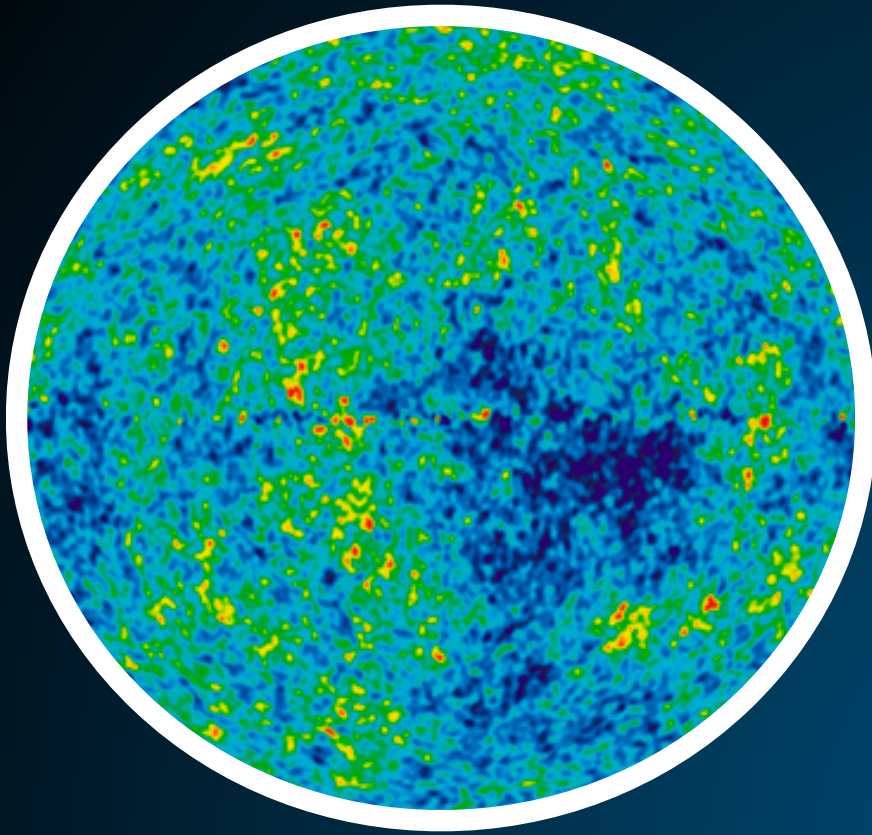


THE TRUTH IS OUT THERE





Everything We Know About Dark Matter



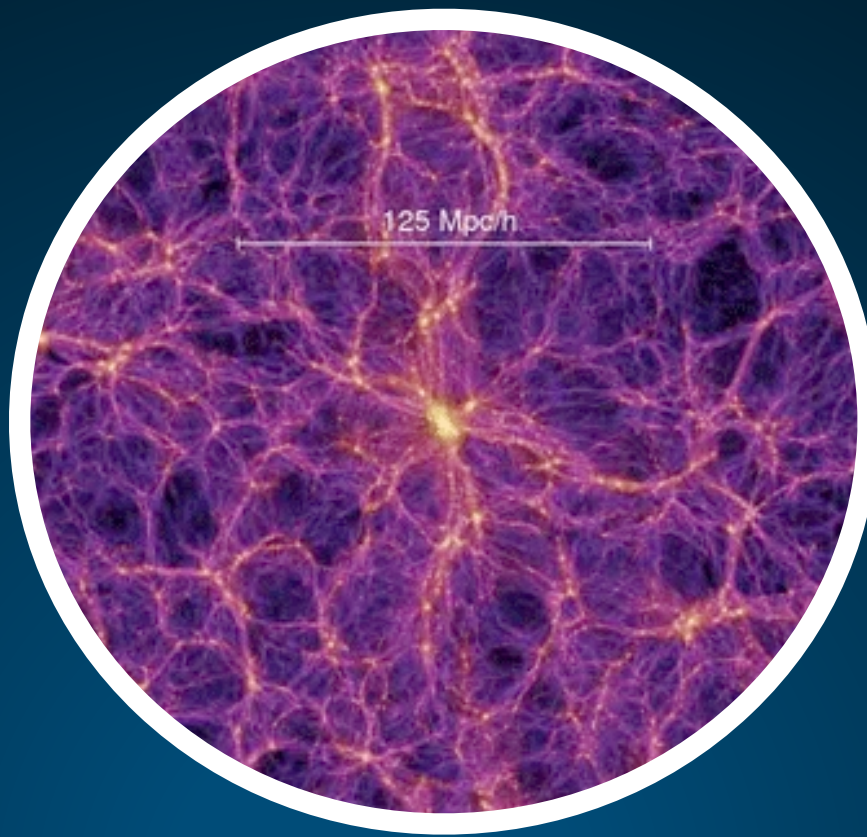
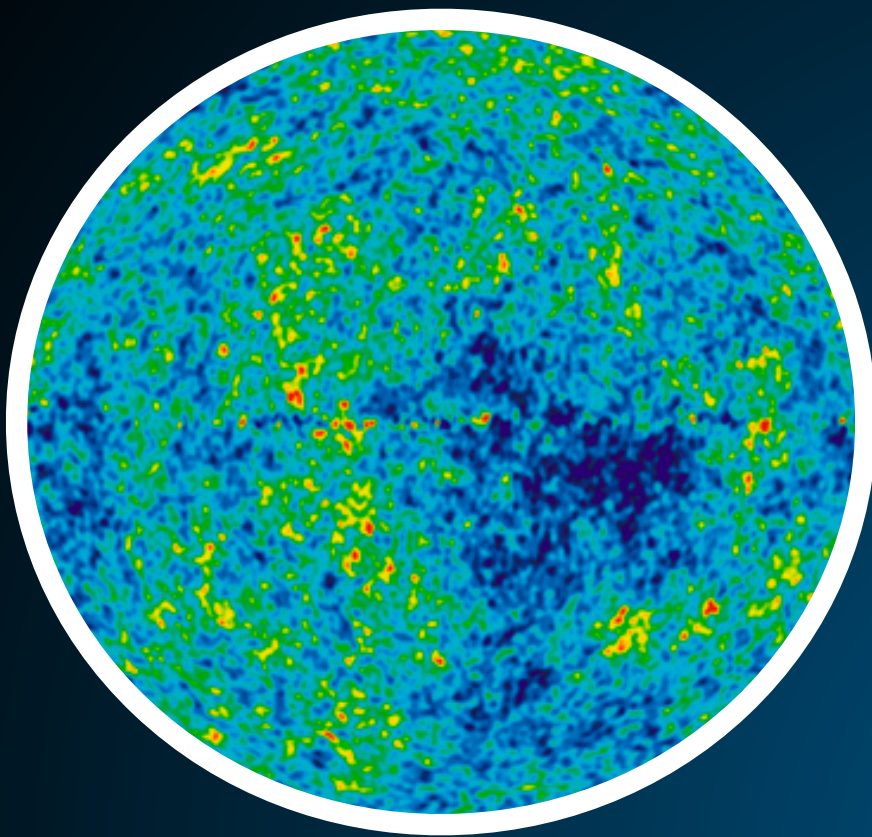
Stable - *on cosmological timescales*

Dark - *negligible electromagnetic cross-section*

Cold - *not relativistic*

~5.3x as prevalent as baryonic matter

Everything We Know About Dark Matter



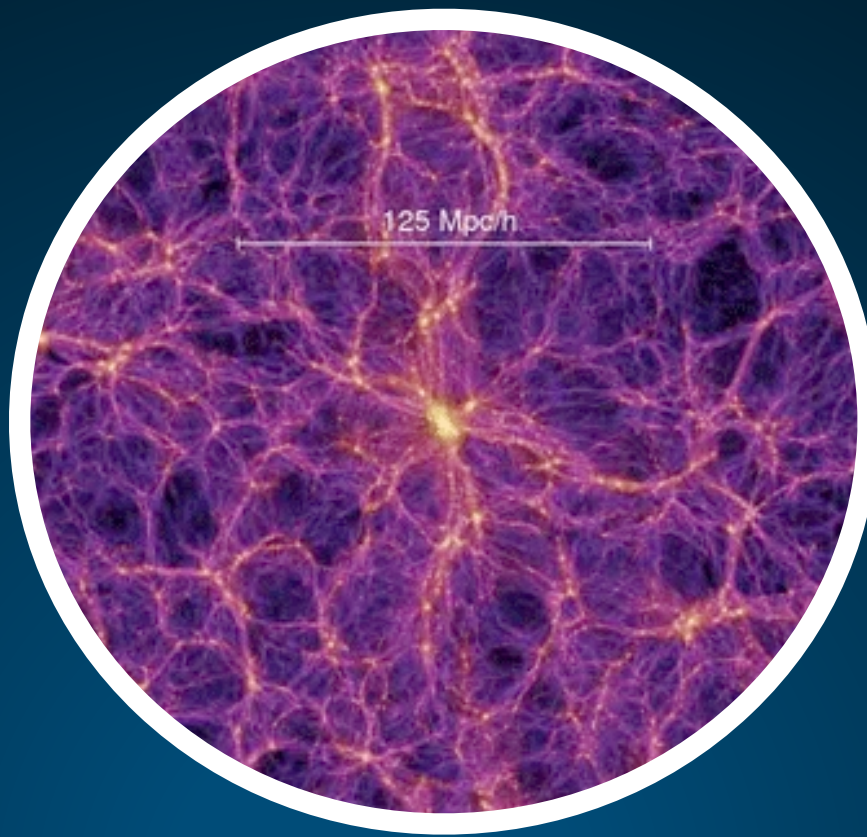
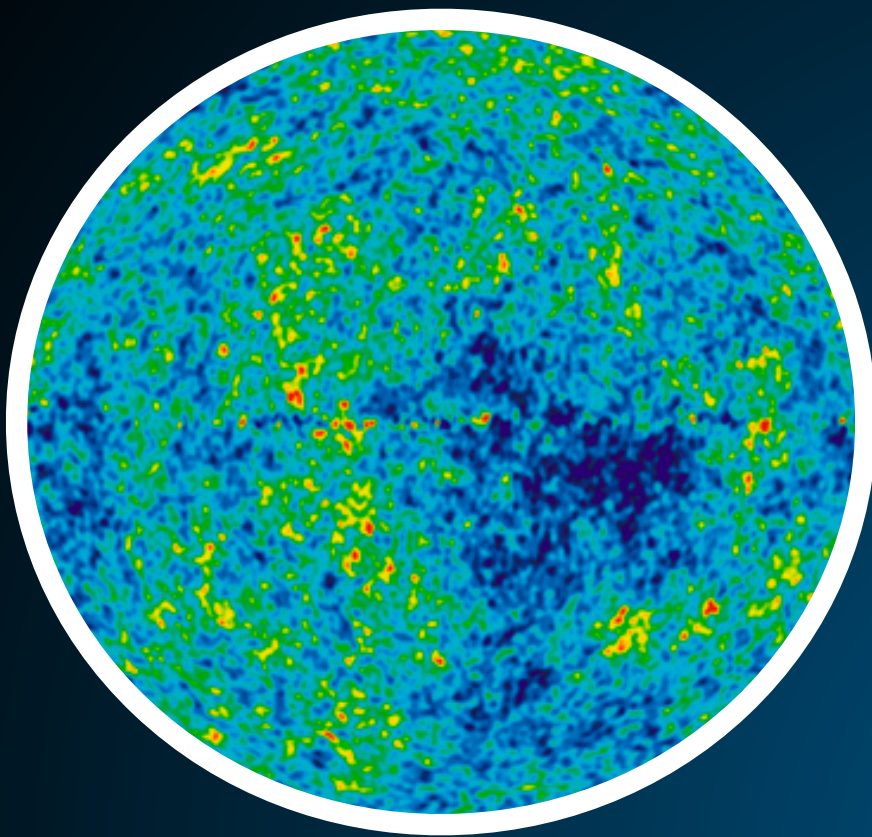
10^{-25} GeV

$\sigma_X > R_{\text{UFD}}$

10^{62} GeV

$m_X > M_{\text{UFD}}$

Everything We Know About Dark Matter



~100 GeV

**Thermal
Dark Matter**

10^{-25} GeV

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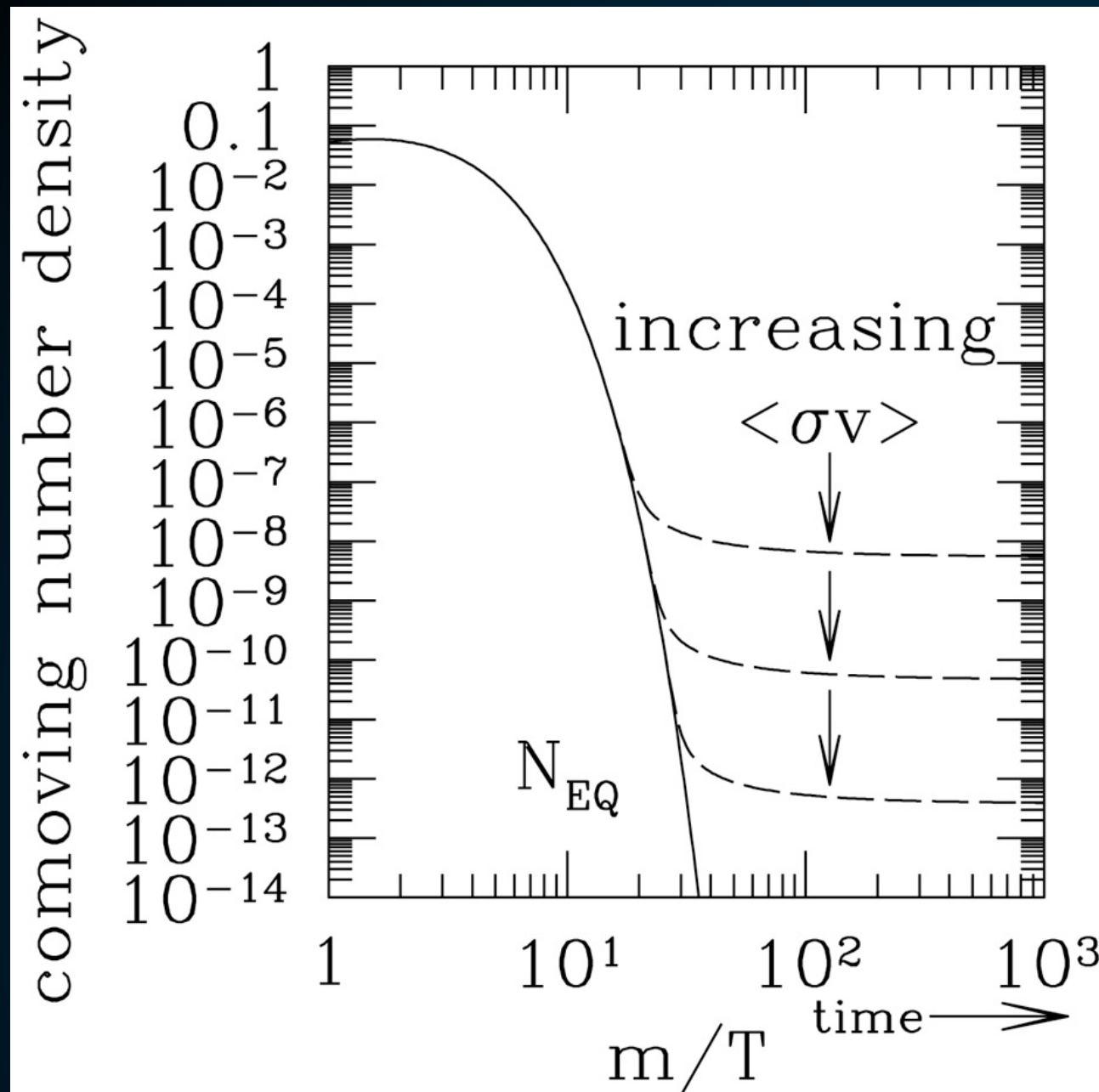
$m_X > M_{\text{UFD}}$

Thermal Dark Matter

artist: Sarah Szabo



The WIMP Miracle



A particle with a weak interaction cross-section and a mass on the weak scale is expected to naturally obtain the correct relic abundance through thermal freeze-out in the Earth universe.

$$\frac{\Omega_\chi}{0.2} = \frac{x_{f.o}}{20} \left(\frac{10^{-8} \text{ GeV}^{-2}}{\sigma} \right)$$

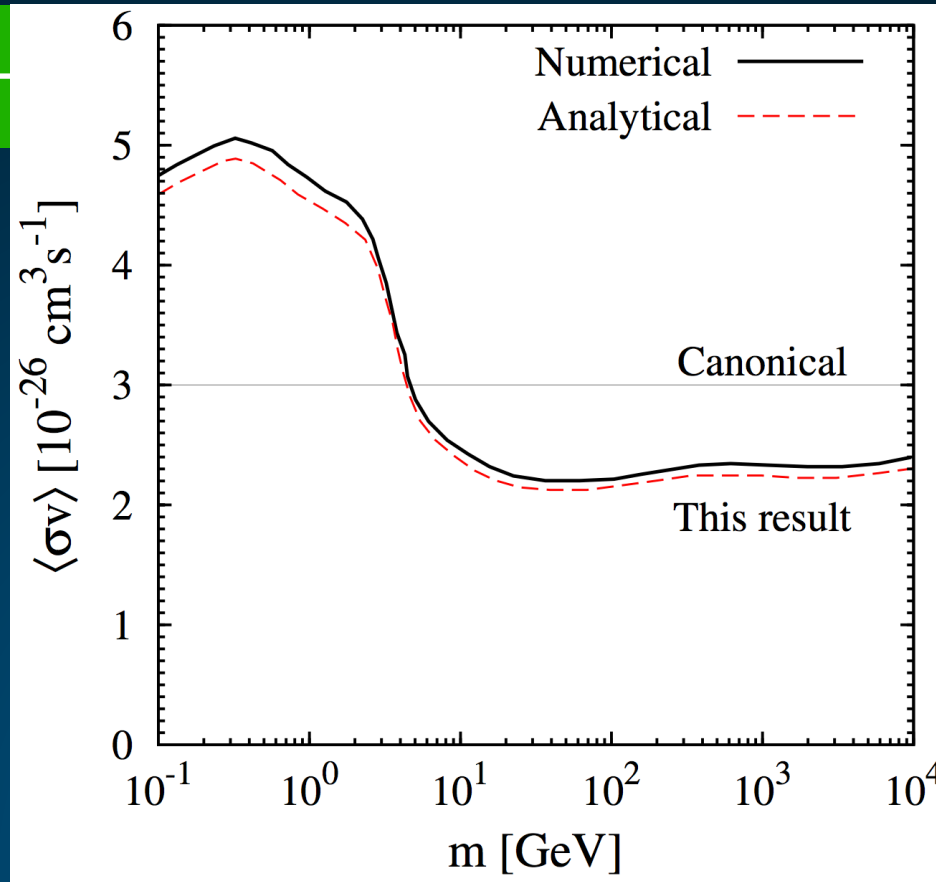
$$\langle \sigma v \rangle \sim 10^{-8} \text{ GeV}^{-2} \left(3 \times 10^{-28} \right) \text{ GeV}^2 \text{ cm}^2 10^{10} \frac{\text{cm}}{\text{s}}$$

$$\langle \sigma v \rangle \sim 3 \times 10^{-26} \frac{\text{cm}^3}{\text{s}}$$

**N_{eff}
>3 MeV**

$$N_{\text{eff}} = N_\nu \left(\frac{4}{11} \right)^{-4/3} \left(\frac{T_\nu}{T_\gamma} \right)^4$$

$$N_{\text{eff}}^{\text{equil}} = N_\nu \left[1 + \frac{1}{N_\nu} \sum_{i=1}^n \frac{g_i}{2} F \left(\frac{m_i}{T_d} \right) \right]^{4/3}$$



Steigman et al. (1204.3622)

**Unitarity
< 124 TeV**

$$\sigma \sim \frac{g^2}{m_\chi^4}$$

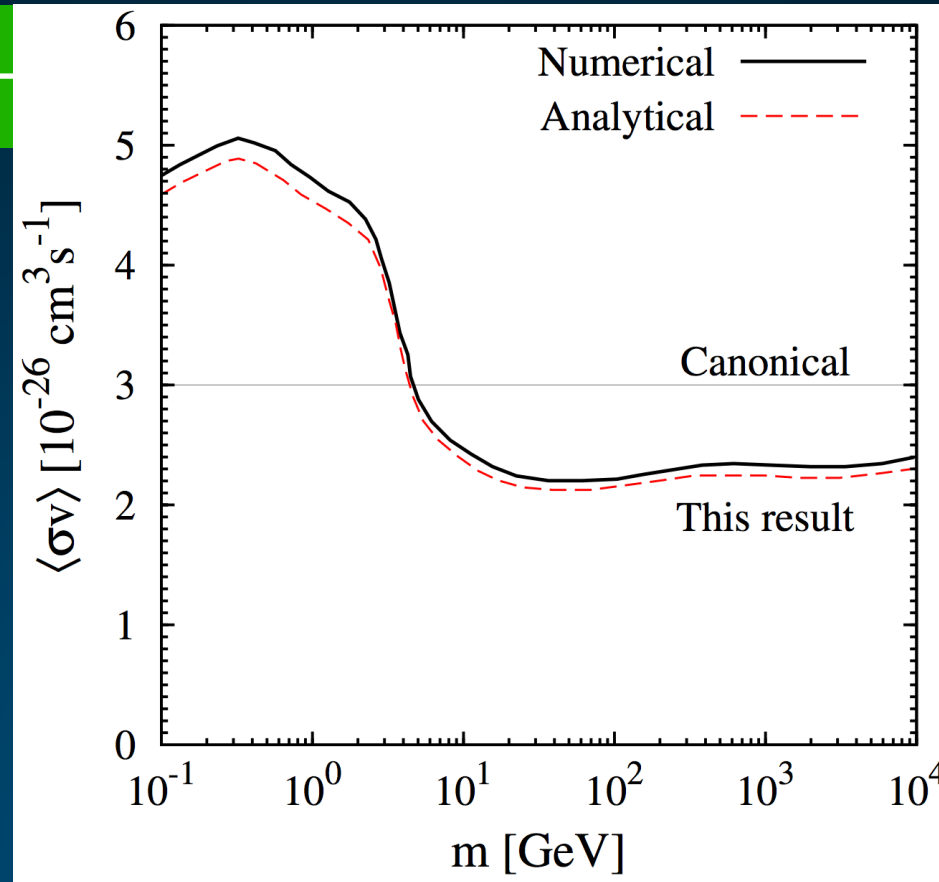
$$g \lesssim 4\pi$$

$$\frac{\Omega_\chi}{0.2} \gtrsim 10^{-8} \text{ GeV}^{-2} \left(\frac{m_\chi^2}{4\pi} \right)$$

$$x_f = \ln \left(\frac{m_\chi^3 \langle \sigma v \rangle}{H_{x=1}} \right)$$

**Lee-Weinberg
>10 GeV**

if interaction is due
to the weak force



Steigman et al. (1204.3622)

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$$g \lesssim 4\pi$$

$$\frac{\Omega_\chi}{0.2} \gtrsim 10^{-8} \text{ GeV}^{-2} \left(\frac{m_\chi^2}{4\pi} \right)$$

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Thermal WIMPs - The Most Boring Model

1 new particle (can be motivated by more-complex physics)

1 new conserved quantity ("dark matter-ness", r-parity)

1 (maybe 0) new forces

Thermal WIMPs - The Most Boring Model

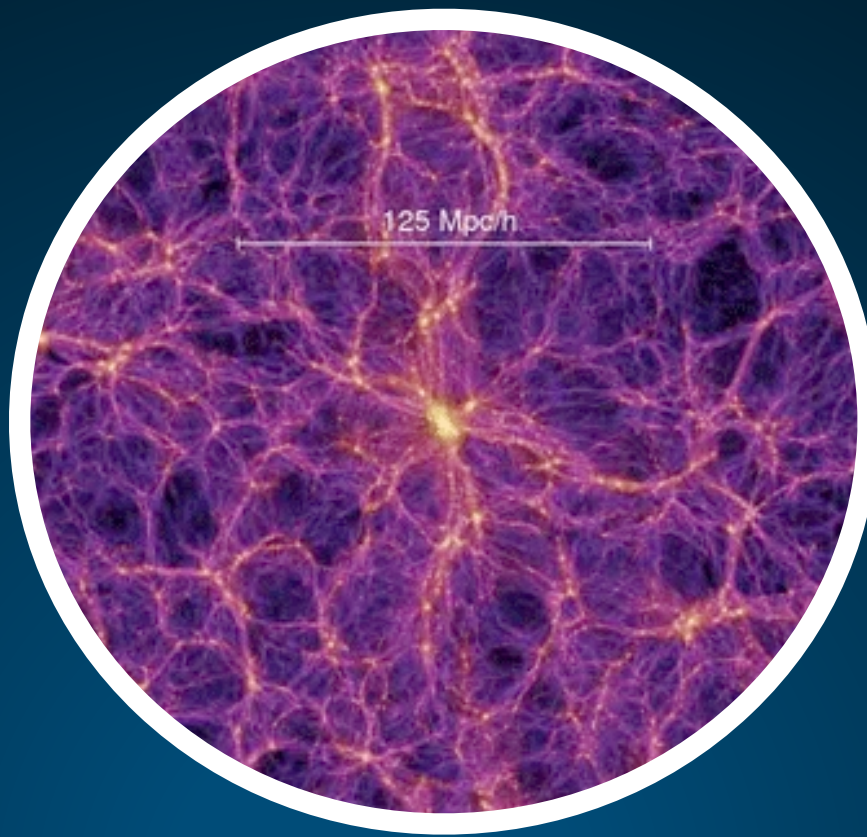
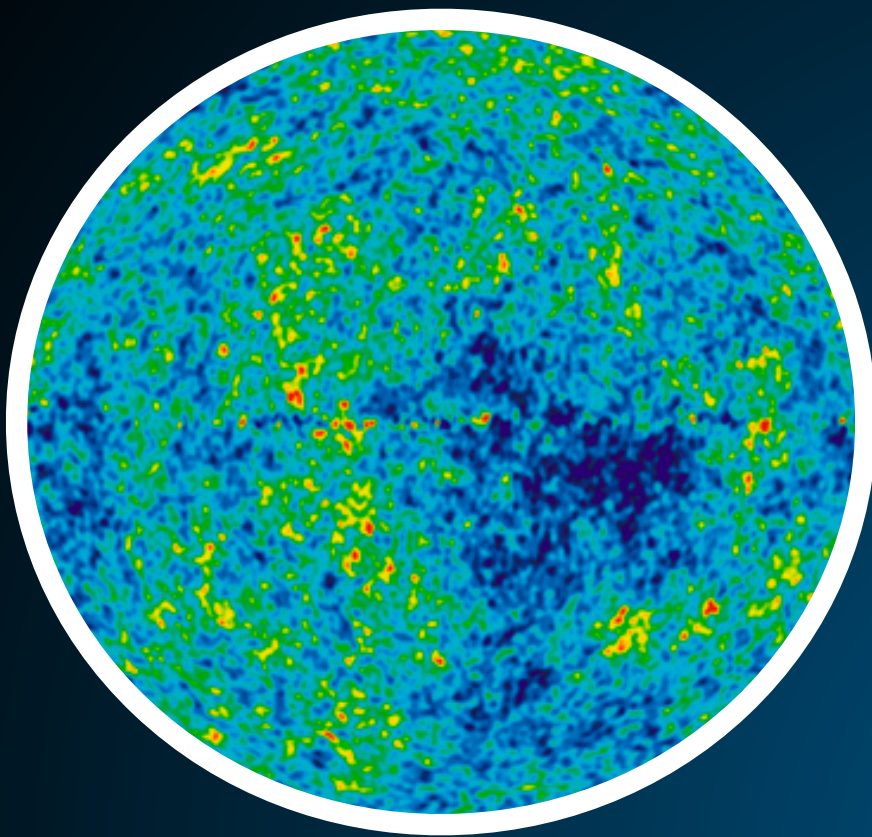
1 new particle (can be motivated by more-complex physics)

1 new conserved quantity ("dark matter-ness", r-parity)

1 (maybe 0) new forces

**Ruling out this model leaves only
*more interesting possibilities.***

Everything We Know About Dark Matter



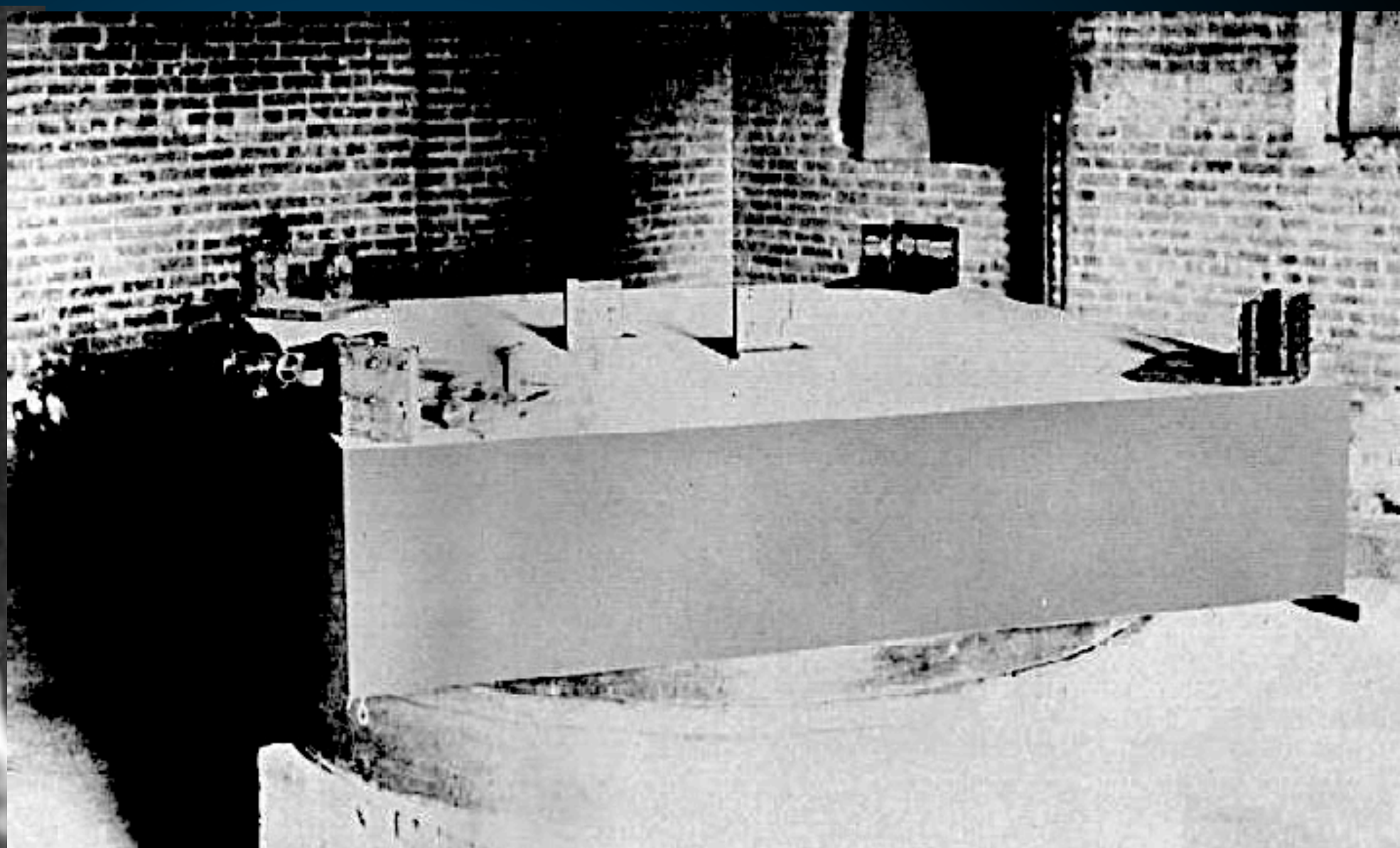
10^{-25} GeV

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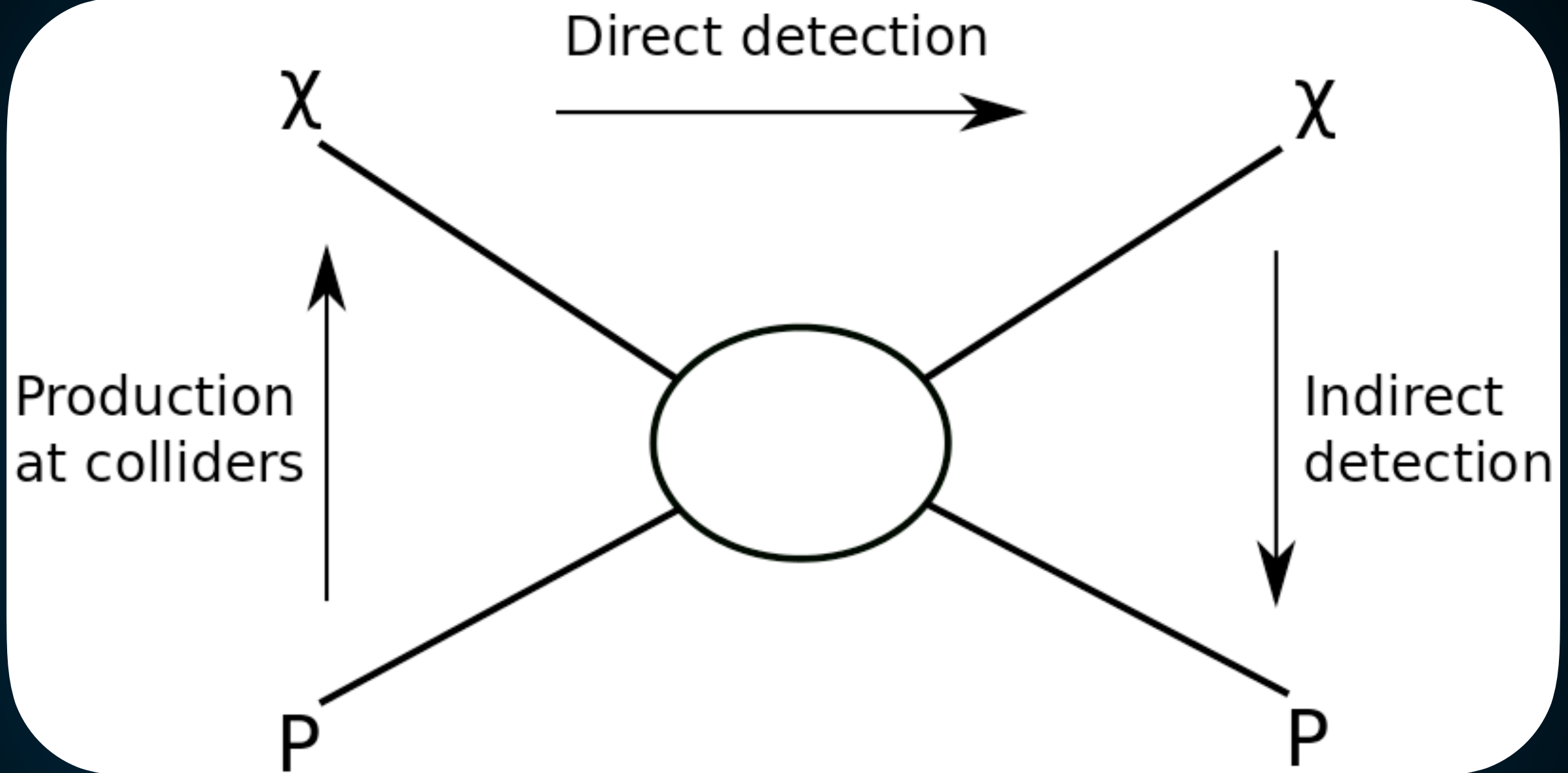
BORING

10^{62} GeV

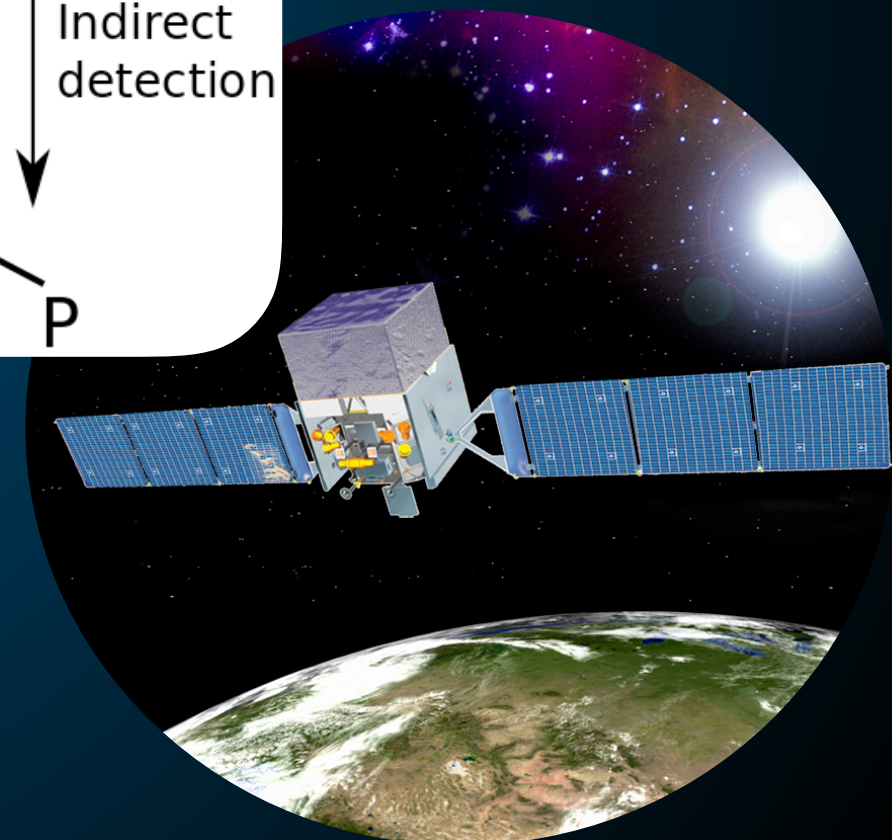
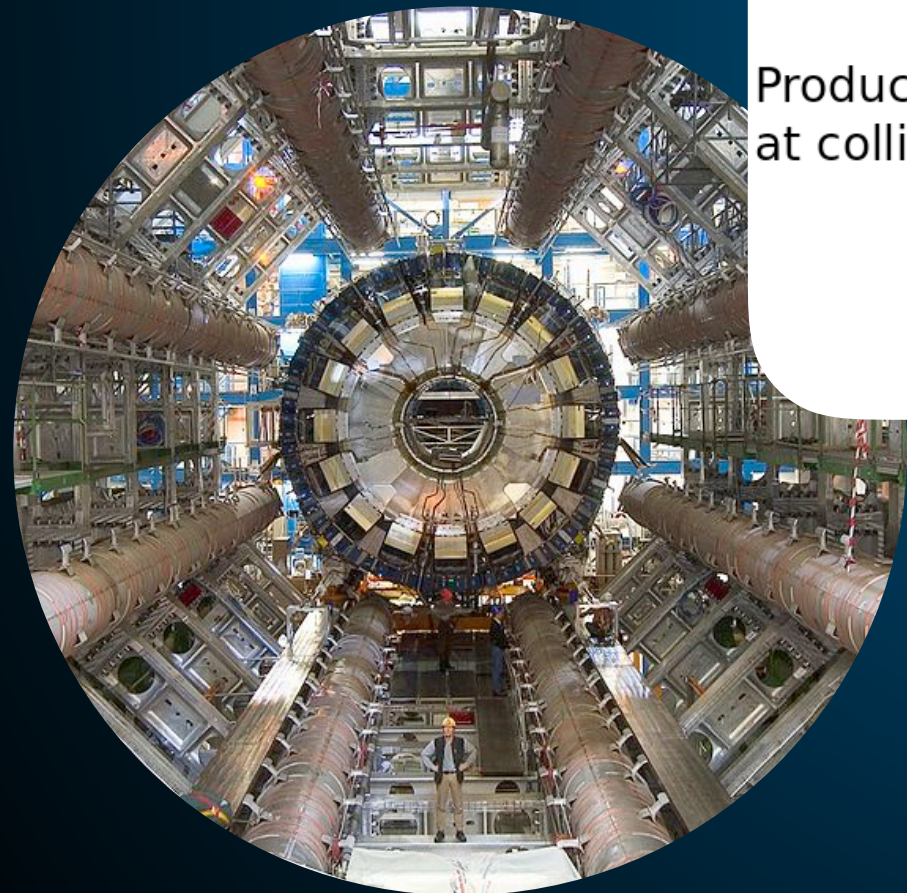
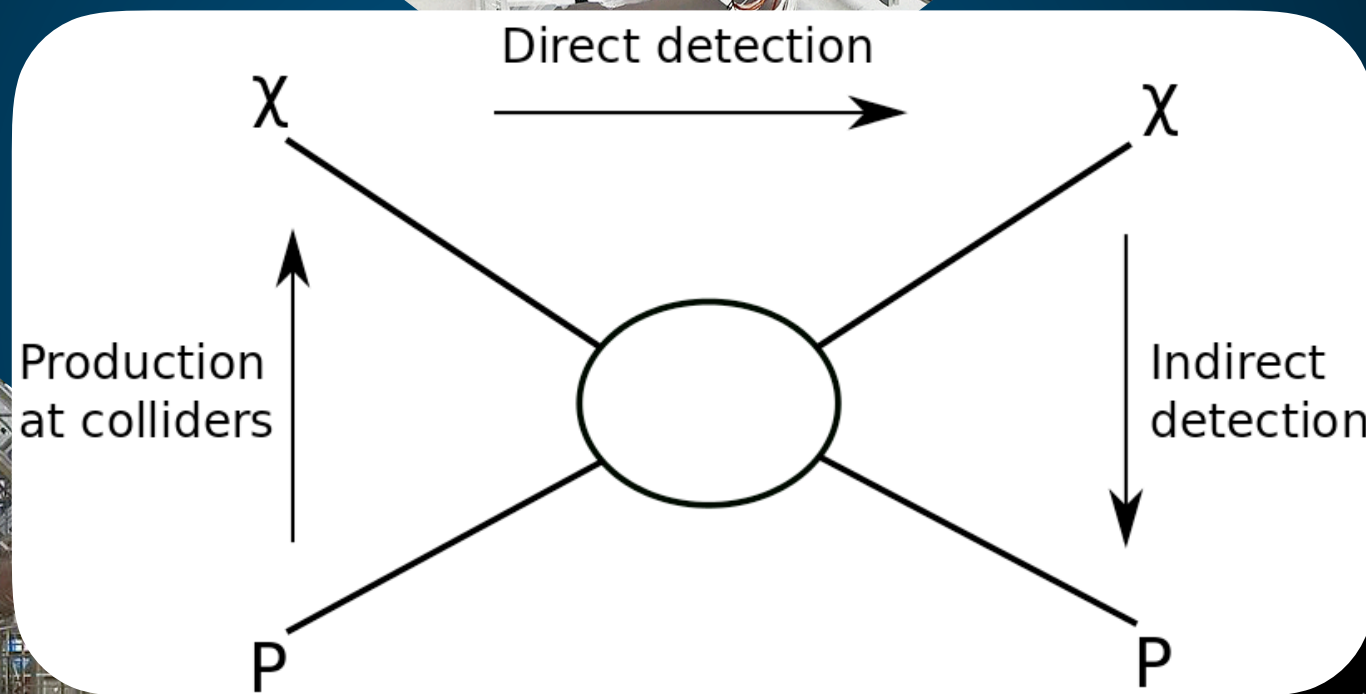
$m_X > M_{\text{UFD}}$



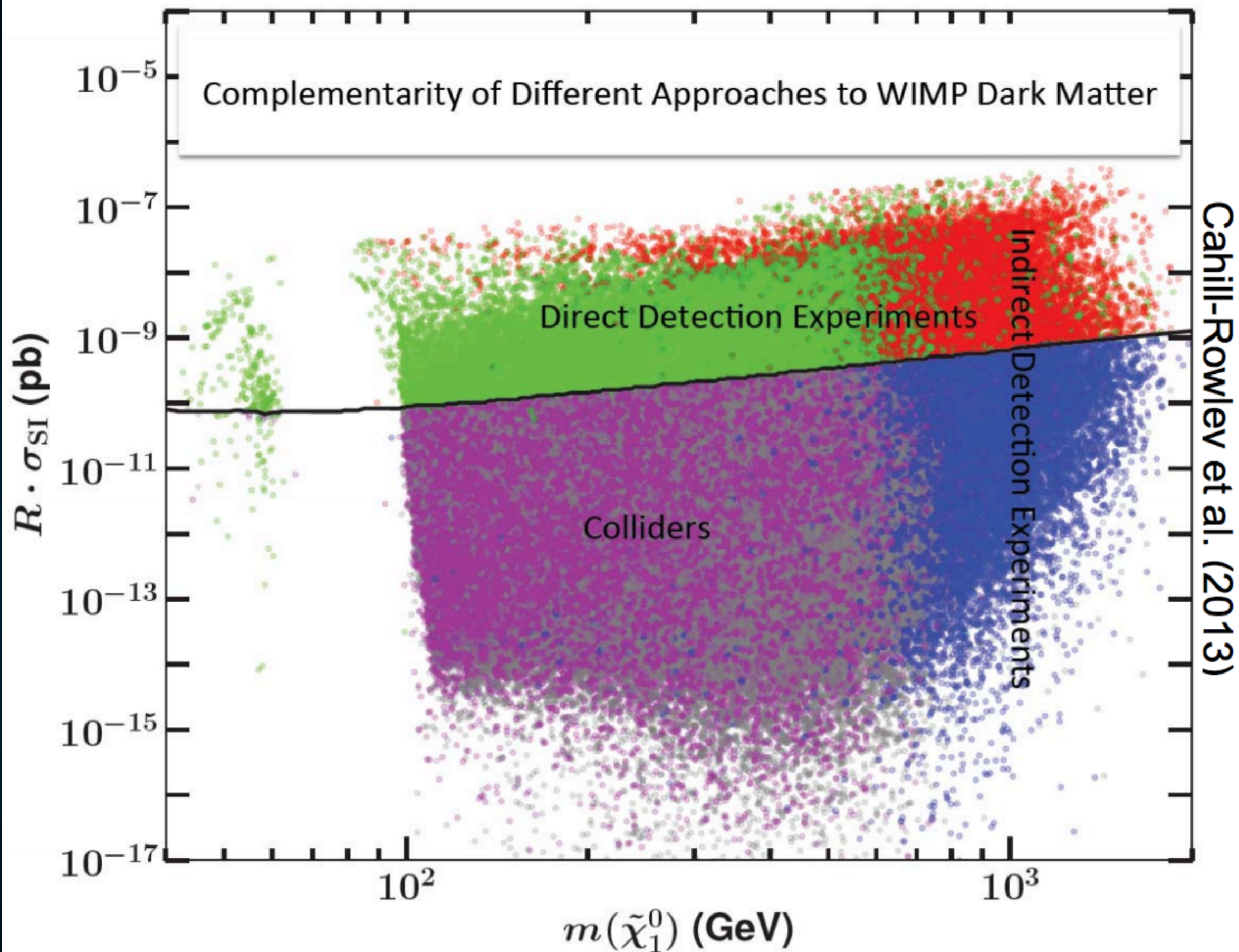
Dark Matter in Thermal Equilibrium

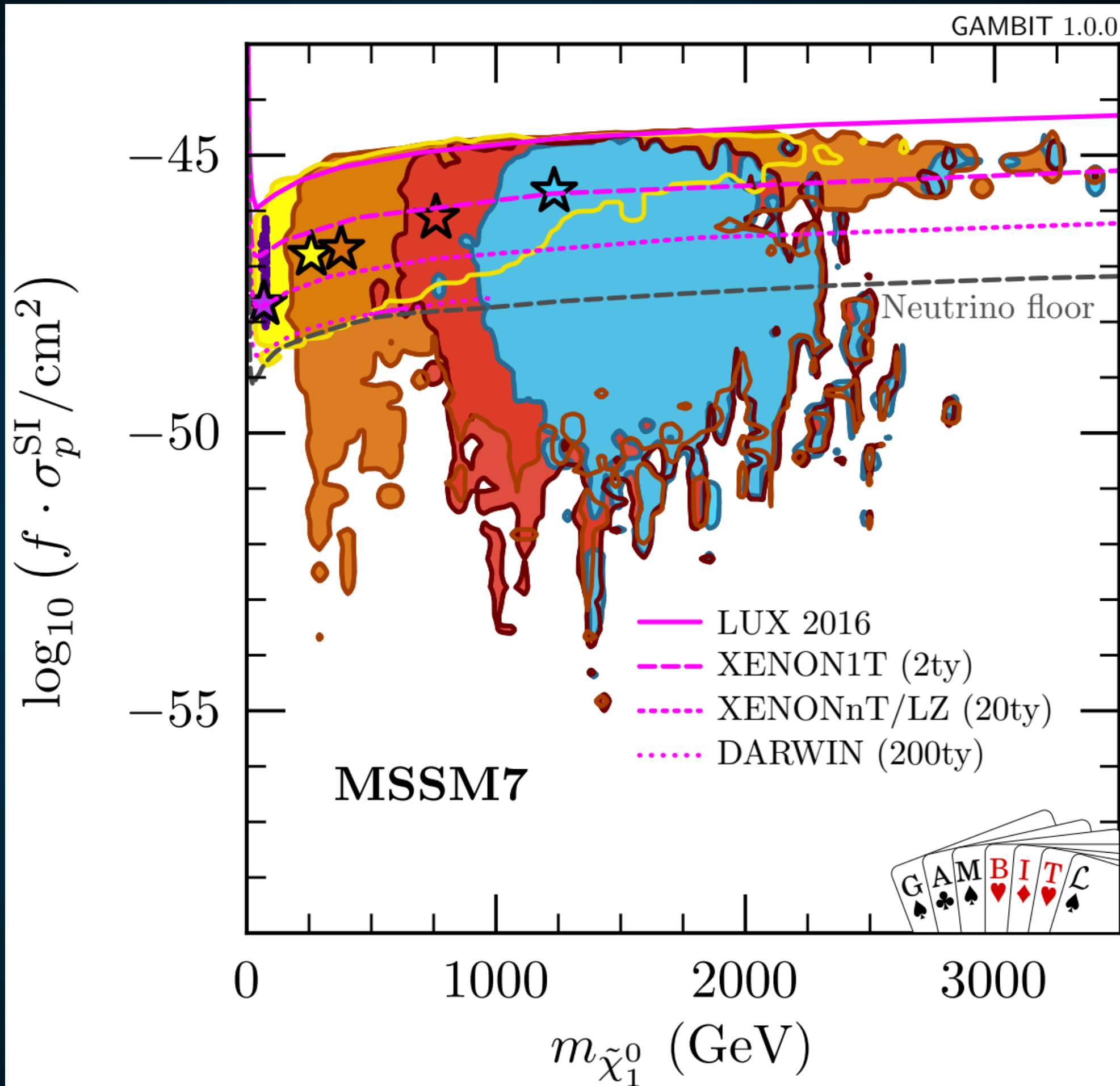


Dark Matter Complementarity



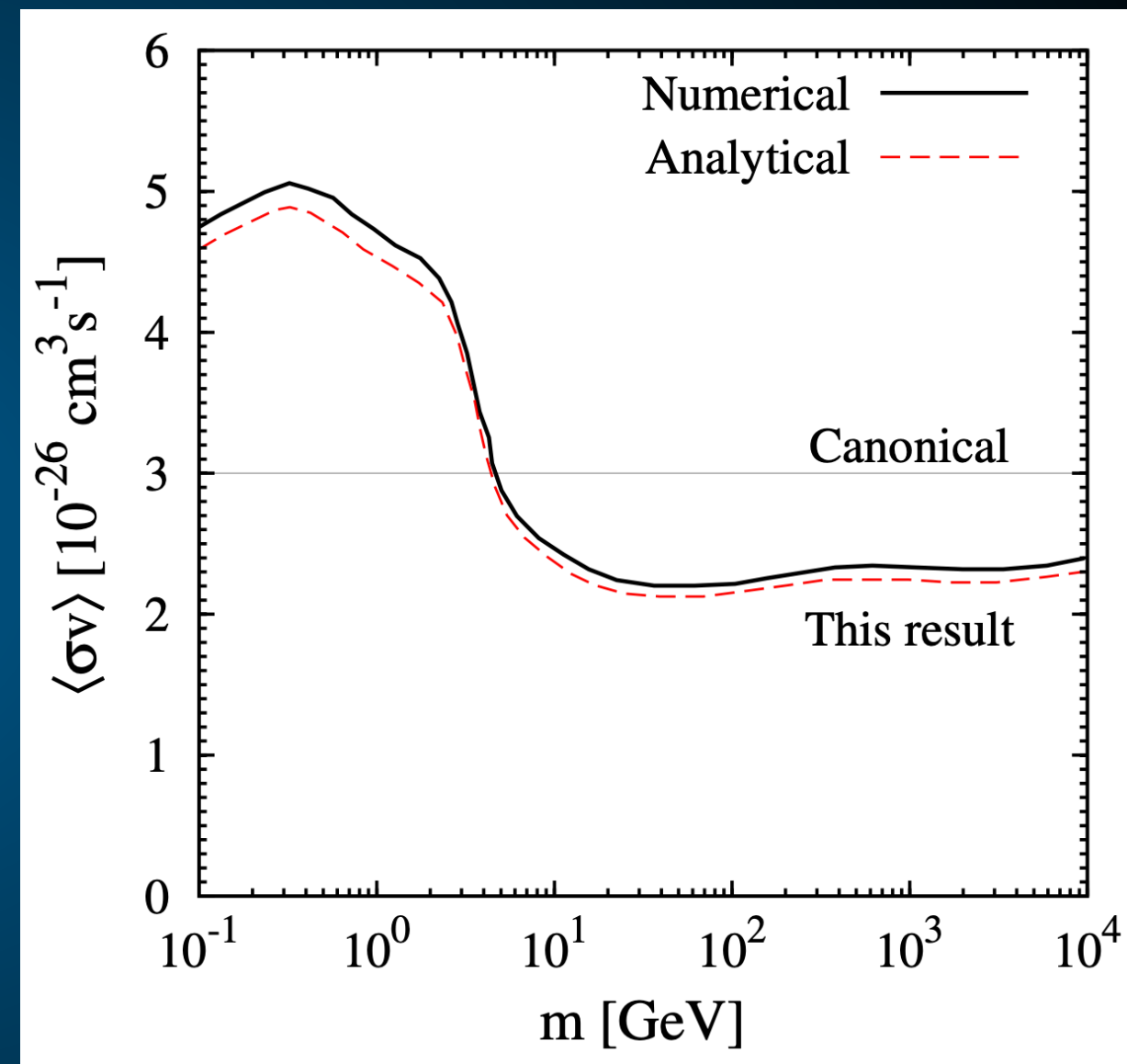
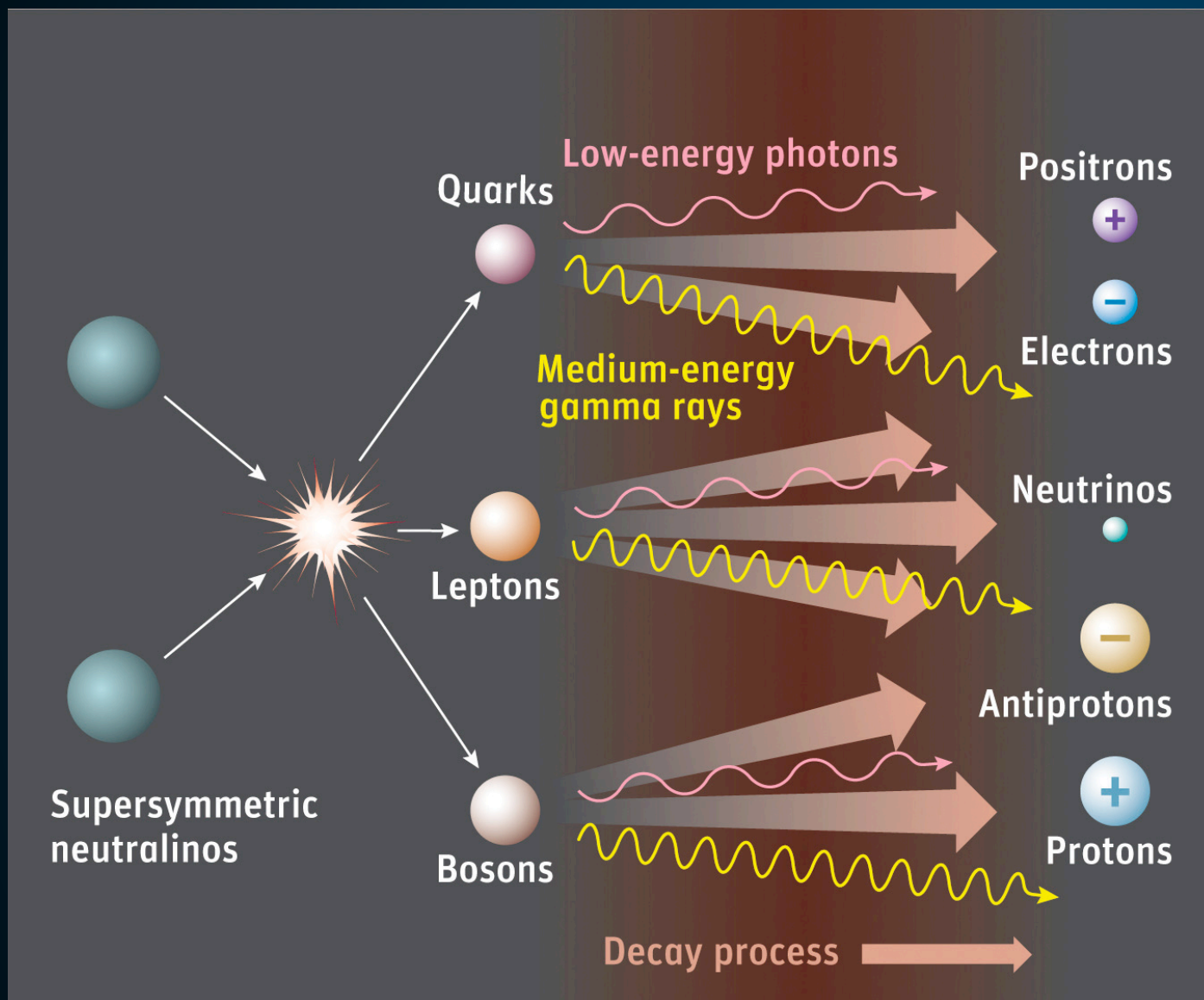
Dark Matter Complementarity





Why Indirect Detection?

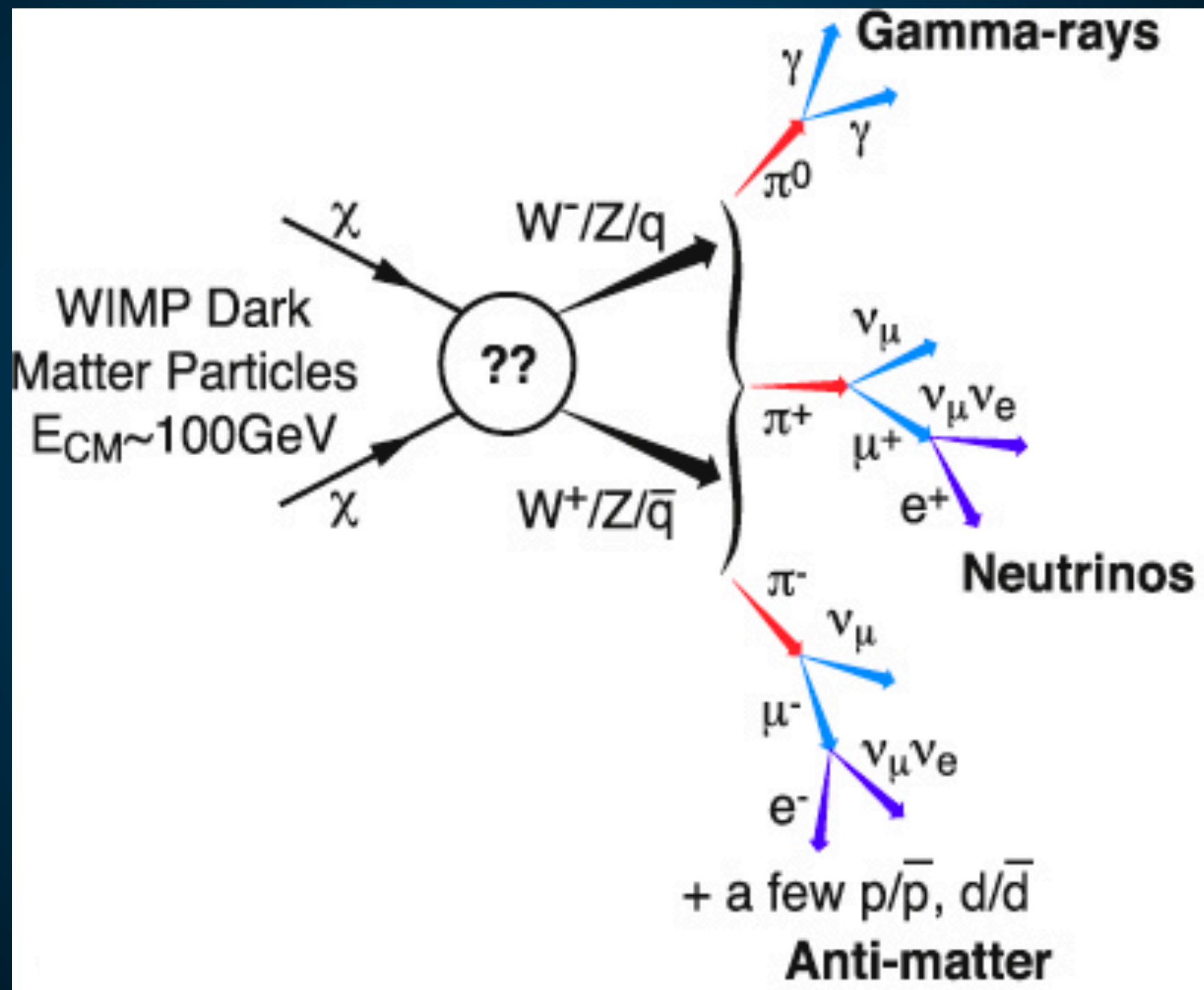
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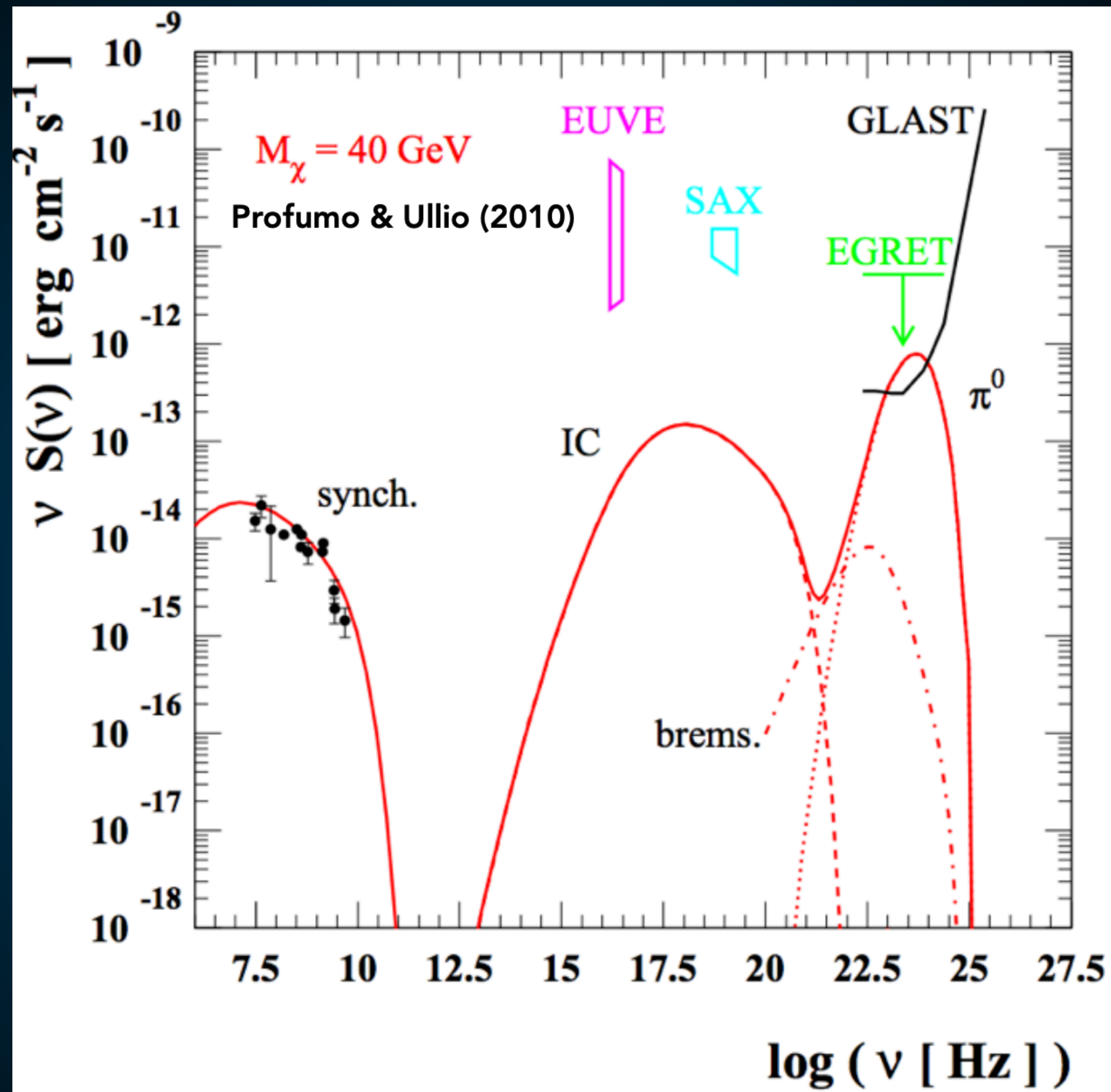
Steigman et al. (1204.3622)

How Do We Know Anything About the Energy of the Signals?

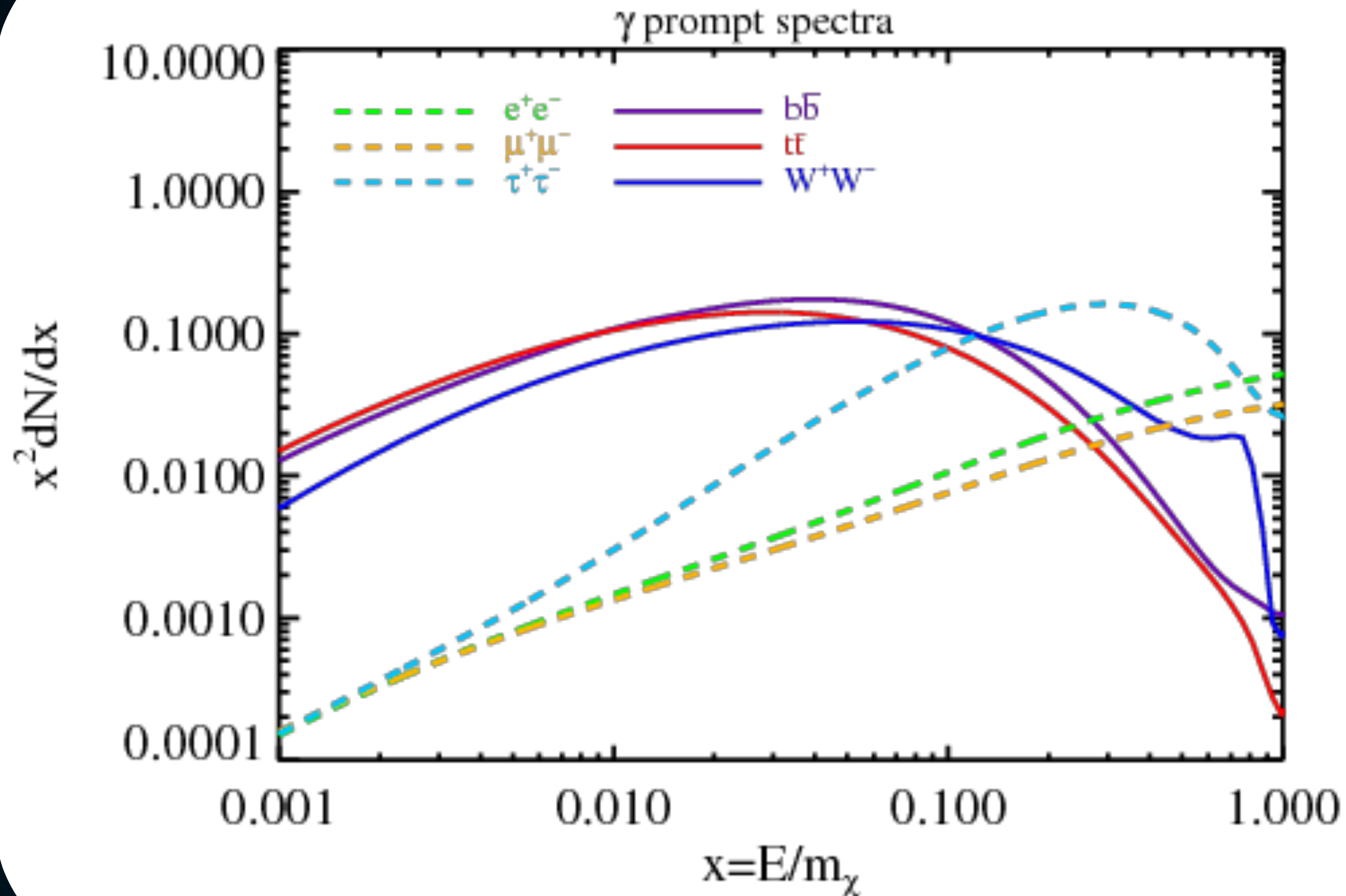
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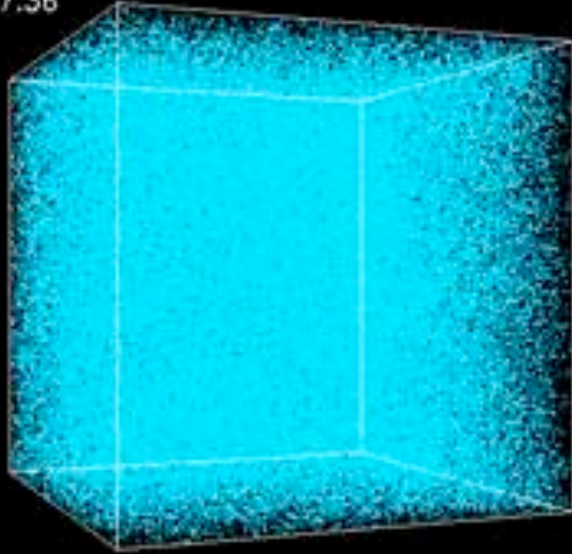


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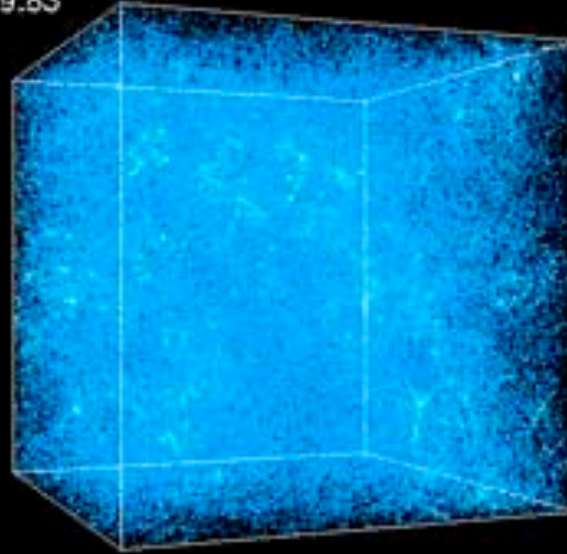


How Do We Know Anything About the Location of the Signals?

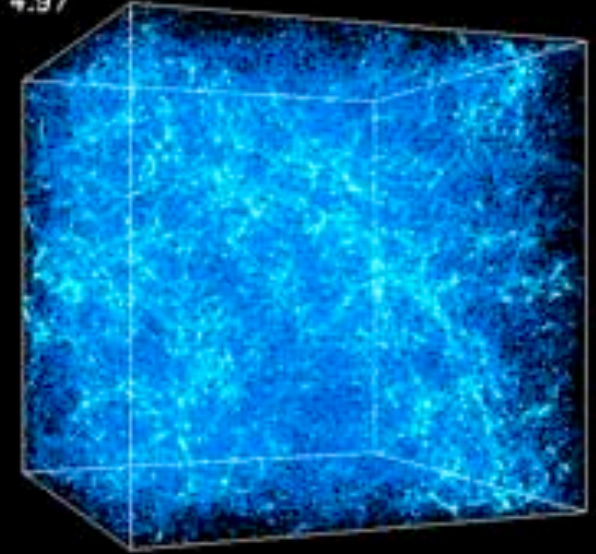
$Z=27.36$



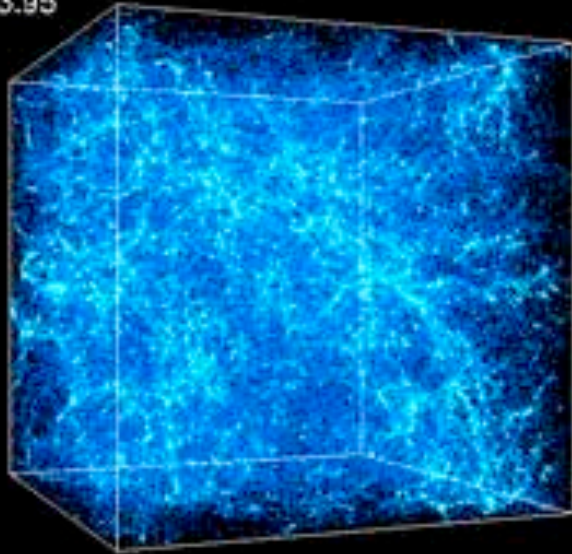
$Z= 9.83$



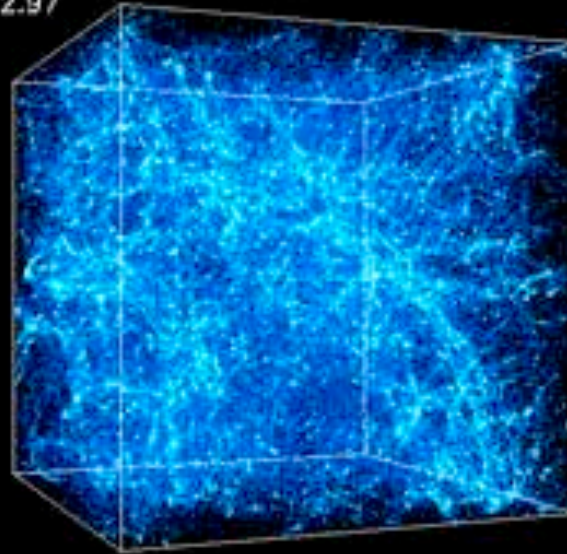
$Z= 4.97$



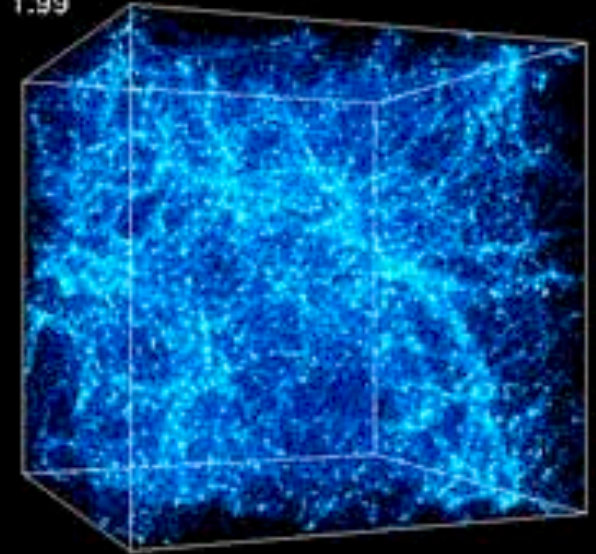
$Z= 3.95$



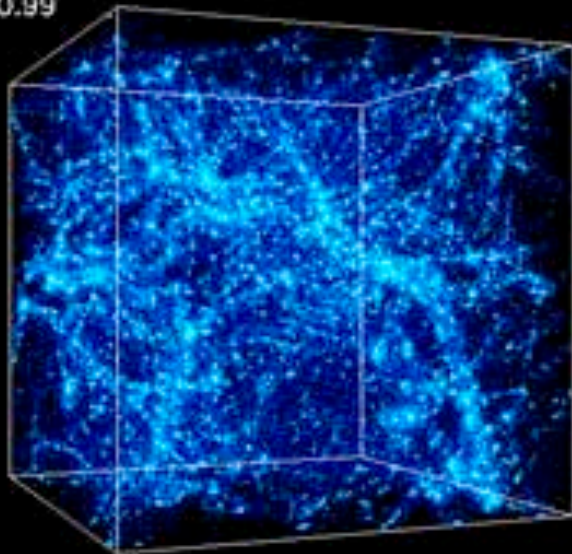
$Z= 2.97$



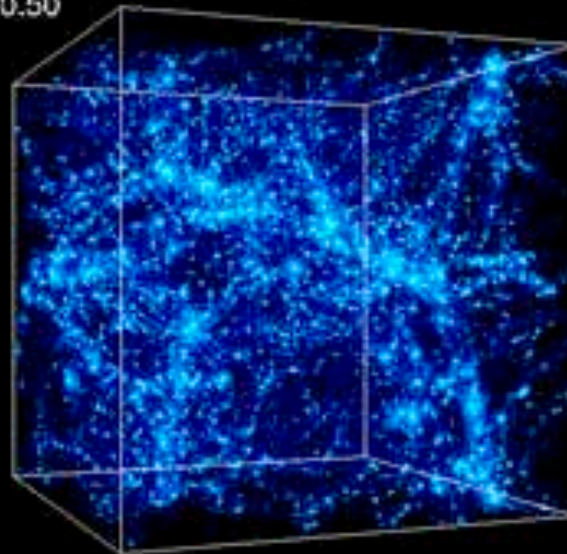
$Z= 1.99$



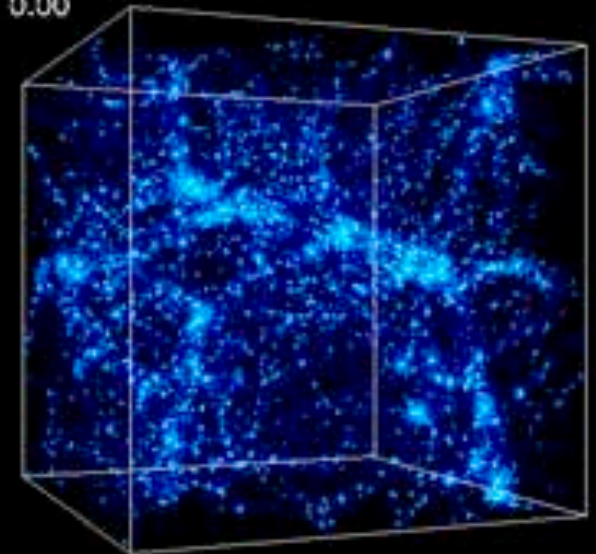
$Z= 0.99$



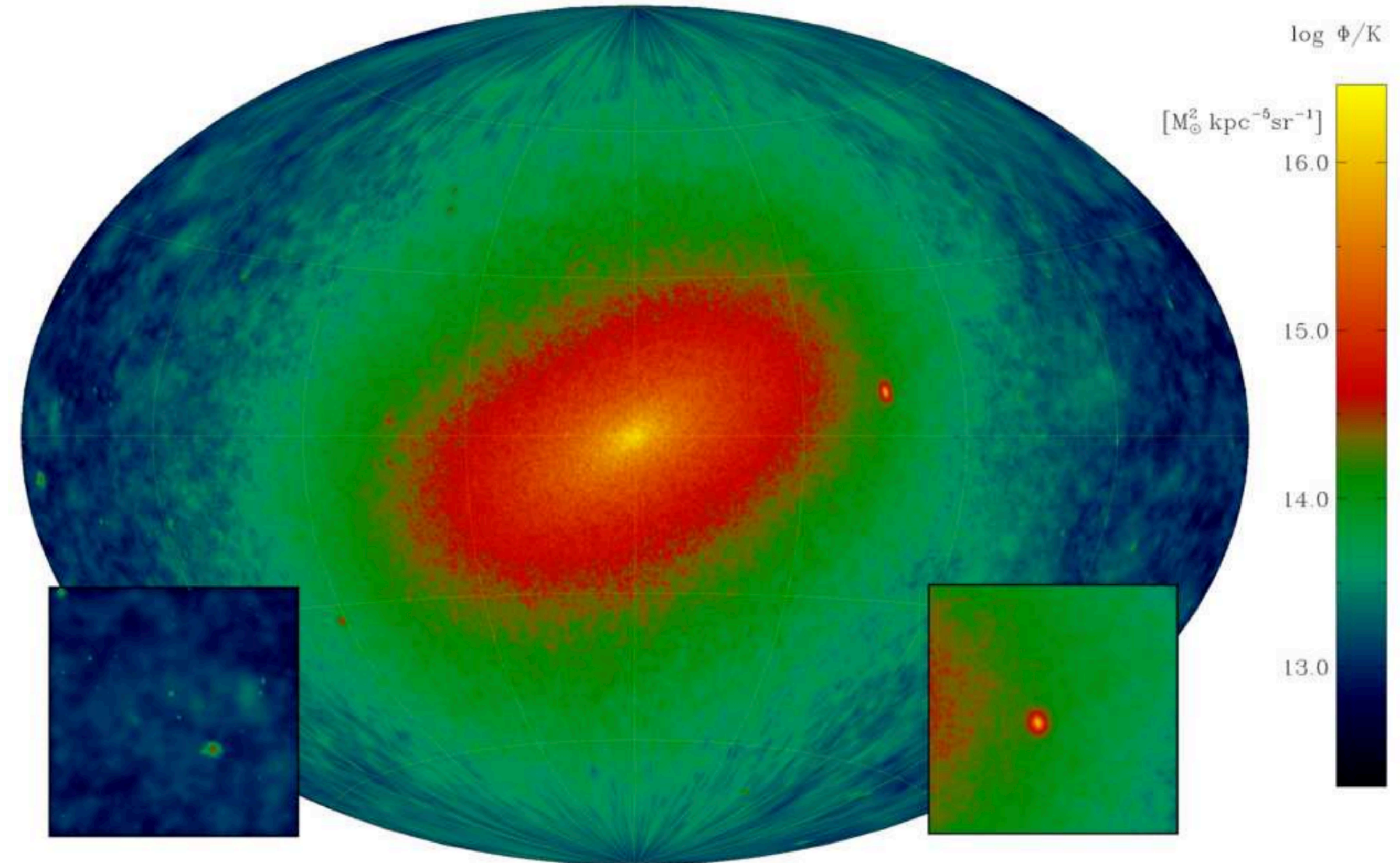
$Z= 0.50$



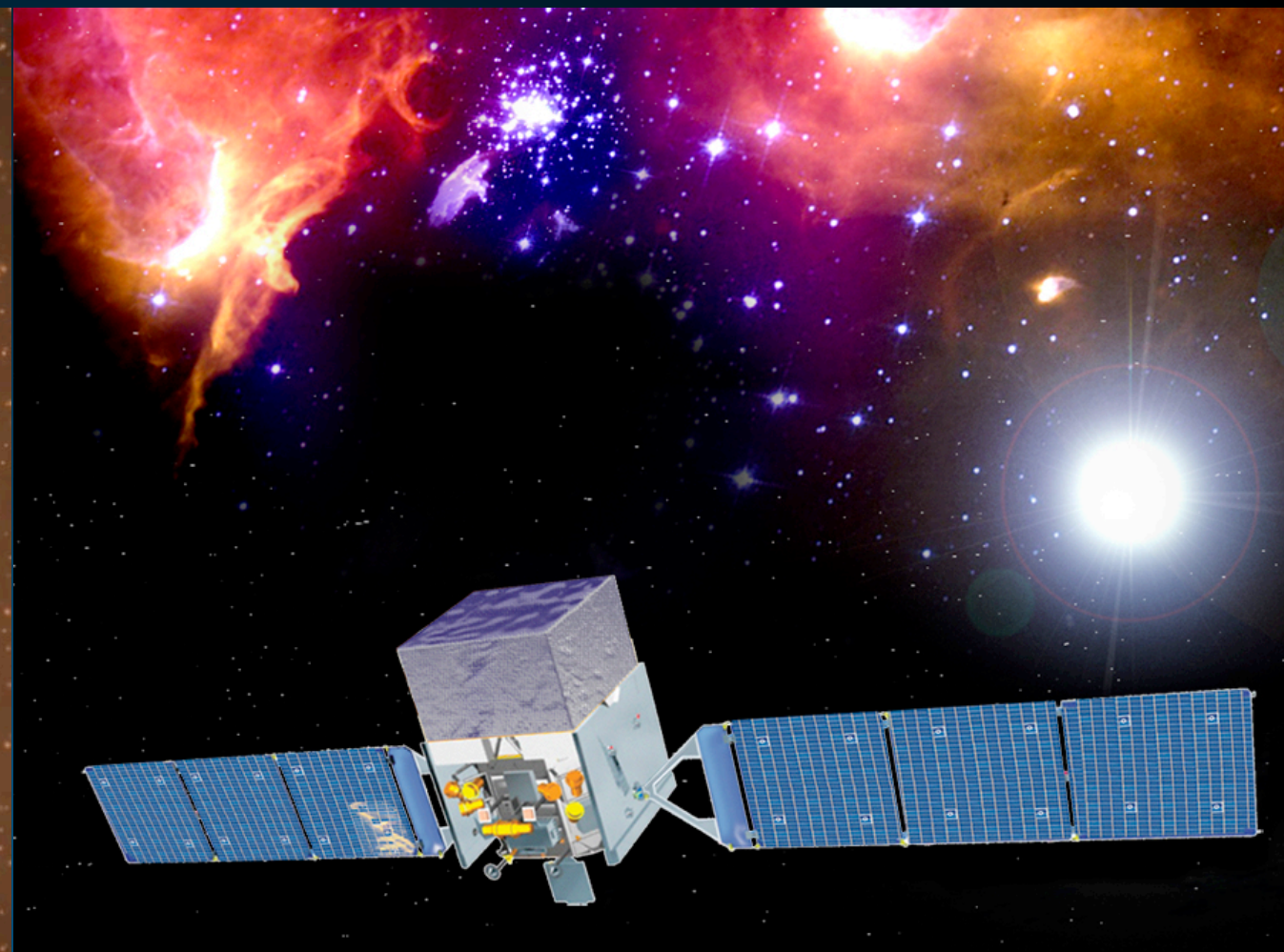
$Z= 0.00$



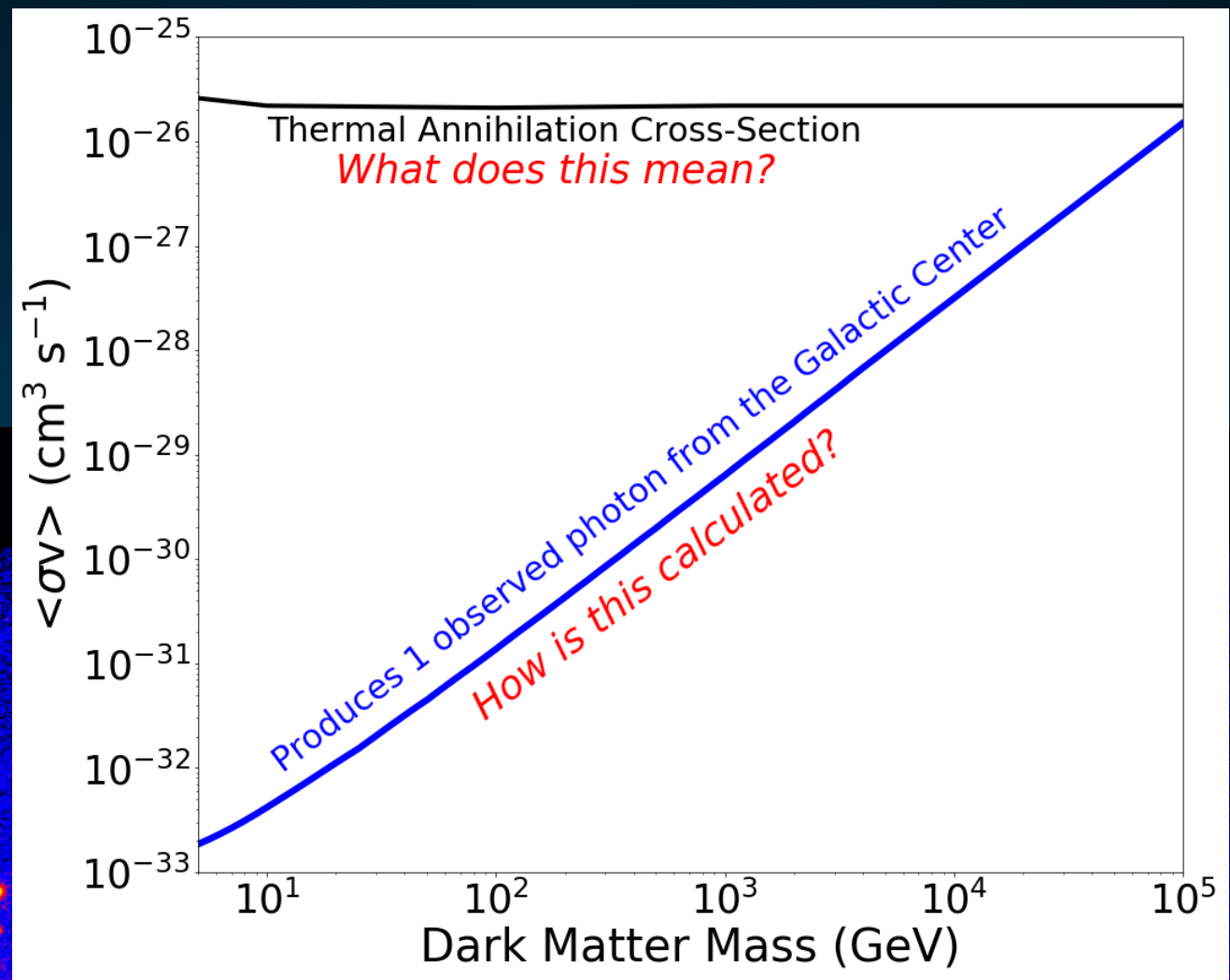
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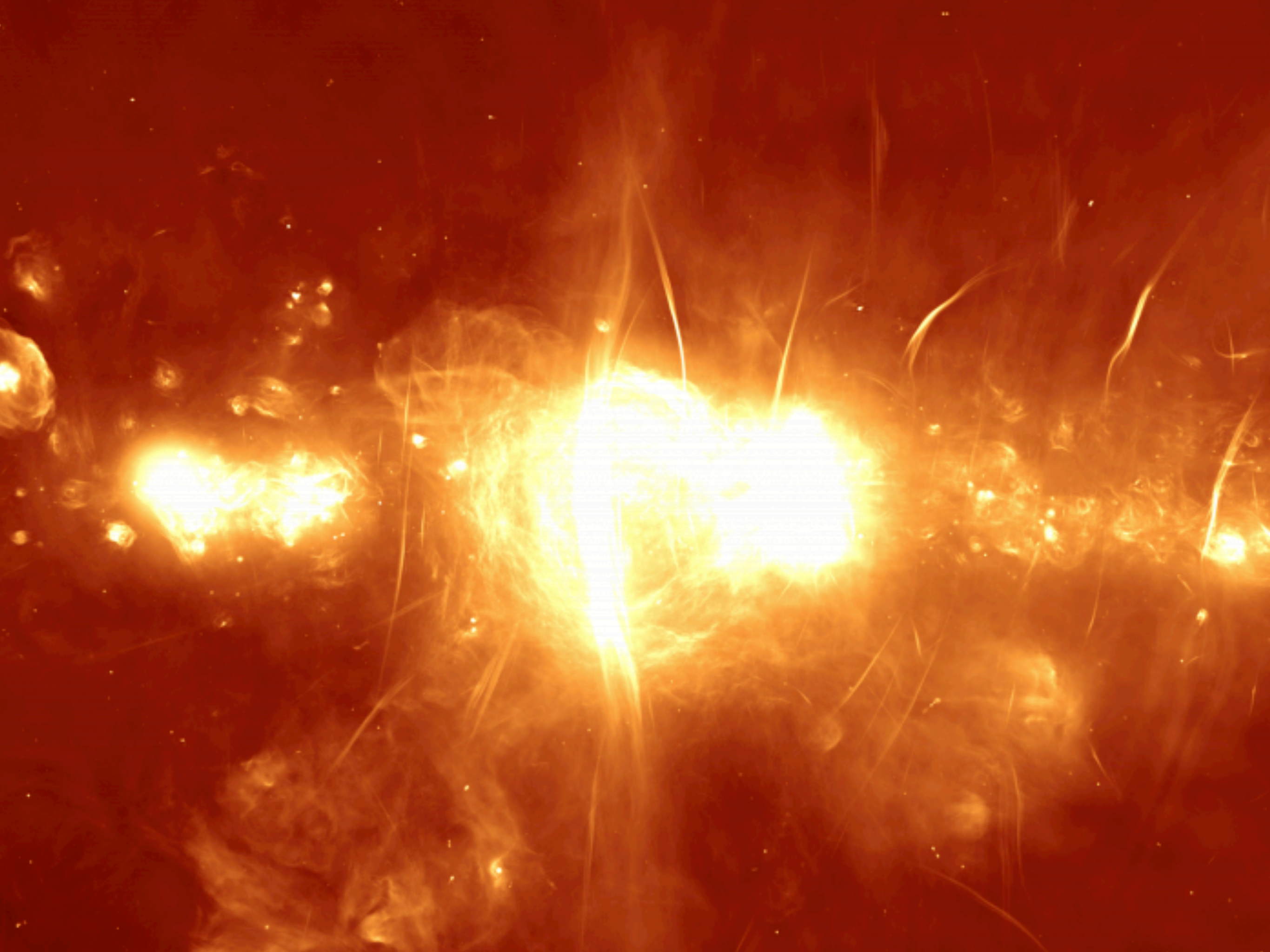


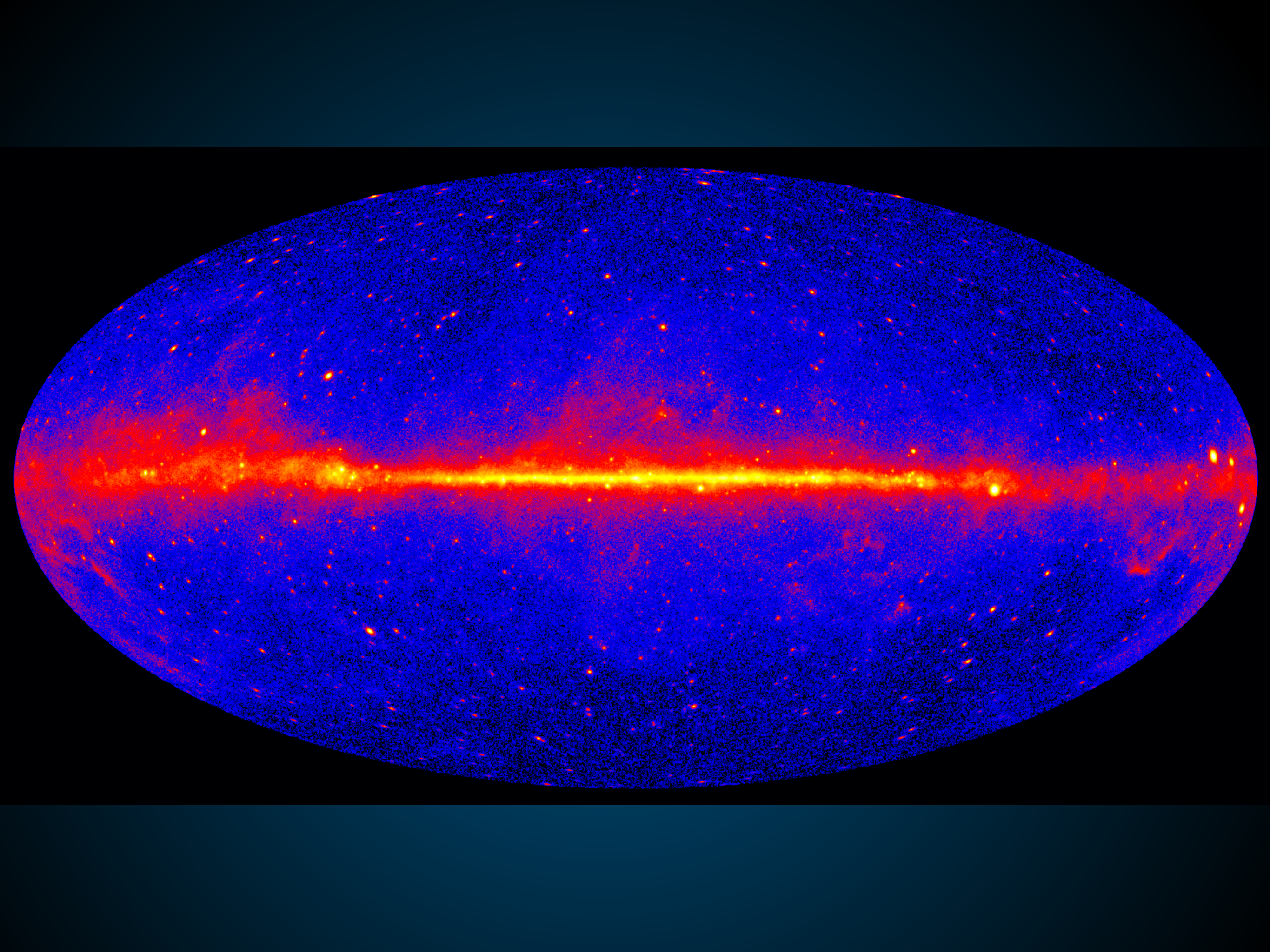
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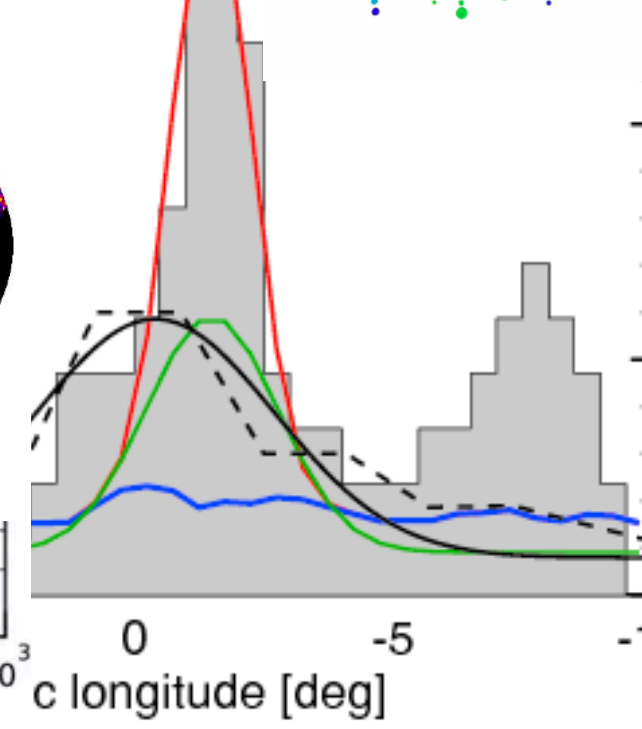
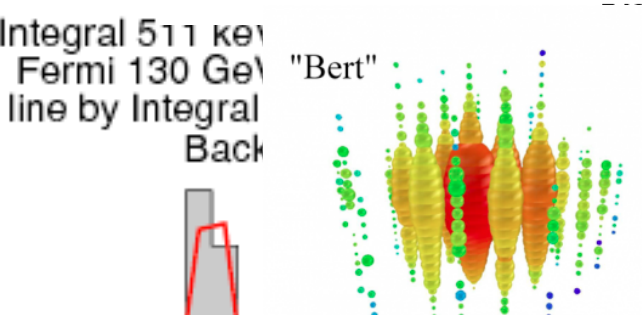
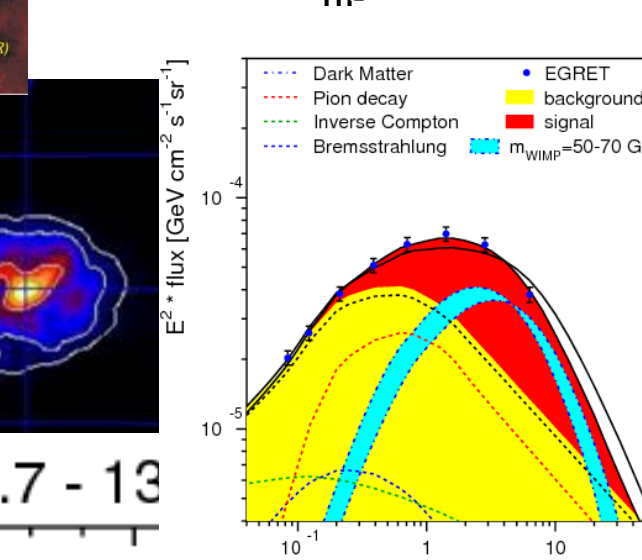
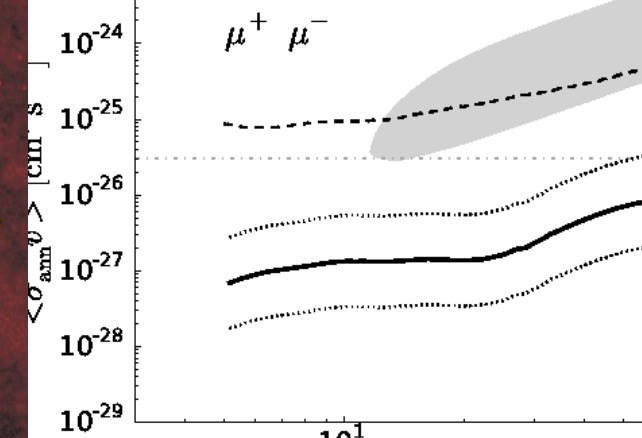
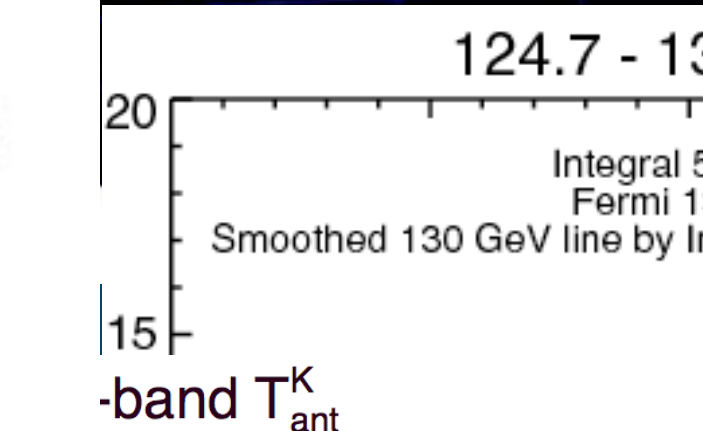
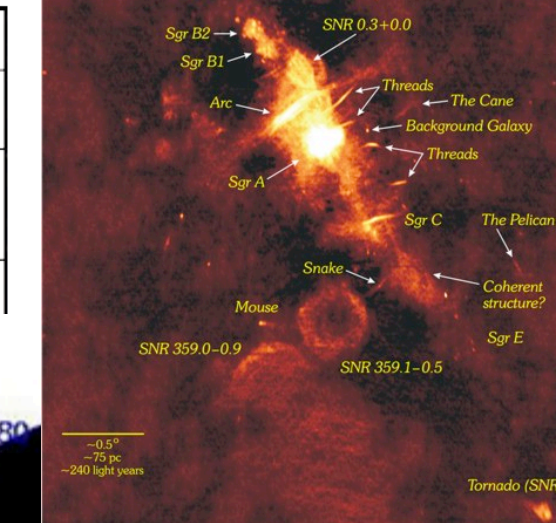
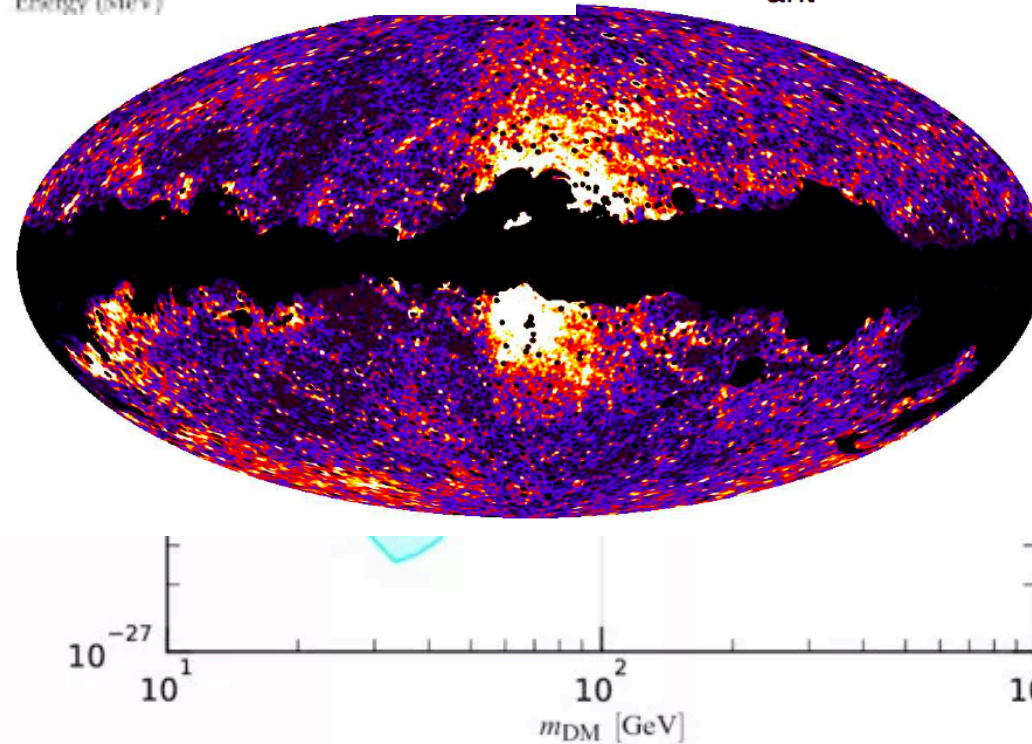
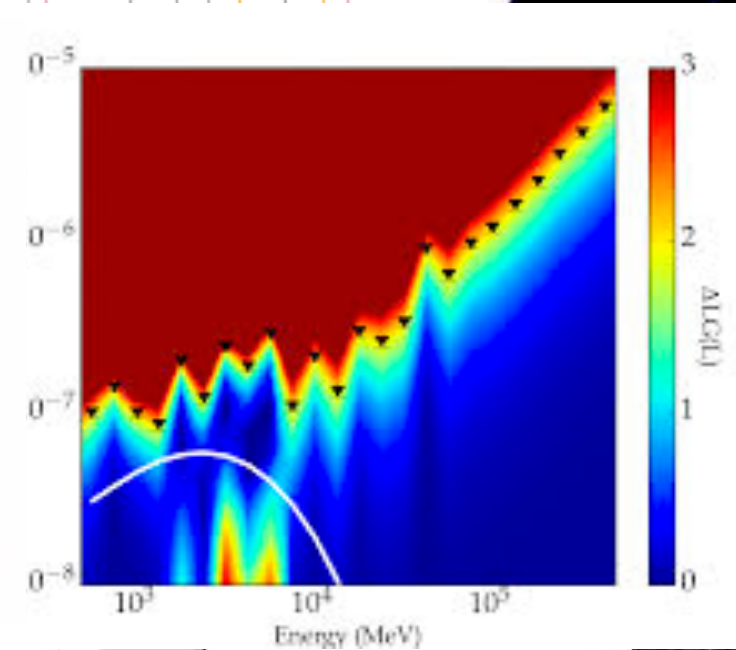
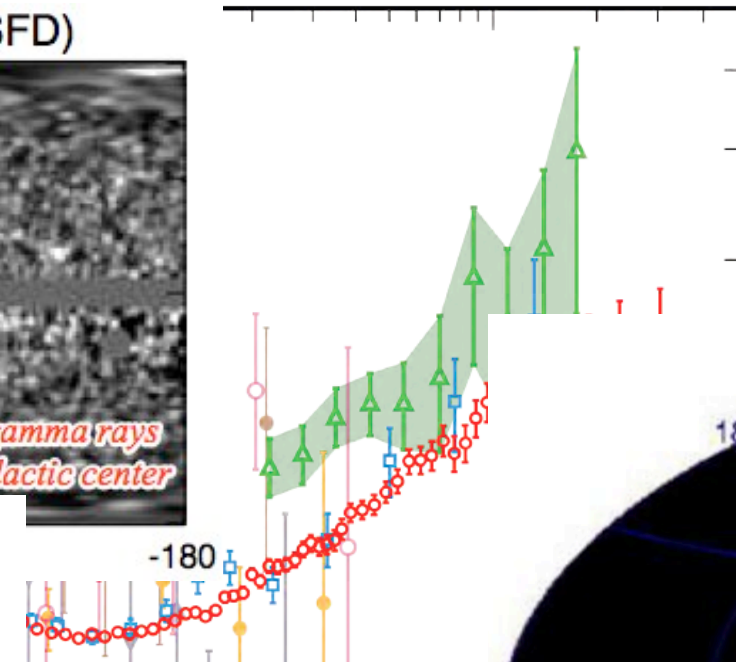
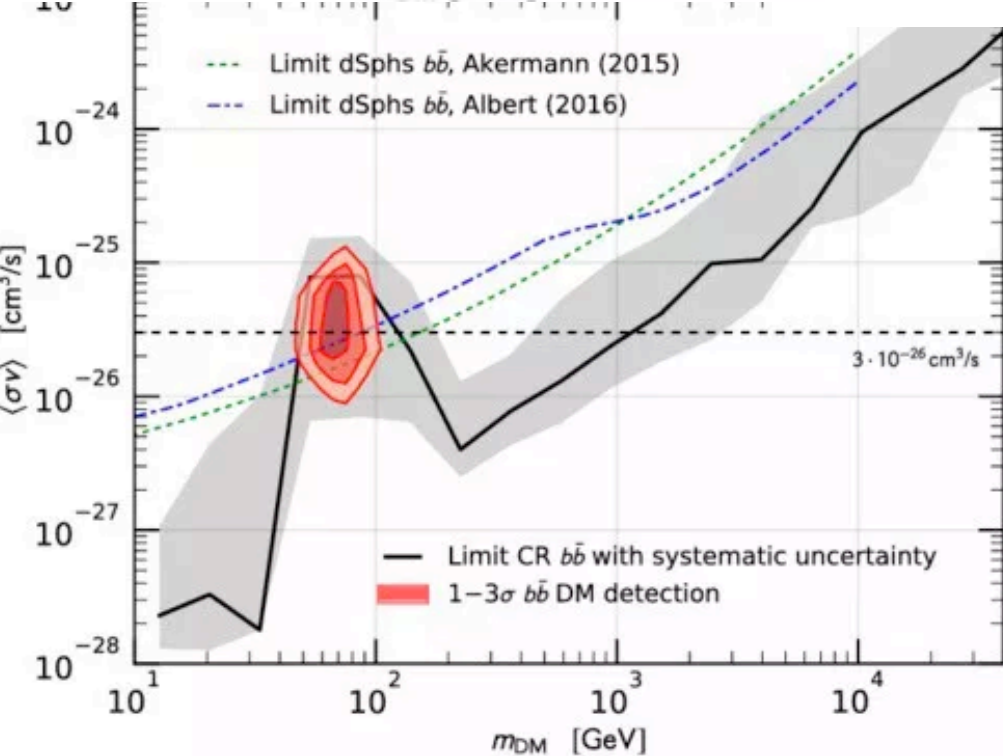
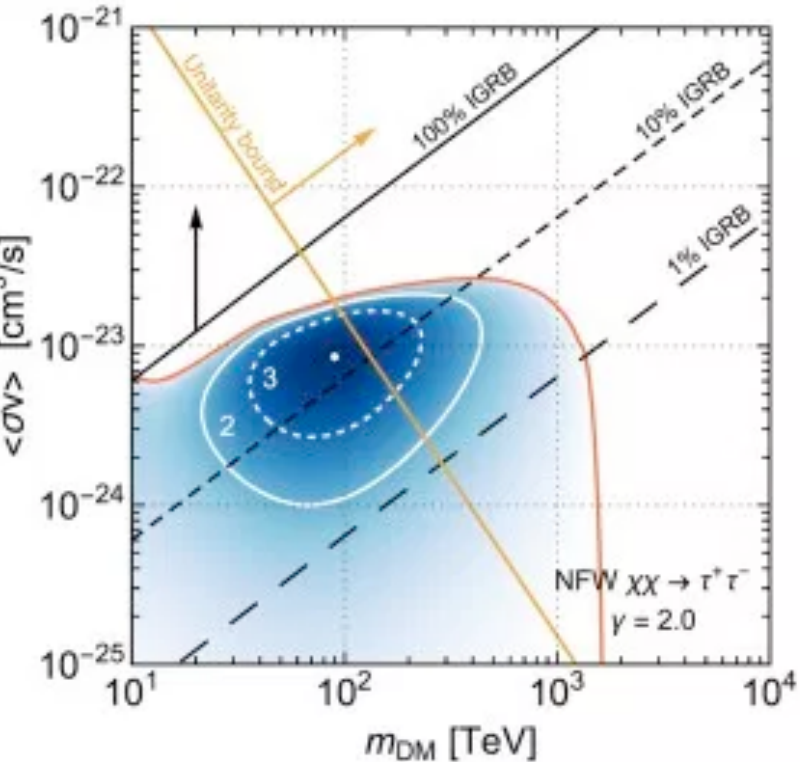
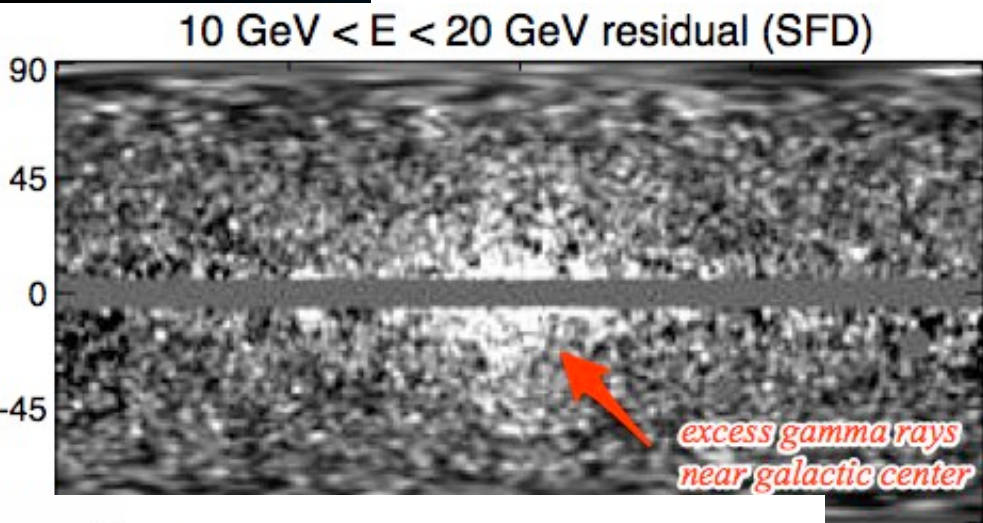


- For a standard dark matter density profile, the annihilation rate within 5° of the Galactic center is $\sim 1 \times 10^{38} \text{ ann s}^{-1}$.
- For a 1 m^2 instrument, this produces a flux of $10^{-4} \text{ ann s}^{-1}$.

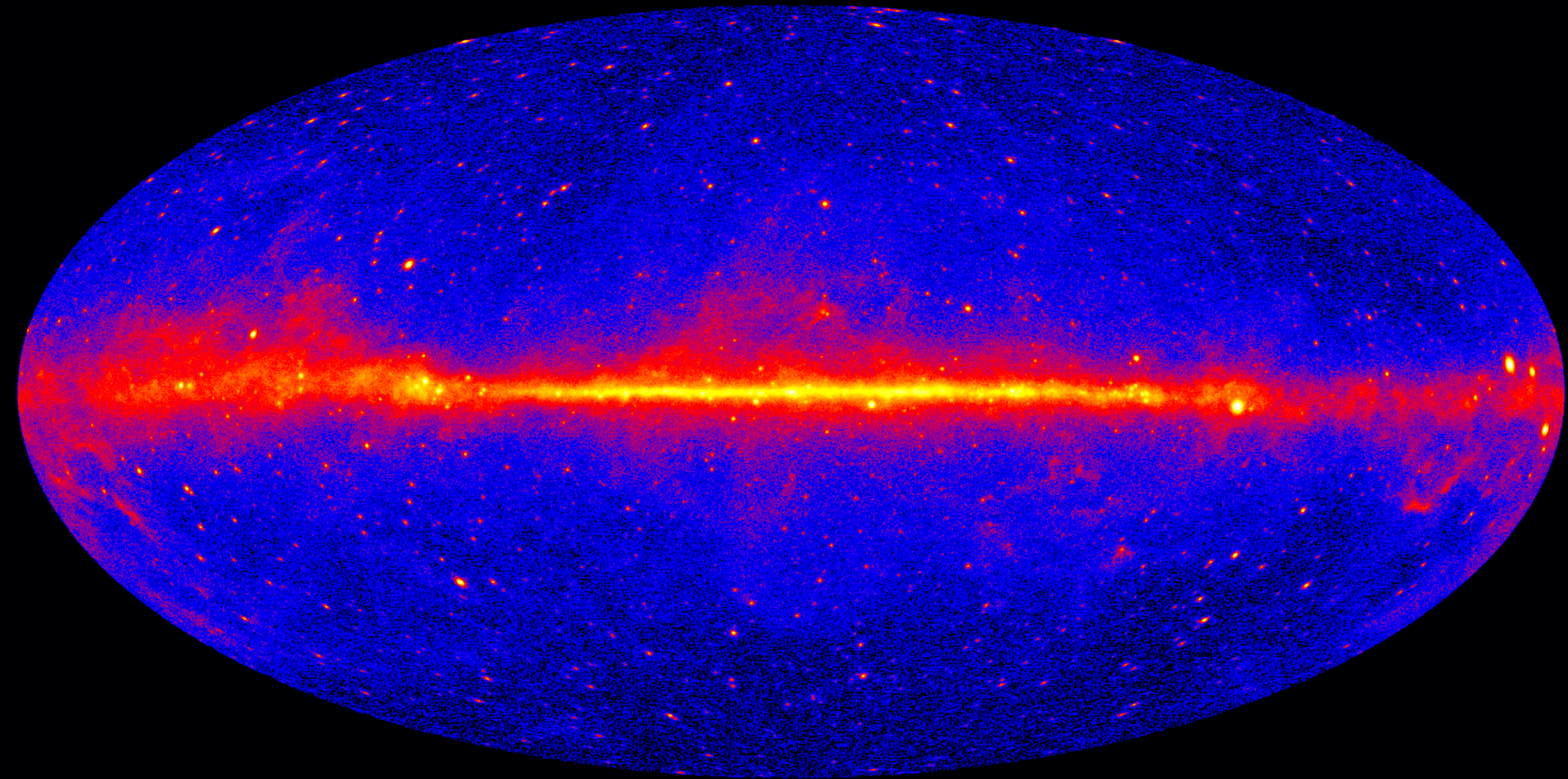




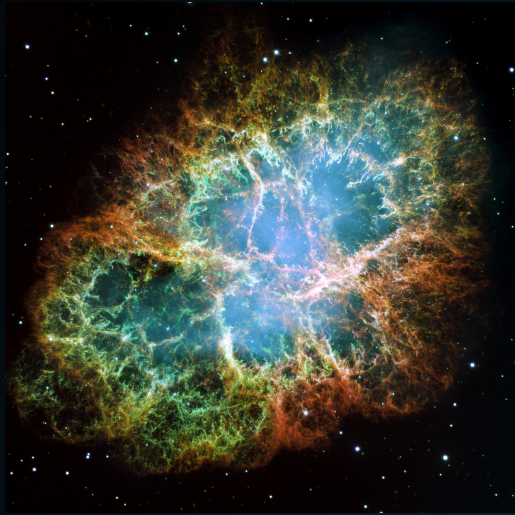




What produces the background?



COSMIC-RAY ACCELERATION AND PROPAGATION



Start with a source of relativistic cosmic-rays

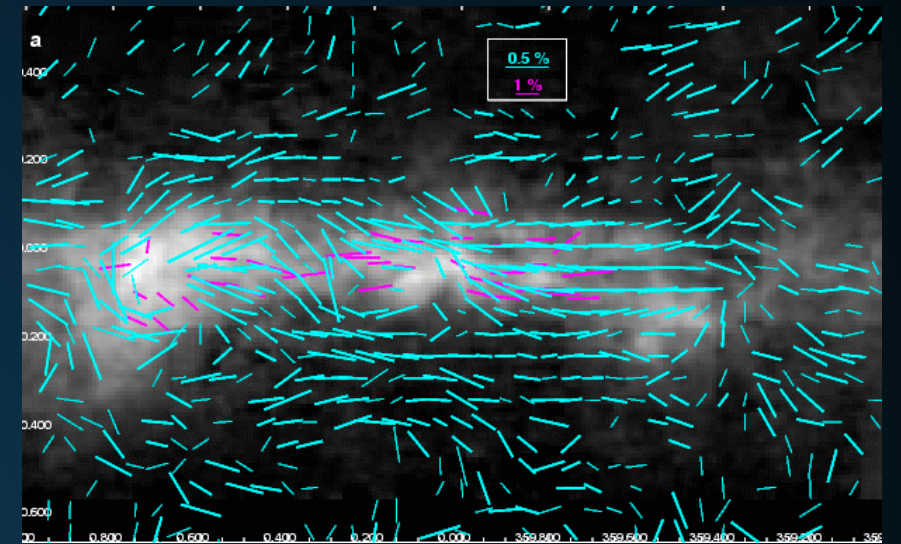
- **Supernova Explosions**
- **Supernova Remnants**
- **Pulsars**
- **Shocks/Mergers**

COSMIC-RAY ACCELERATION AND PROPAGATION



Start with a source of relativistic cosmic-rays

cosmic rays propagate



$$\frac{\partial \psi}{\partial t} = q(\vec{r}, p) + \vec{\nabla} \cdot (D_{xx} \vec{\nabla} \psi - \vec{V} \psi) + \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial}{\partial p} \frac{1}{p^2} \psi - \frac{\partial}{\partial p} \left[p \psi - \frac{p}{3} (\vec{\nabla} \cdot \vec{V}) \psi \right] - \frac{1}{\tau_f} \psi - \frac{1}{\tau_r} \psi$$

Solved Numerically:
e.g. Galprop

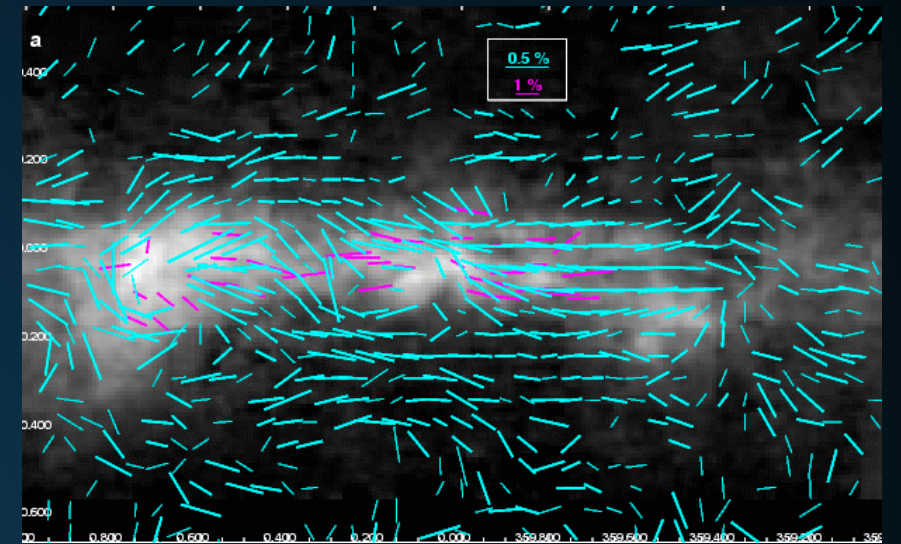
- If they propagate to Earth, can be detected:
 - AMS-02/PAMELA
 - CREAM/HEAT/CAPRICE

COSMIC-RAY ACCELERATION AND PROPAGATION



Start with a source of relativistic cosmic-rays

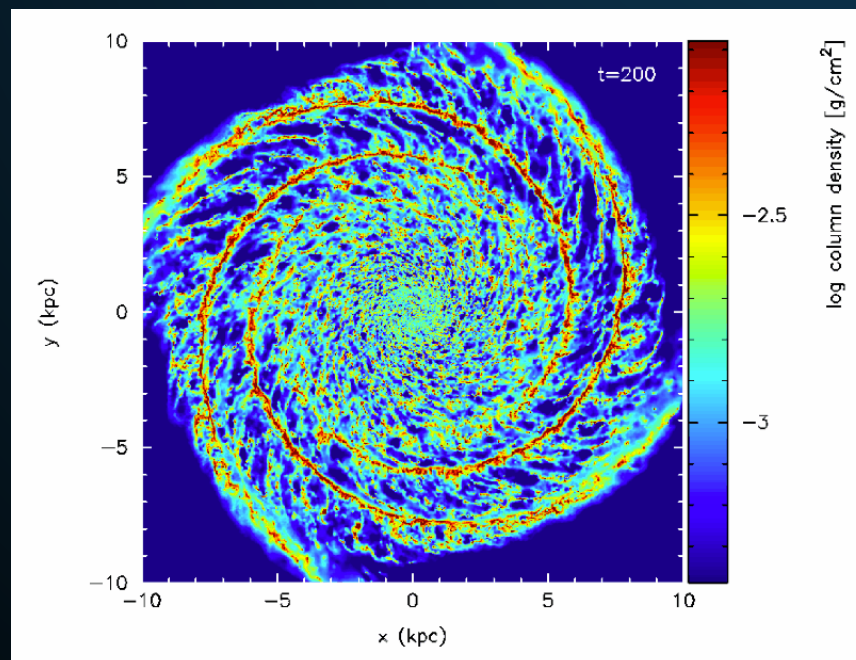
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Solved Numerically:
e.g. Galprop

Gas/ISRF



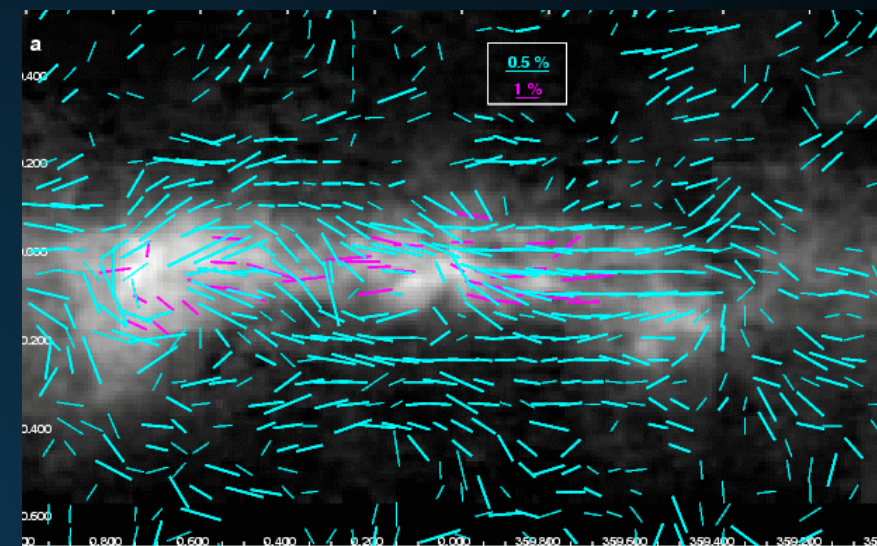
- Alternatively can collide with Galactic gas or the interstellar radiation field.

COSMIC-RAY ACCELERATION AND PROPAGATION



Start with a source of relativistic cosmic-rays

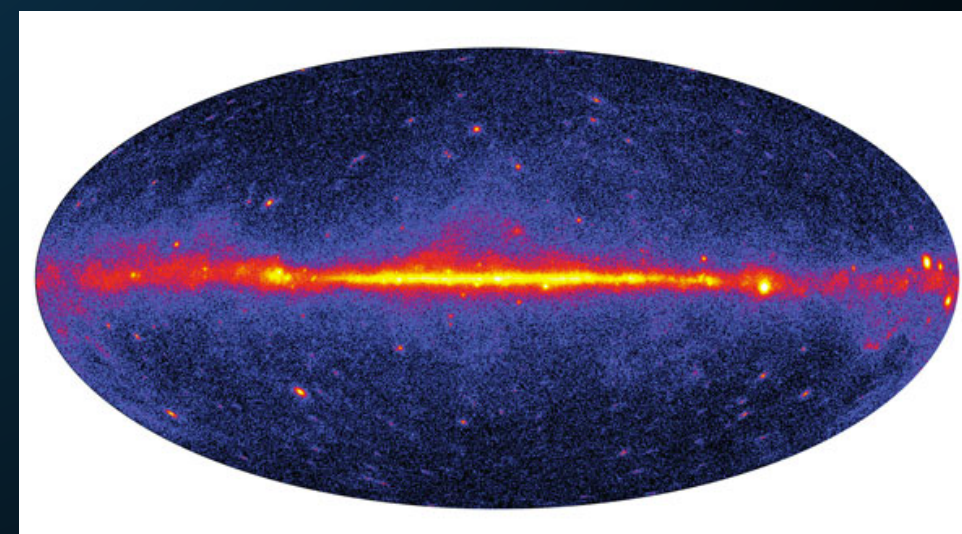
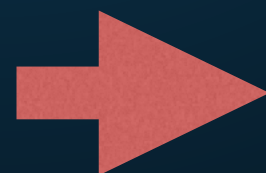
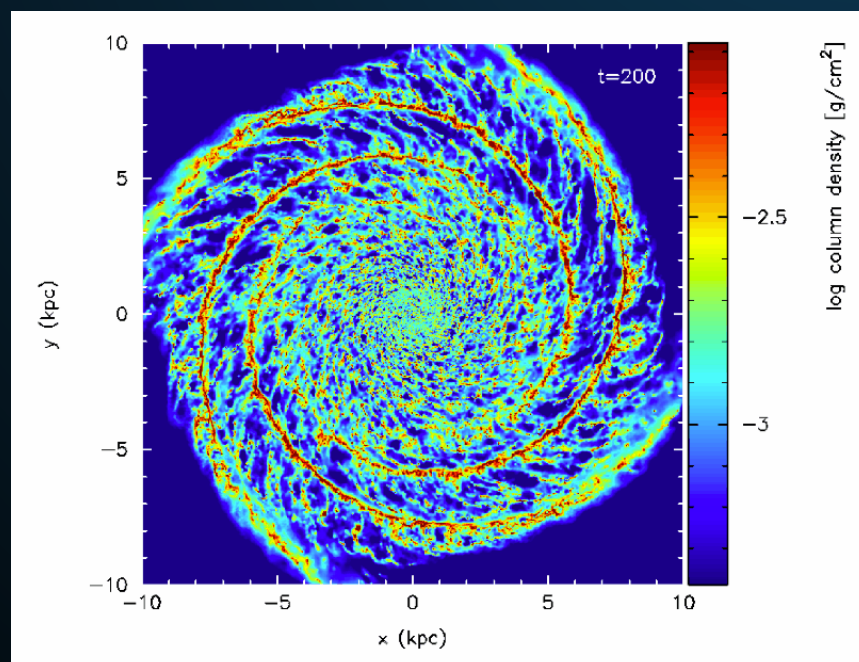
cosmic rays propagate



$$\frac{\partial \psi}{\partial t} = q(\vec{r}, p) + \vec{\nabla} \cdot (D_{xx} \vec{\nabla} \psi - \vec{V} \psi) + \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial}{\partial p} \frac{1}{p^2} \psi - \frac{\partial}{\partial p} \left[p \psi - \frac{p}{3} (\vec{\nabla} \cdot \vec{V}) \psi \right] - \frac{1}{\tau_f} \psi - \frac{1}{\tau_r} \psi$$

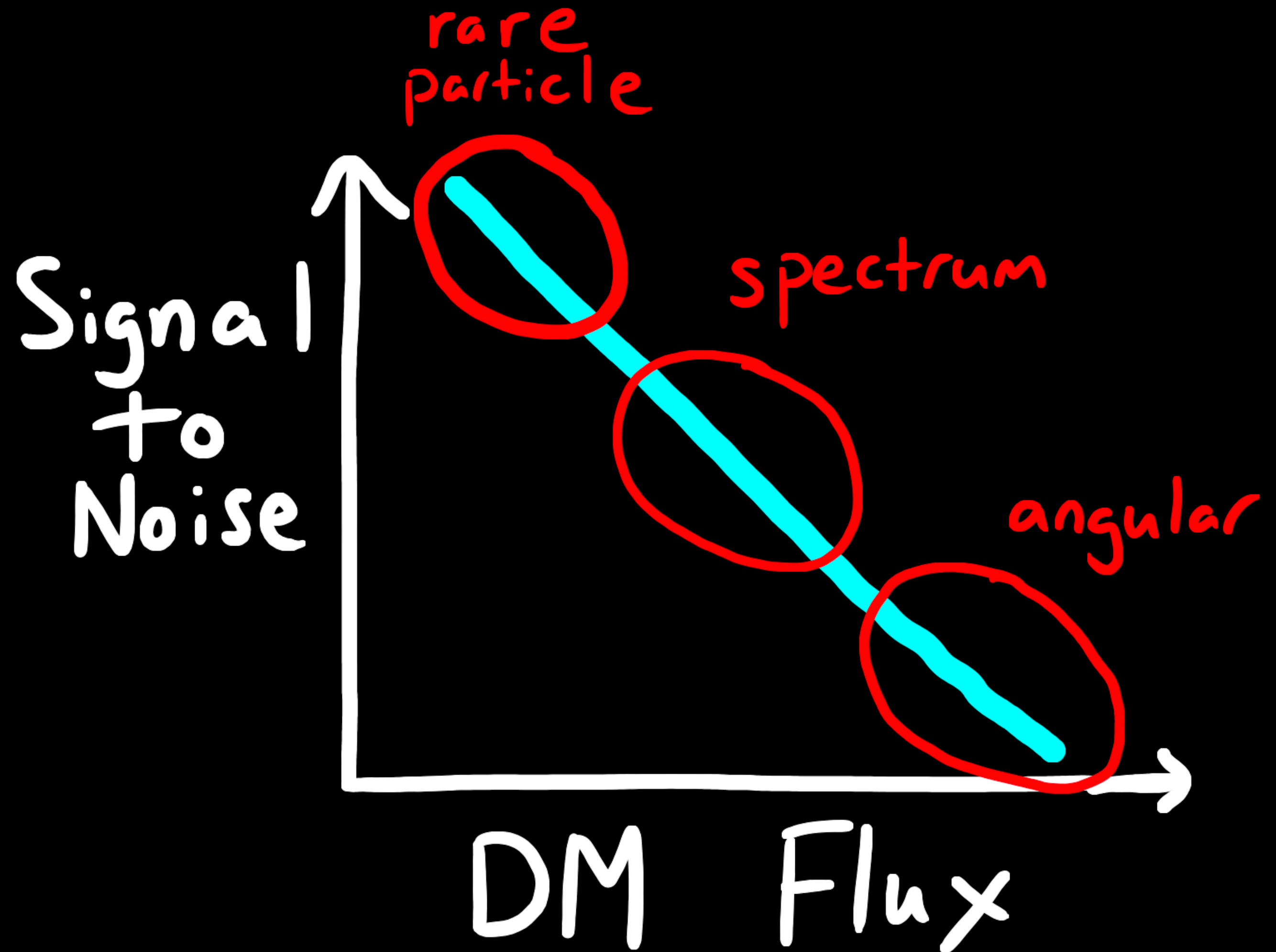
Solved Numerically:
e.g. Galprop

Gas/ISRF



The Techniques

- **Three Methods to Separate Dark Matter Signals:**
 - **Rare Particle Detection**
 - **Spectral Features**
 - **Angular Mapping**

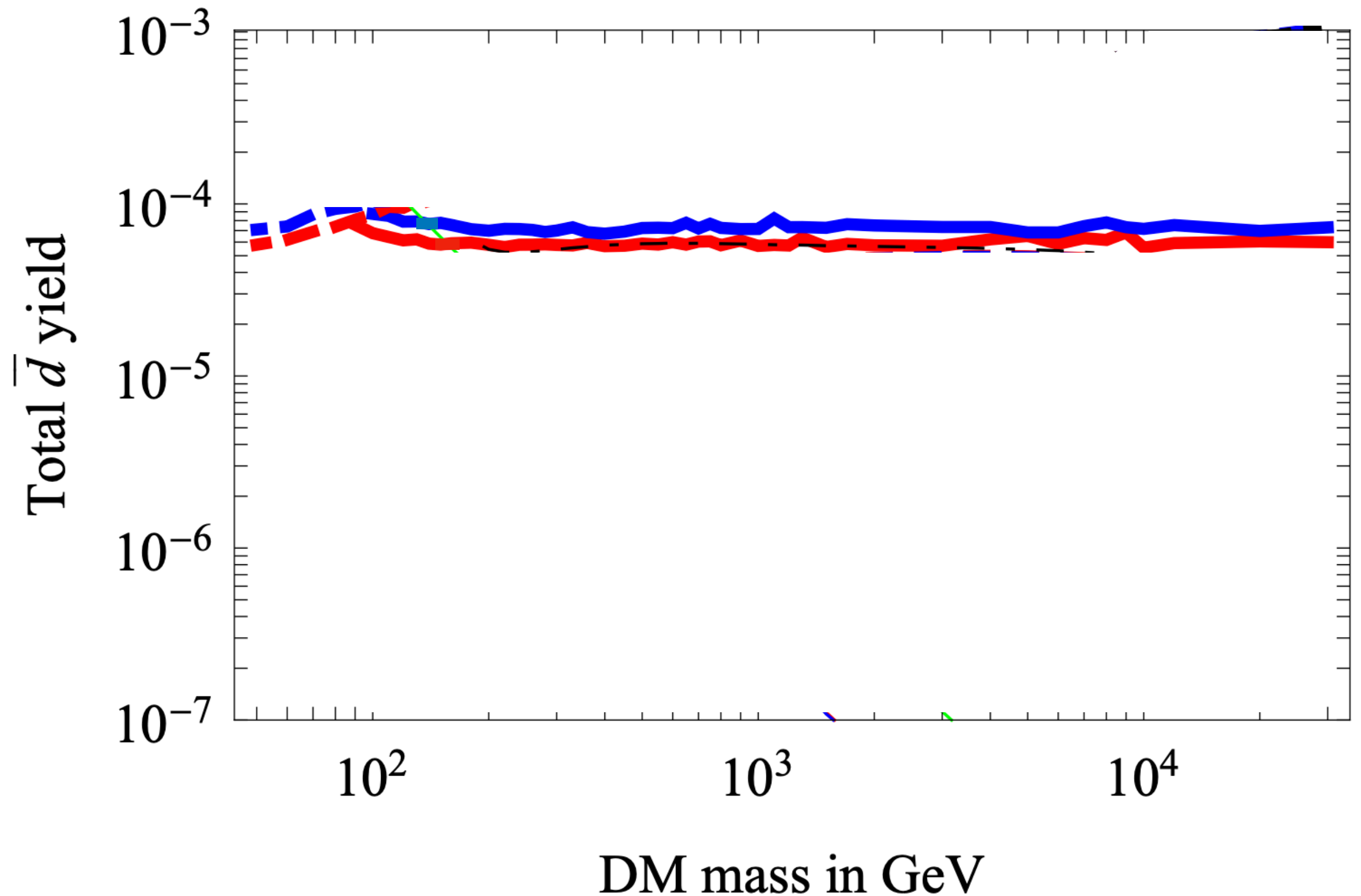


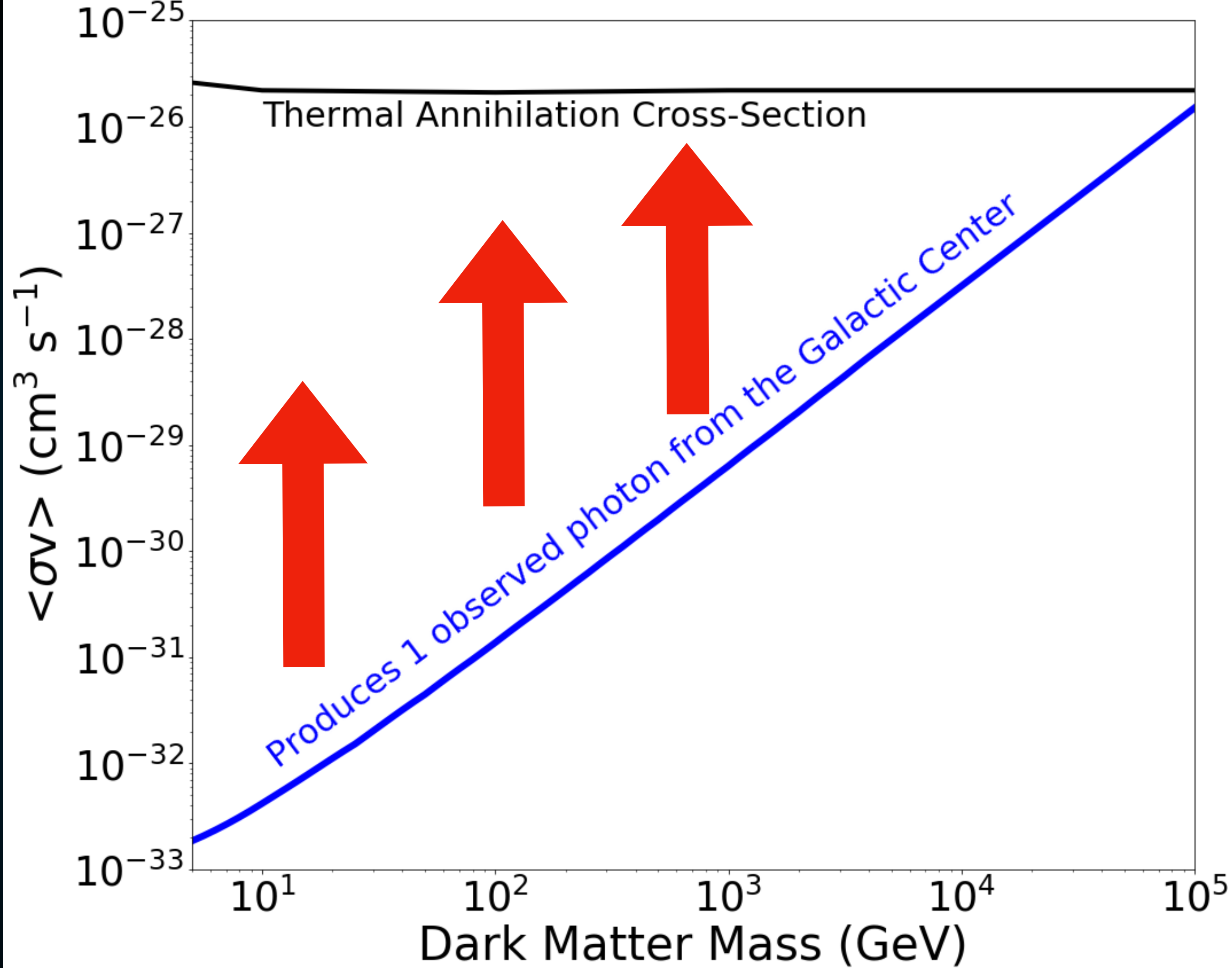
Rare Particle Detection

Exploiting the fact that the universe is mostly matter



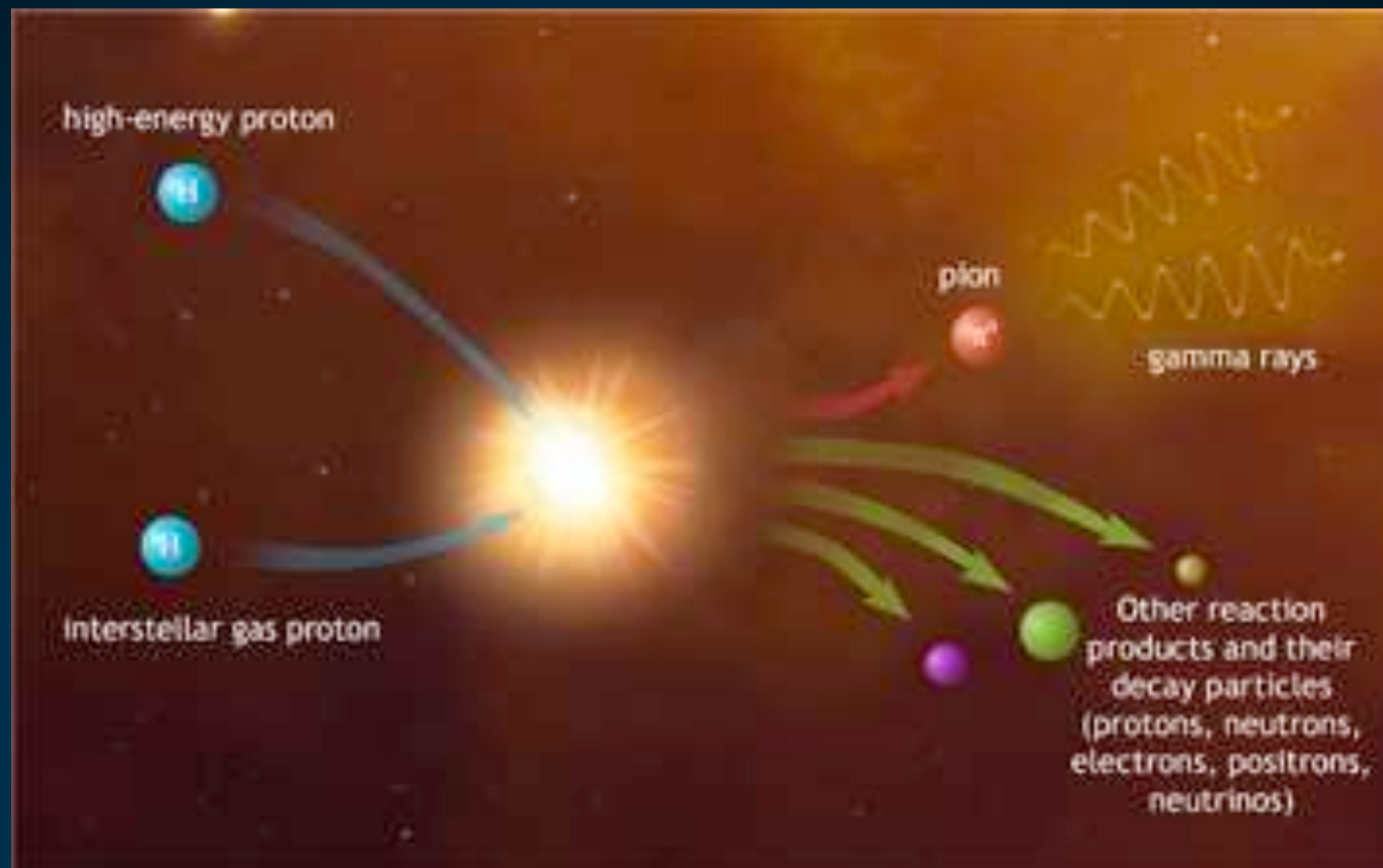
Exploiting the fact that the universe is mostly matter





Rare Particle Detection

Exploiting the fact that the universe is mostly matter



$$p + p \longrightarrow p + p + p + \bar{p} + n + \bar{n}$$

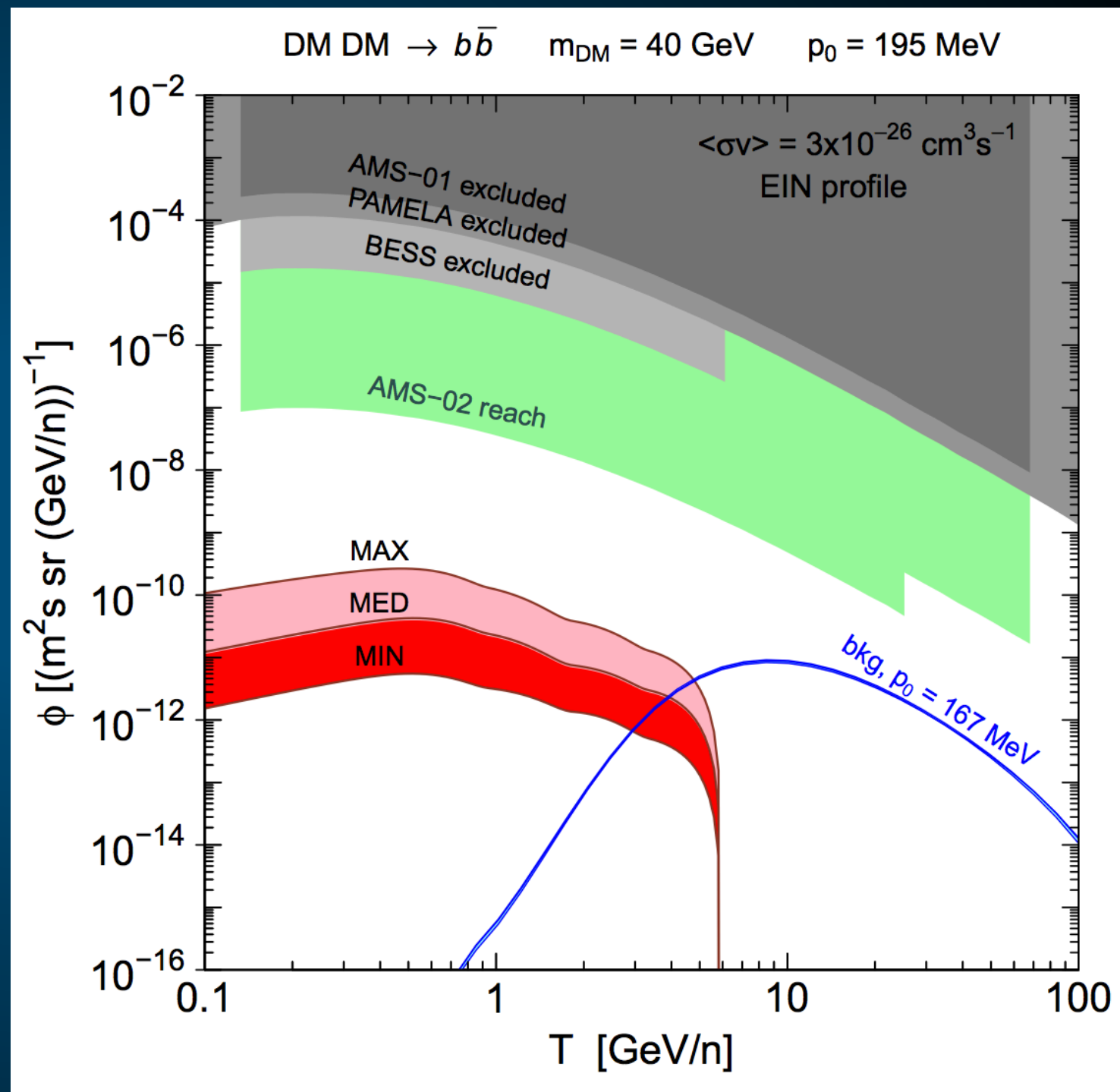
- **Kinematic threshold makes background negligible below 10 GeV.**

Rare Particle Detection

Exploiting the fact that the universe is mostly matter

$$\chi + \chi \longrightarrow p + \bar{p} + n + \bar{n}$$

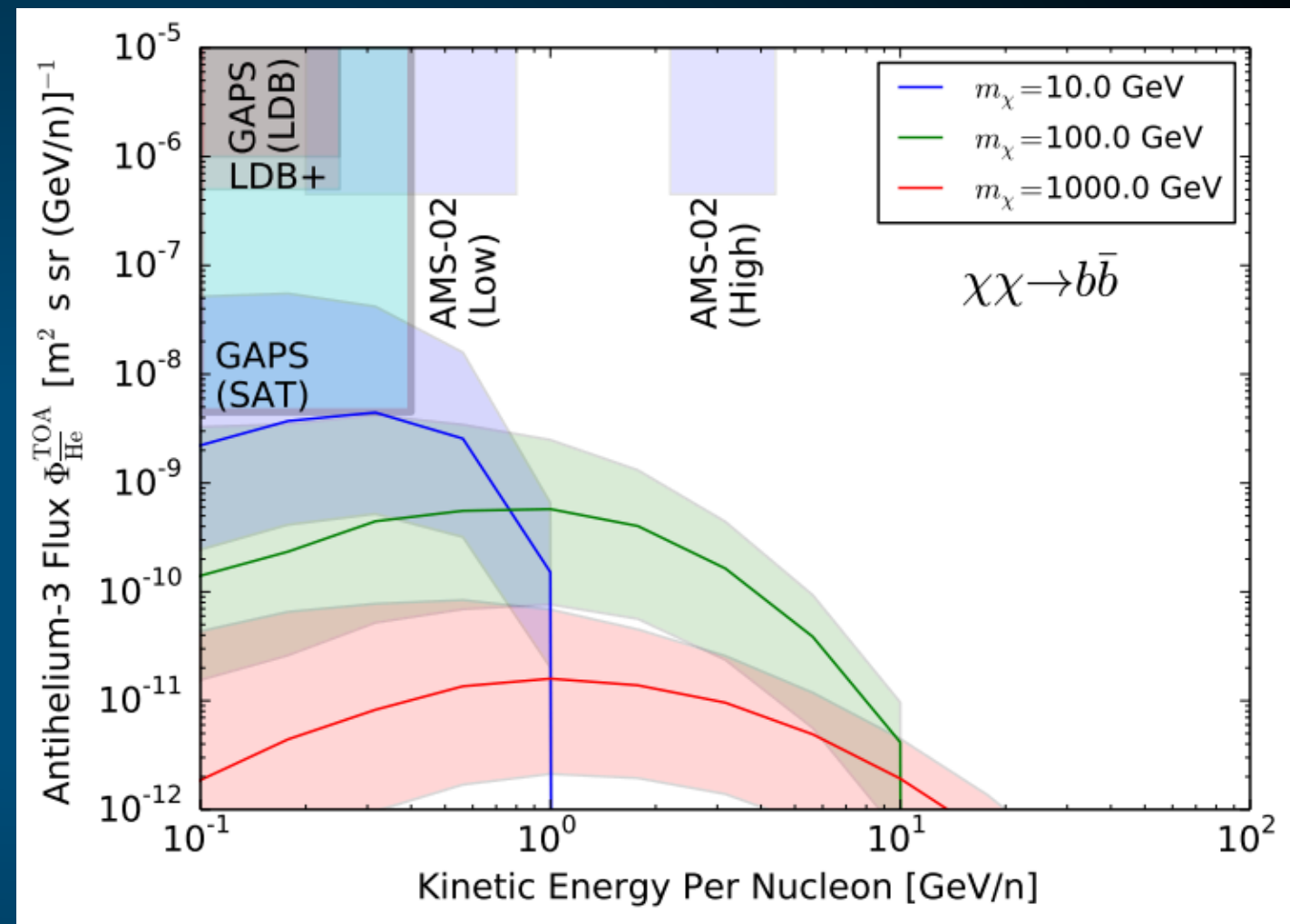
- Dark matter annihilation occurs in the lab frame.
- Dark matter signal dominate at low energies.
- Energies can't change due to propagation!



Rare Particle Detection

Exploiting the fact that the universe is mostly matter

- Dark matter signal is even more dominant in the case of anti-Helium.
- Depending on coupling, anti-Helium signal does not need to be much smaller.



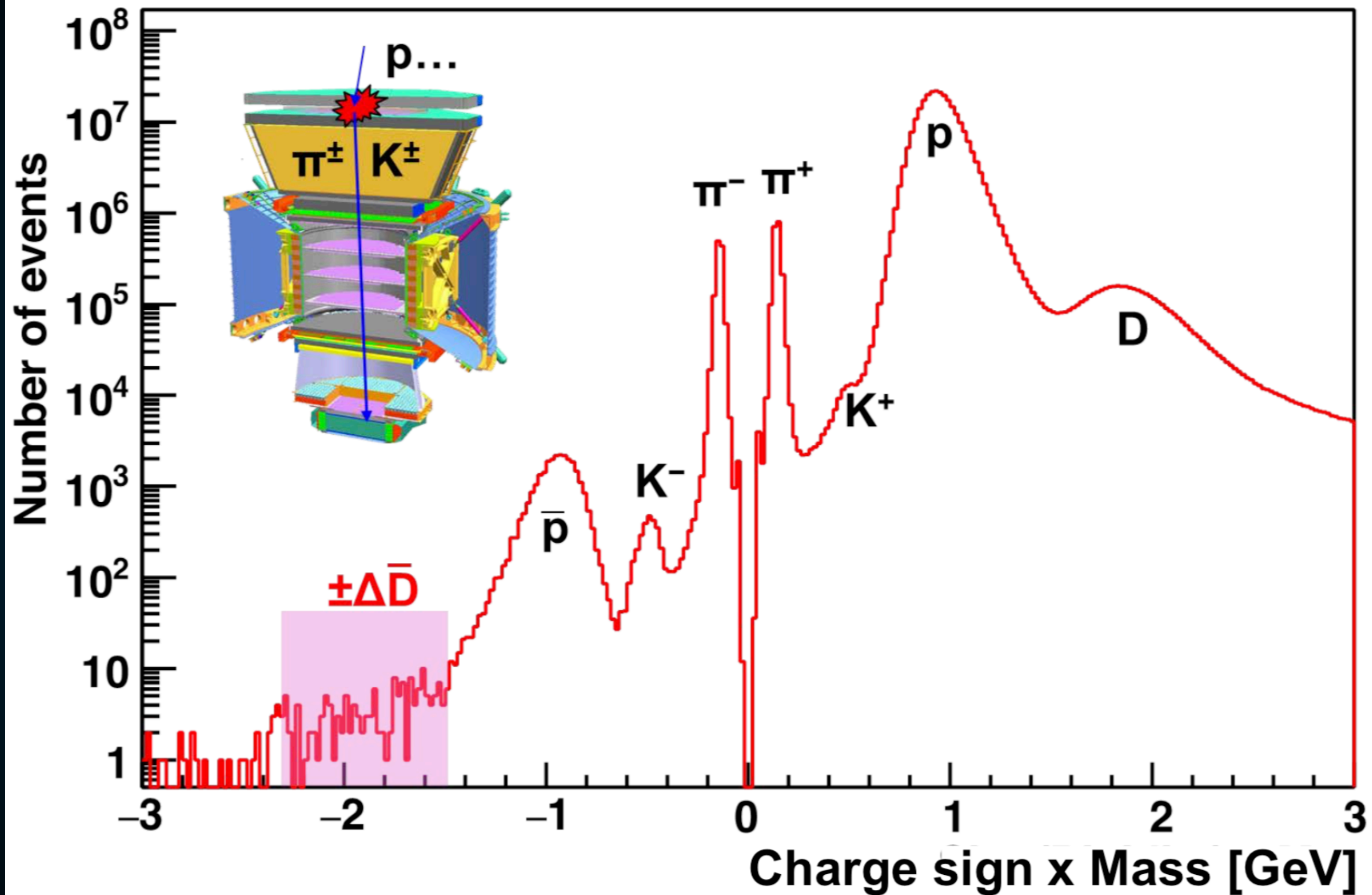
To date, we have observed eight events in the mass region from 0 to 10 GeV with $Z = -2$. All eight events are in the helium mass region.

Currently (having used 50 million core hours to generate 7 times more simulated events than measured events and having found no background events from the simulation), our best evaluation of the probability of the background origin for the eight He events is less than 3×10^{-8} . For the two ${}^4\text{He}$ events our best evaluation of the probability (upon completion of the current 100 million core hours of simulation) will be less than 3×10^{-3} .

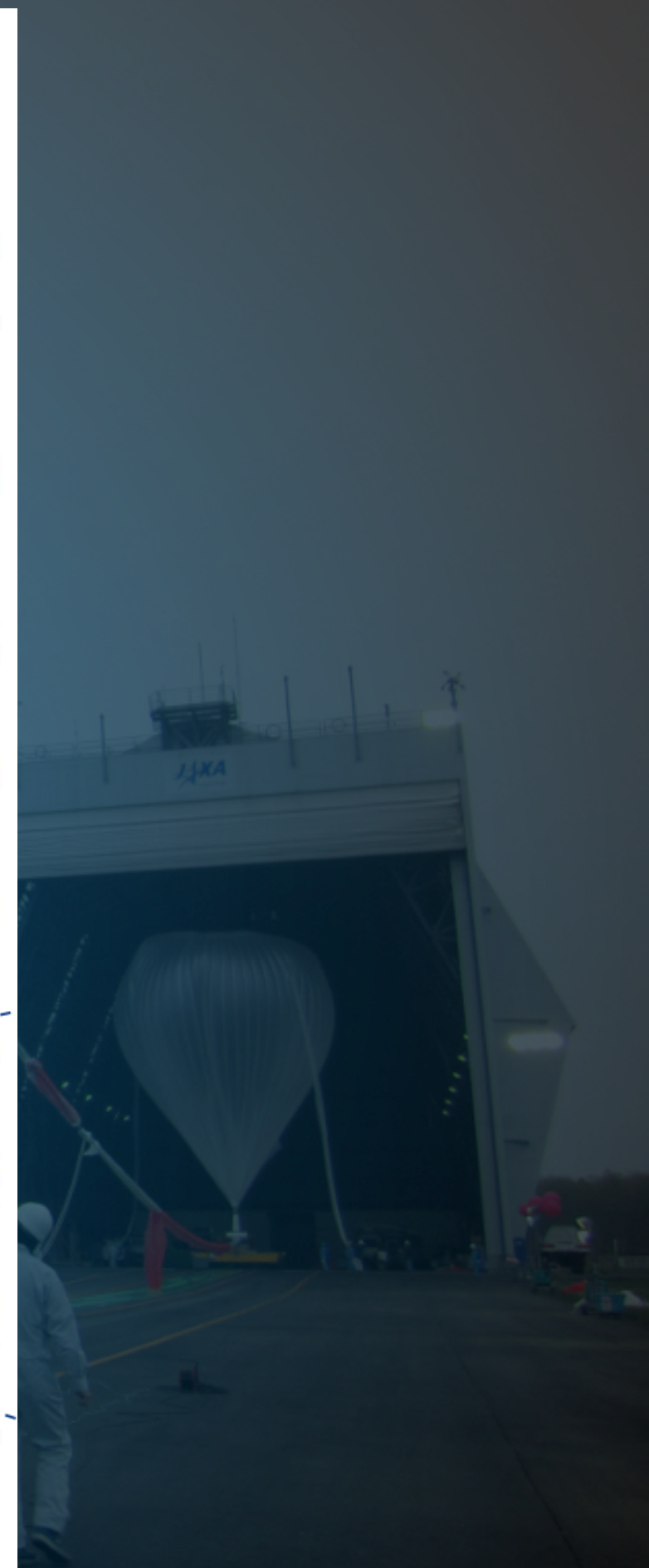
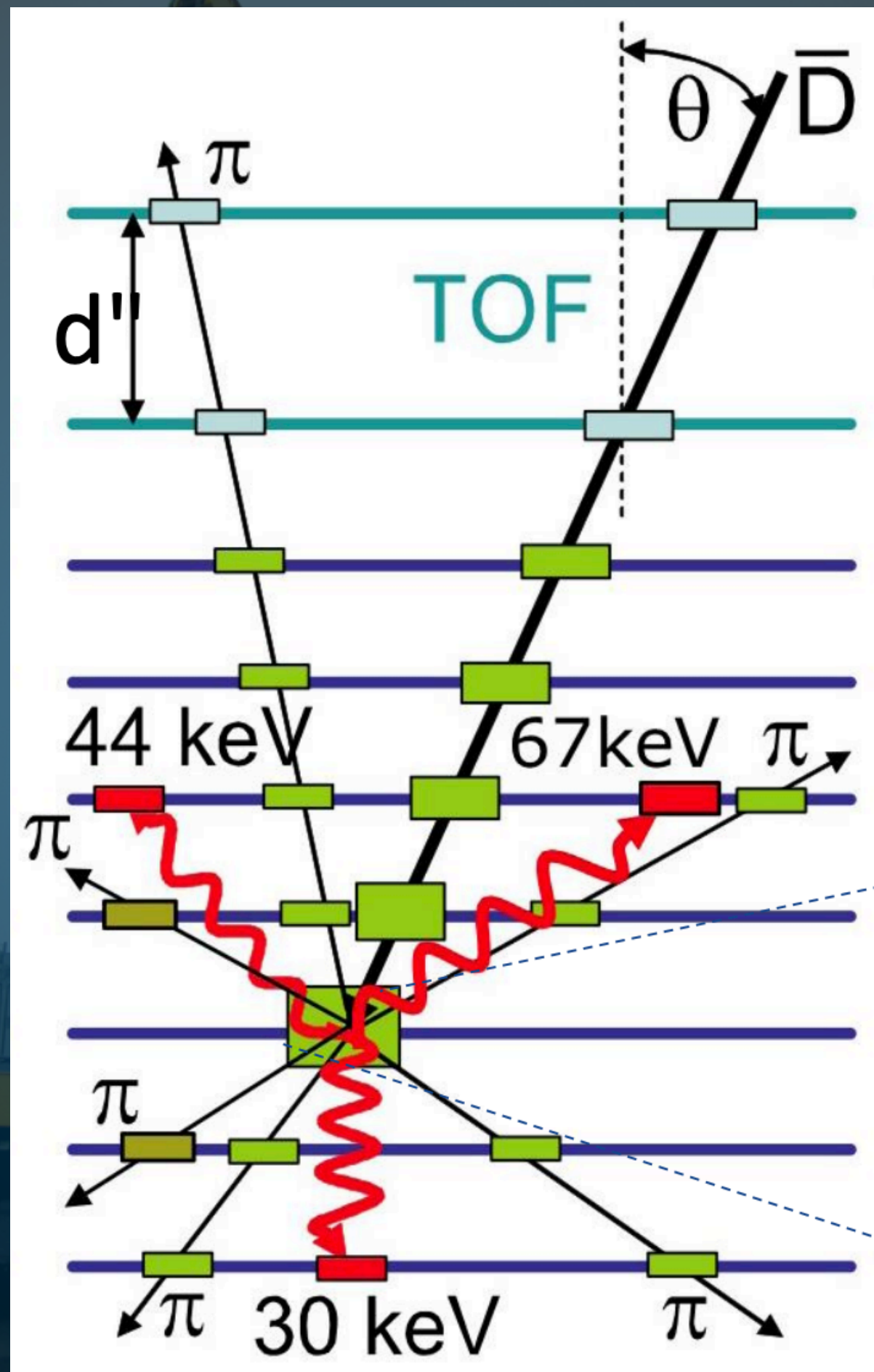
Note that for ${}^4\text{He}$, projecting based on the statistics we have today, by using an additional 400 million core hours for simulation the background probability would be 10^{-4} . Simultaneously, continuing to run until 2023, which doubles the data sample, the background probability for ${}^4\text{He}$ would be 2×10^{-7} , i.e., greater than 5-sigma significance.

Rare Particle Detection

Exploiting the fact that the universe is mostly matter



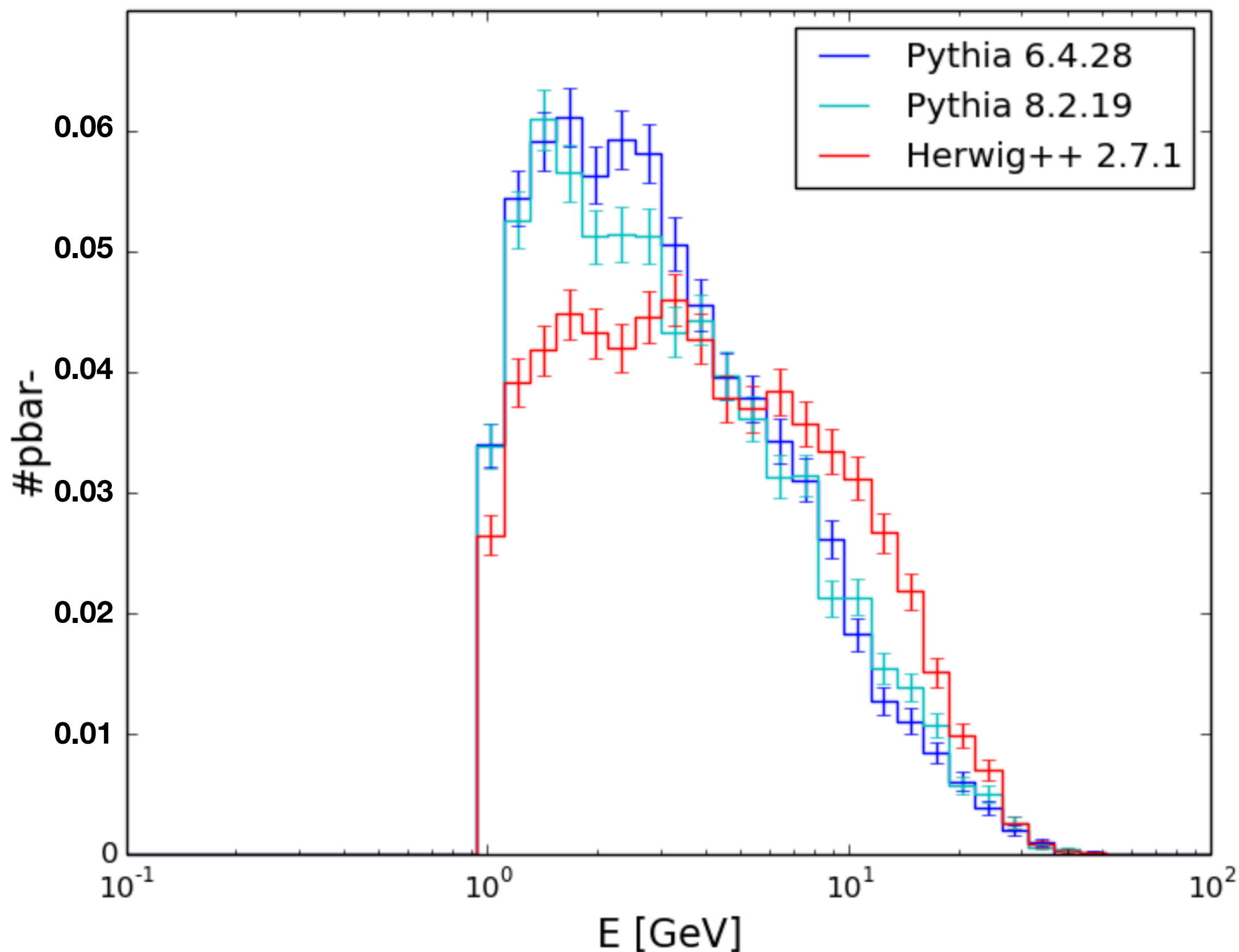




Spectral Features

Joachim Kip (2017)

Things that go bump in the night



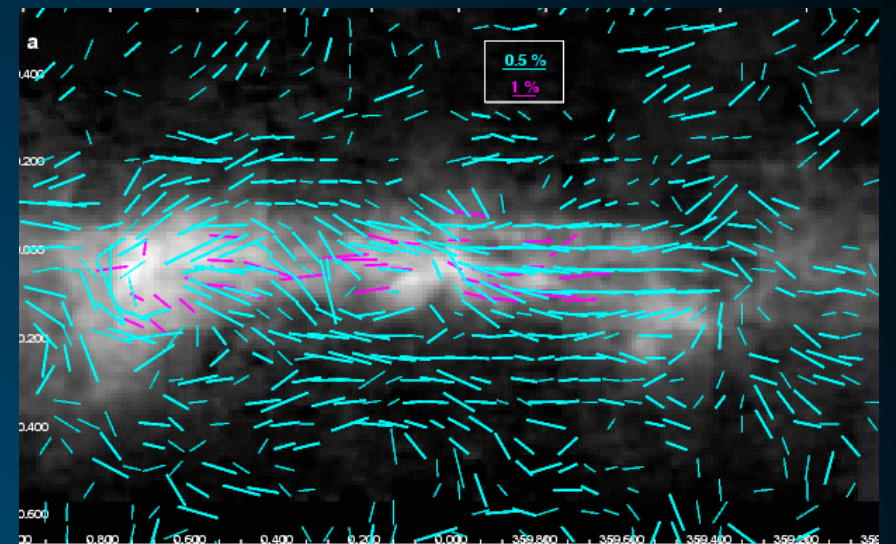
Spectral Features

Things that go bump in the night



Start with a source of relativistic cosmic-rays

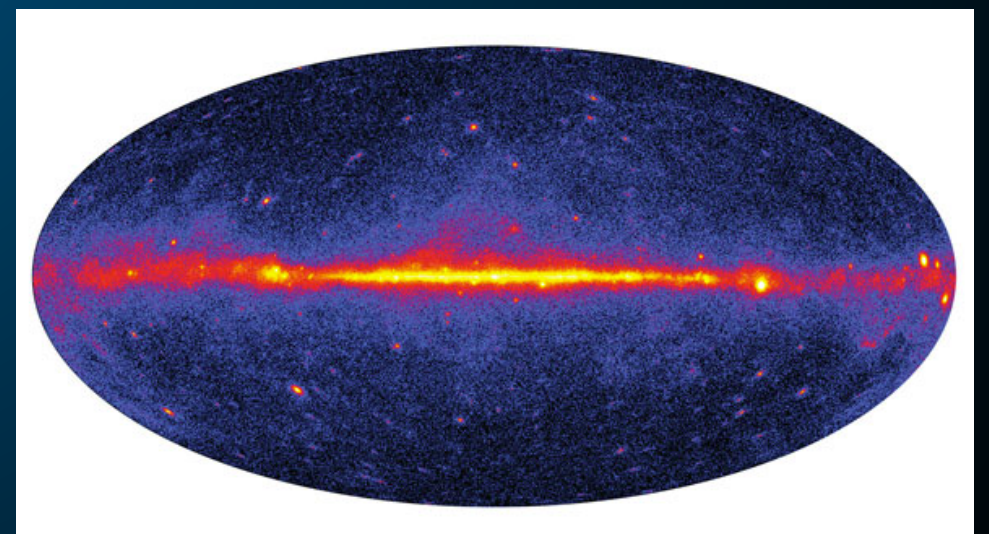
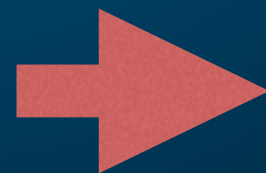
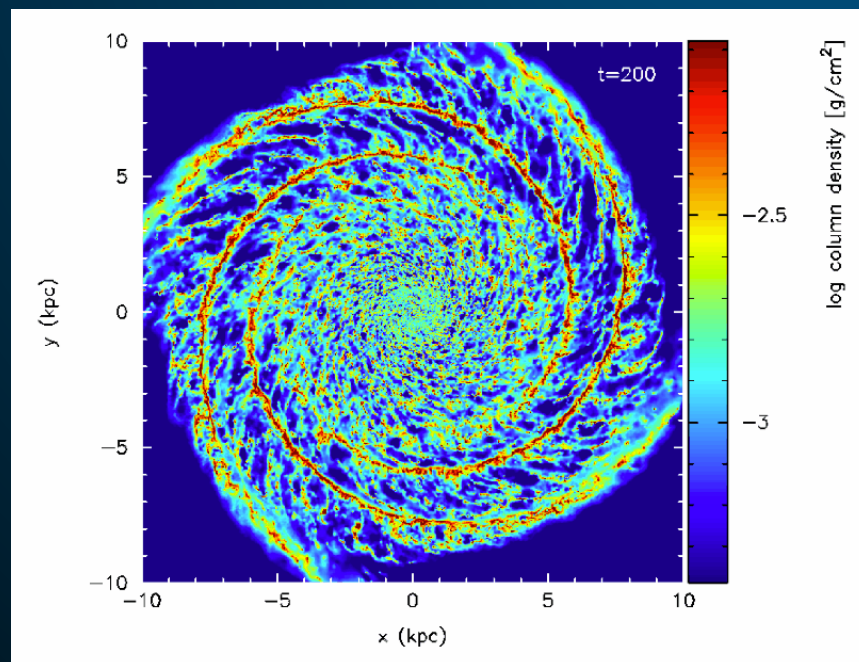
cosmic rays propagate



$$\frac{\partial \psi}{\partial t} = q(\vec{r}, p) + \vec{\nabla} \cdot (D_{xx} \vec{\nabla} \psi - \vec{V} \psi) + \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial}{\partial p} \frac{1}{p^2} \psi - \frac{\partial}{\partial p} \left[p \psi - \frac{p}{3} (\vec{\nabla} \cdot \vec{V}) \psi \right] - \frac{1}{\tau_f} \psi - \frac{1}{\tau_r} \psi$$

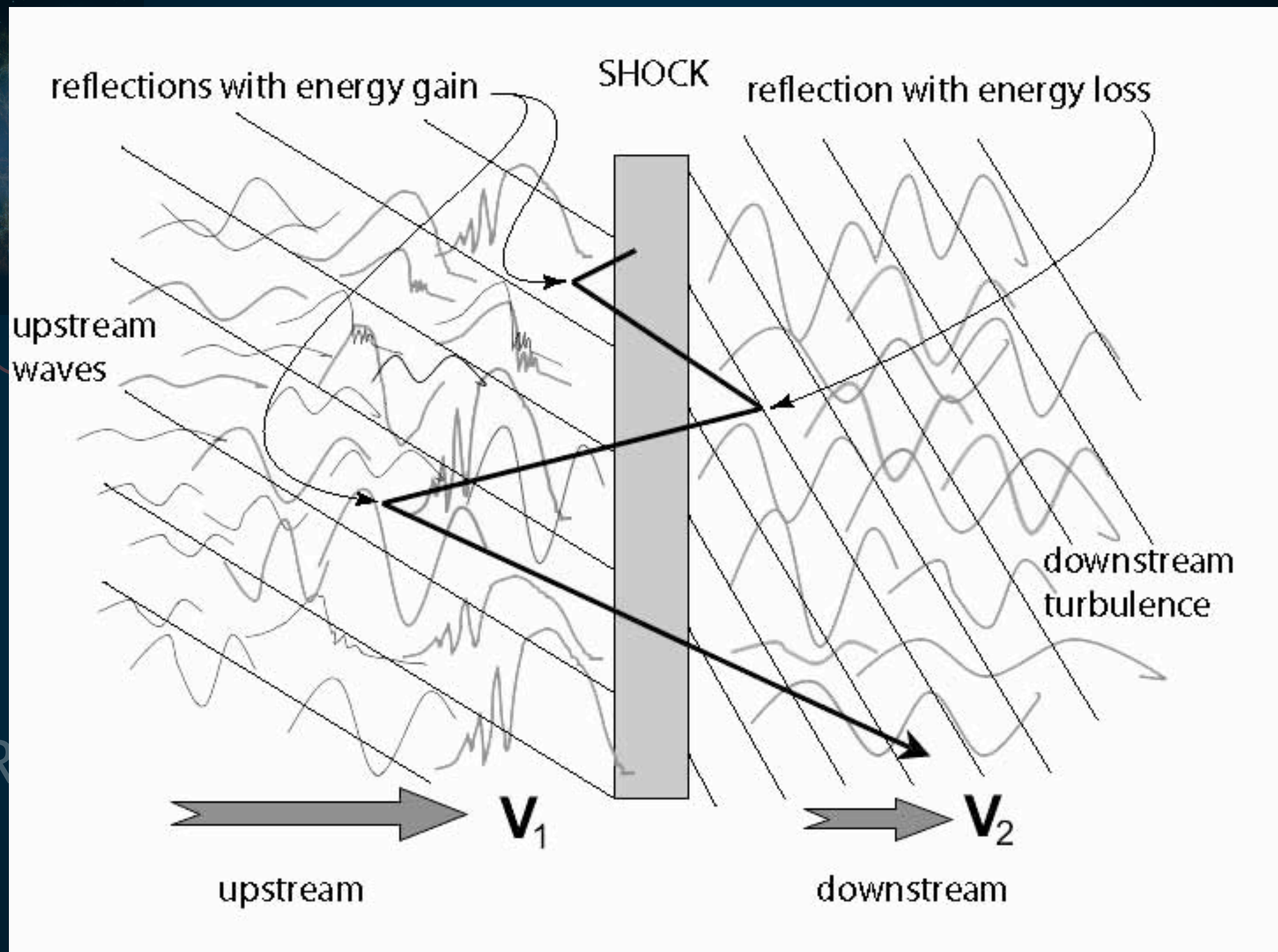
**Solved Numerically:
e.g. Galprop**

Gas/ISRF

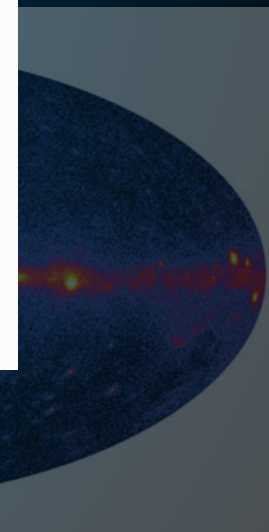


Spectral Features

Things that go bump in the night

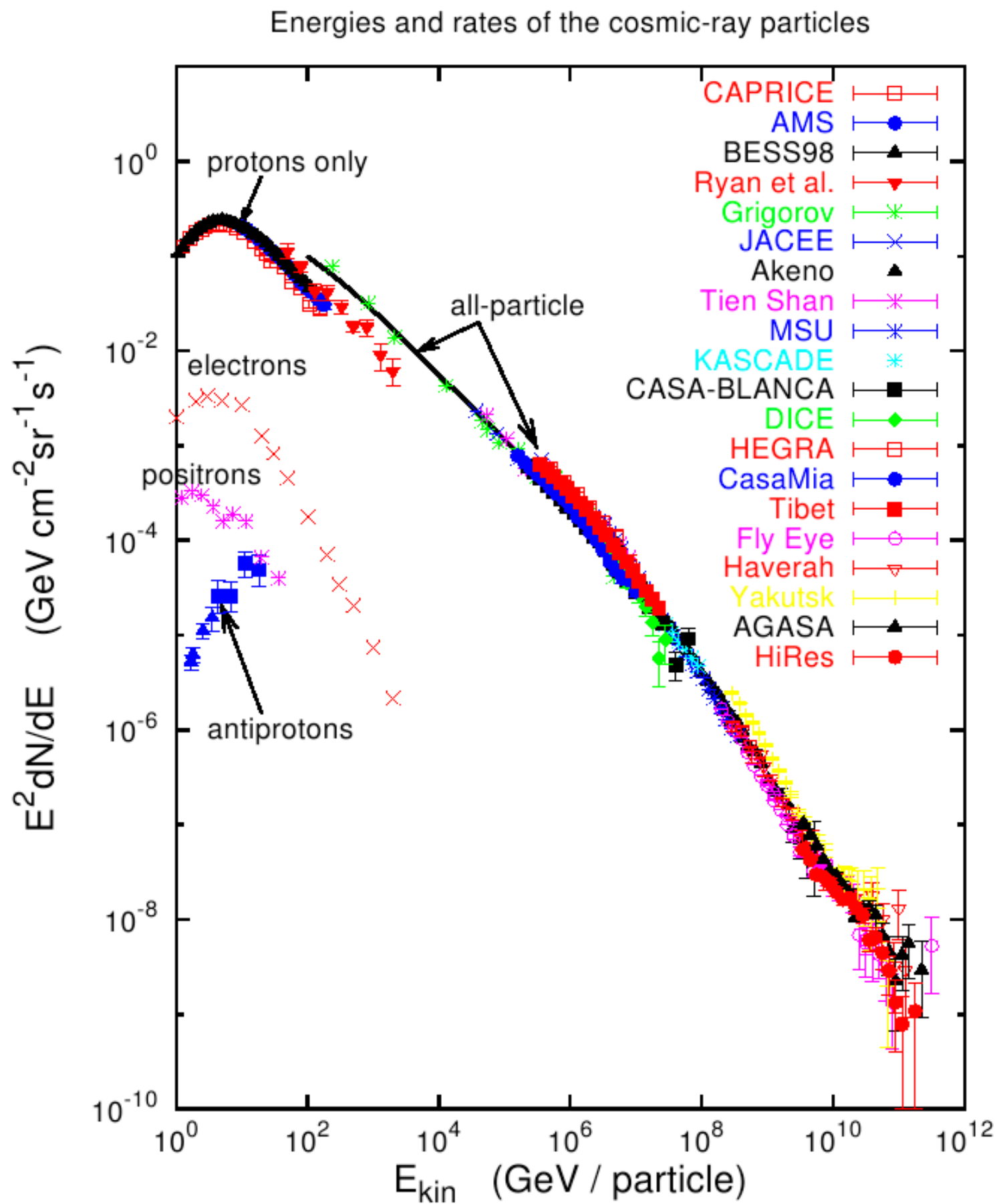


$$\left[b \right] - \frac{1}{\tau_f} \psi - \frac{1}{\tau_r} \psi$$

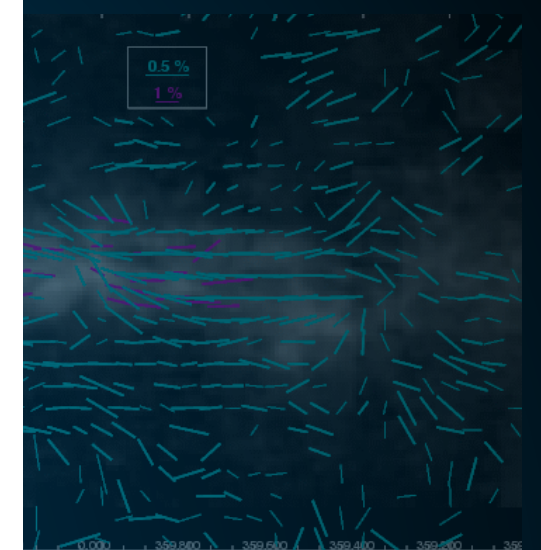


Spectral Features

Things that

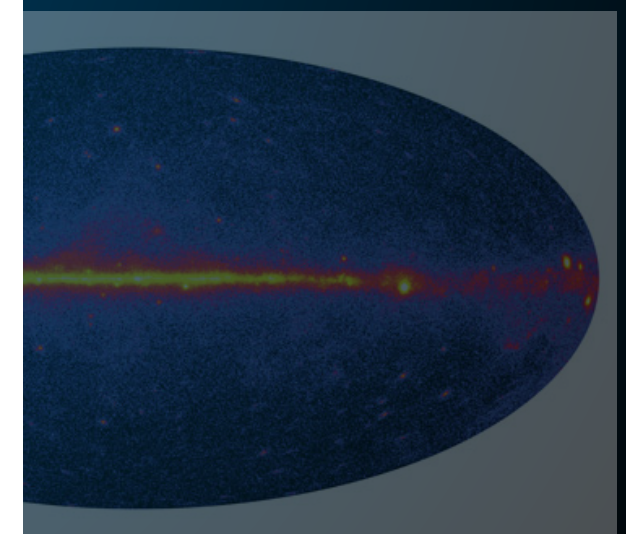


mic-rays



$$-\frac{\partial}{\partial p} \left[\dot{p} \psi - \frac{p}{3} (\vec{\nabla} \cdot \vec{V}) \psi \right] - \frac{1}{\tau_f} \psi - \frac{1}{\tau_r} \psi$$

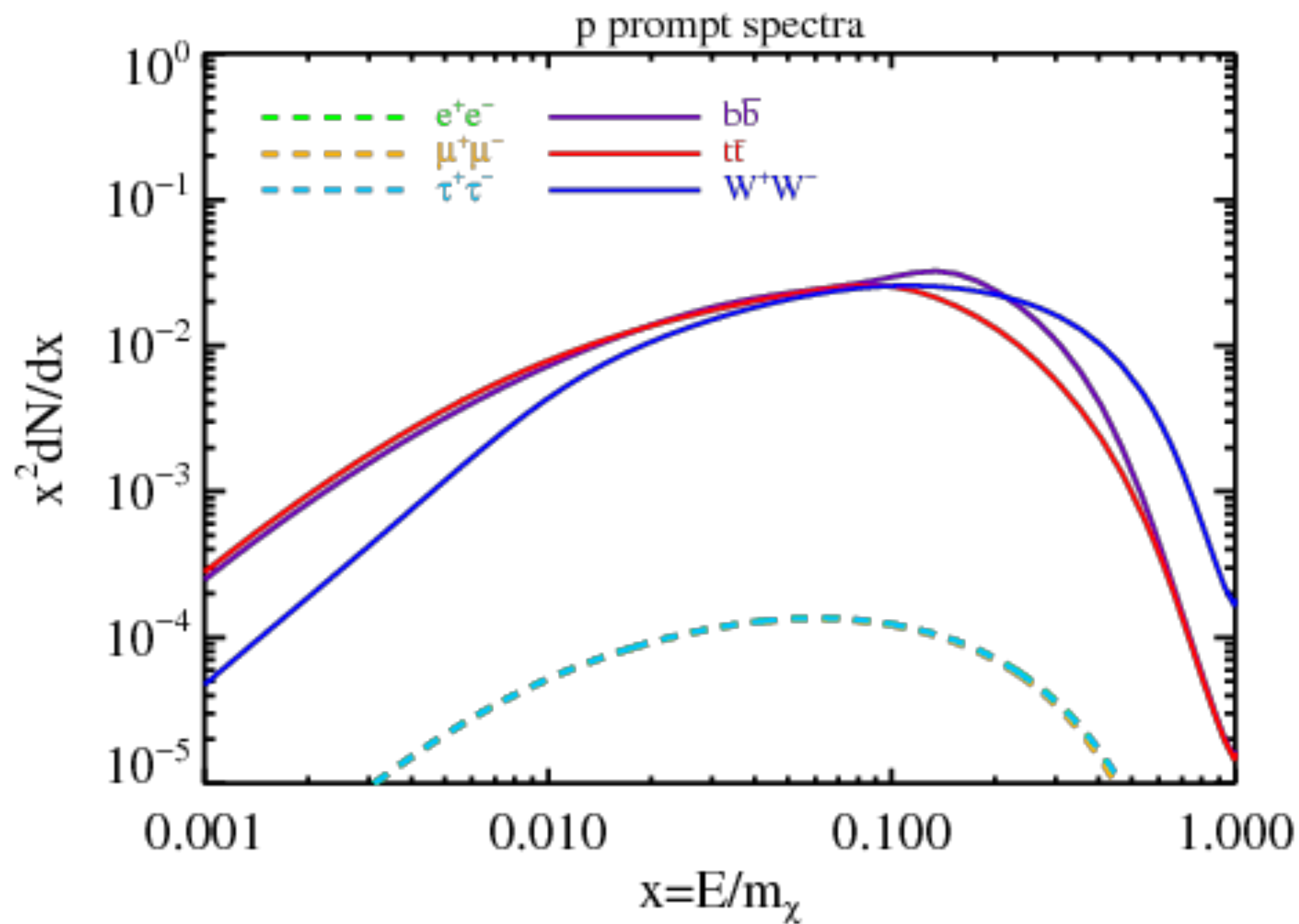
y:



Gas/ISRF

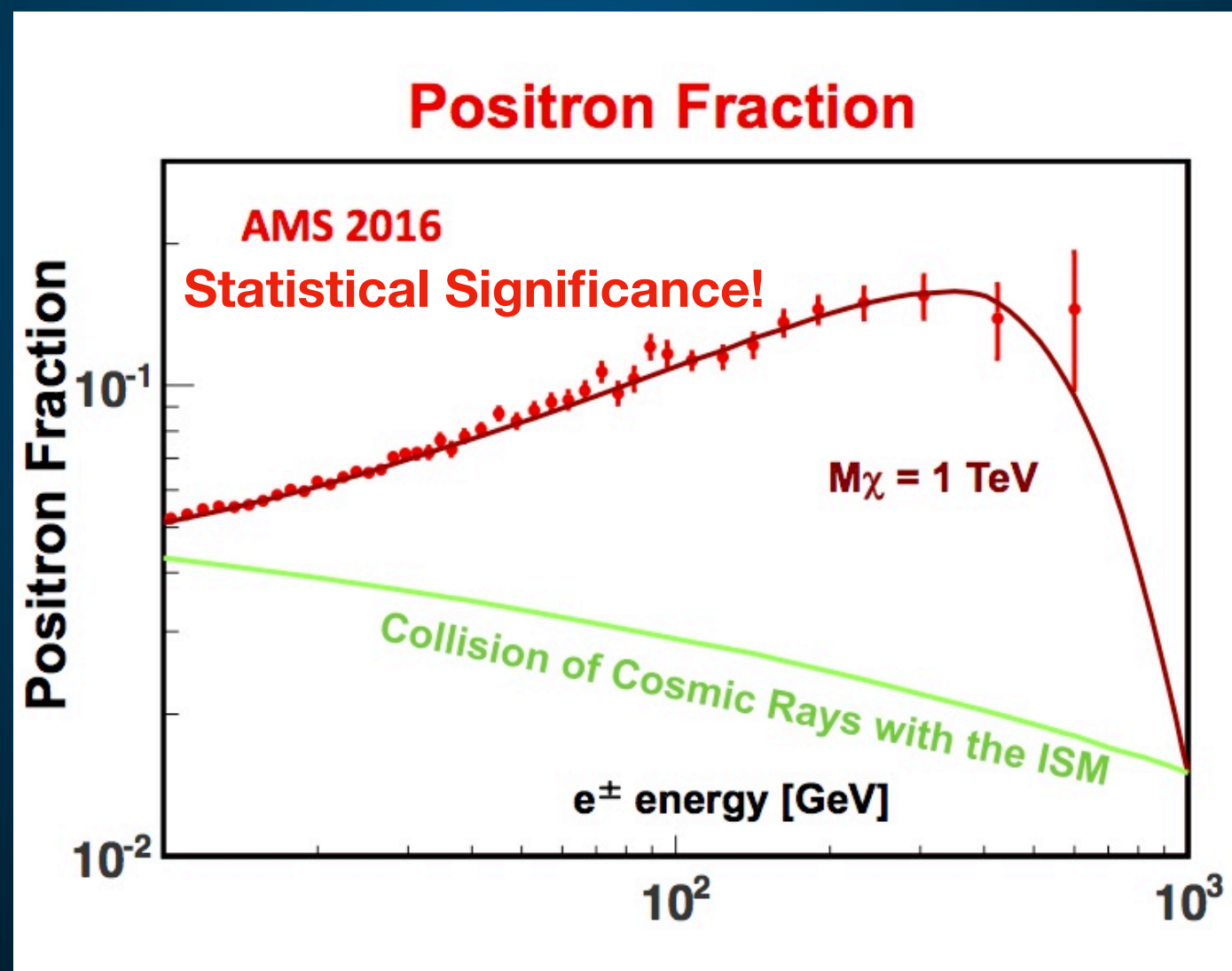
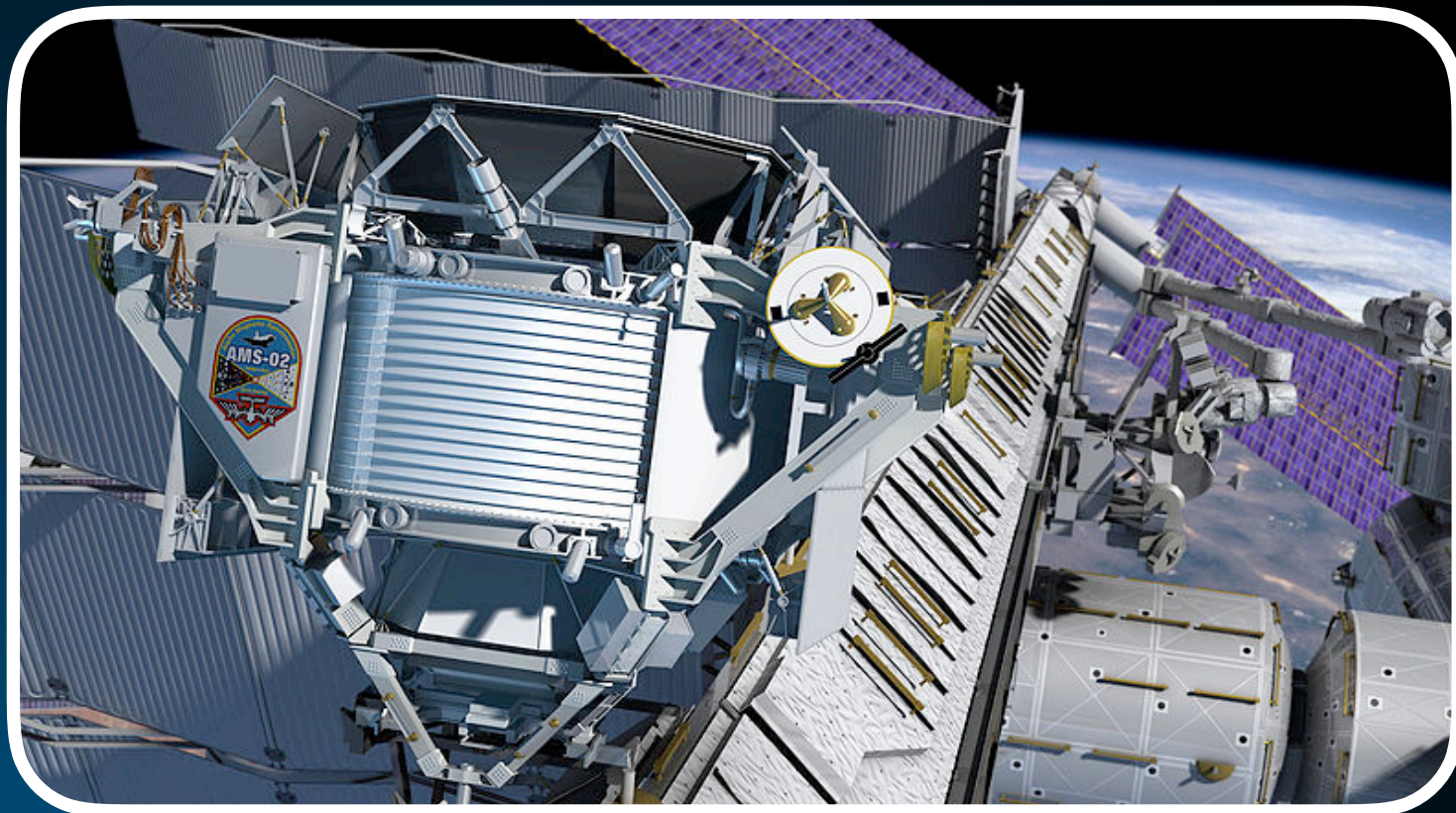
Spectral Features

Things that go bump in the night



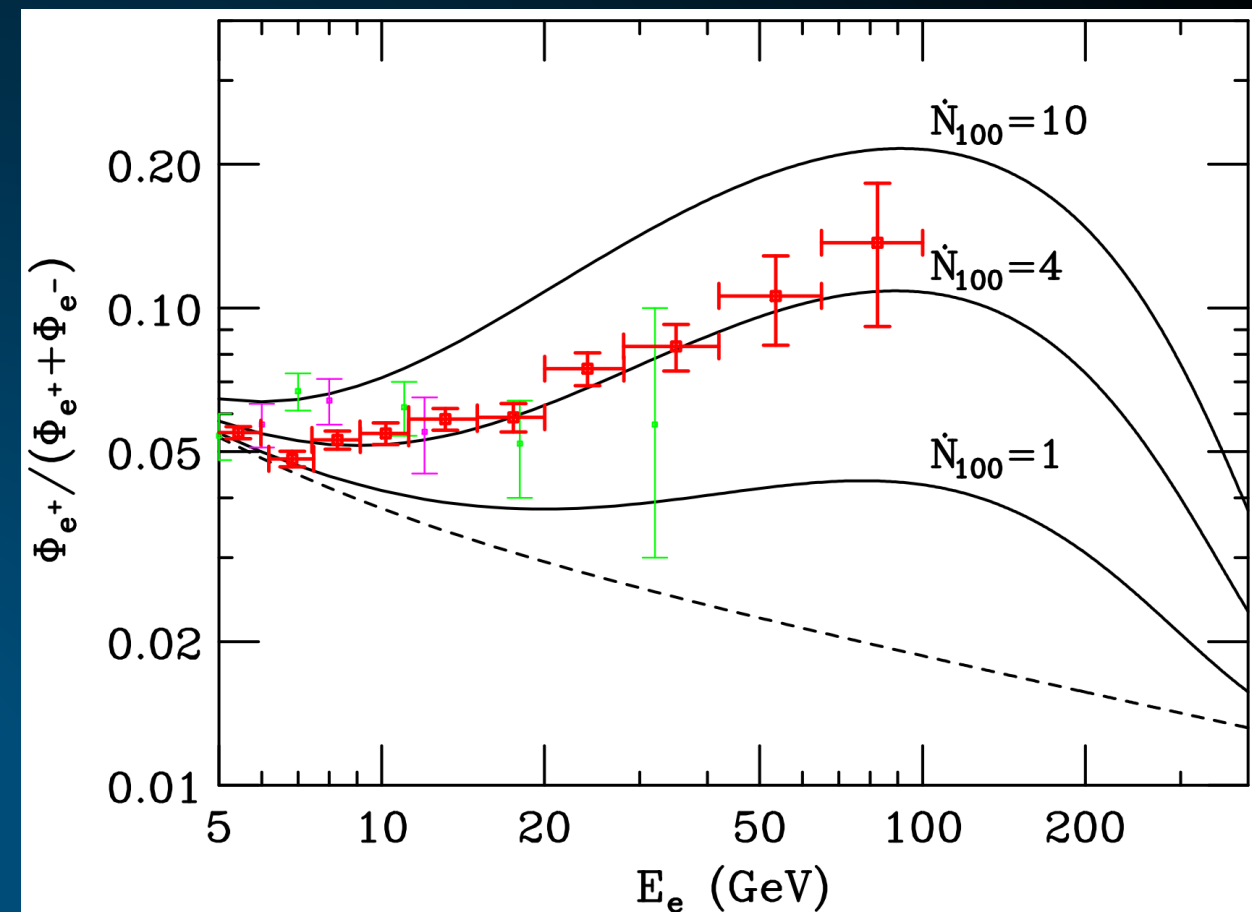
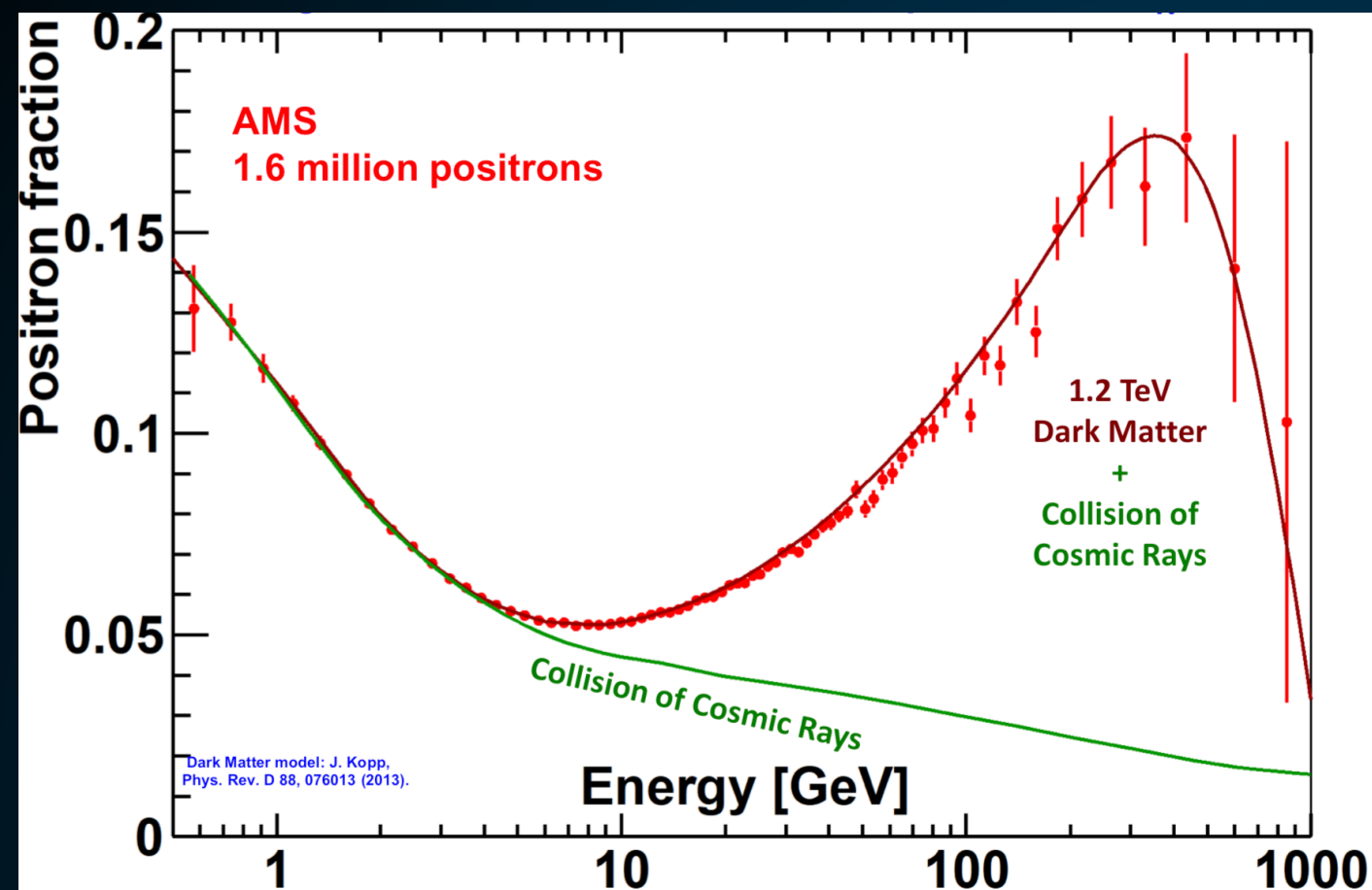
Spectral Features

Part 1: The Positron Excess



Spectral Features

Part 1: The Positron Excess



Dark Matter Models

- *Highly Sommerfeld Enhanced*
- *Leptophilic*

Pulsar Models

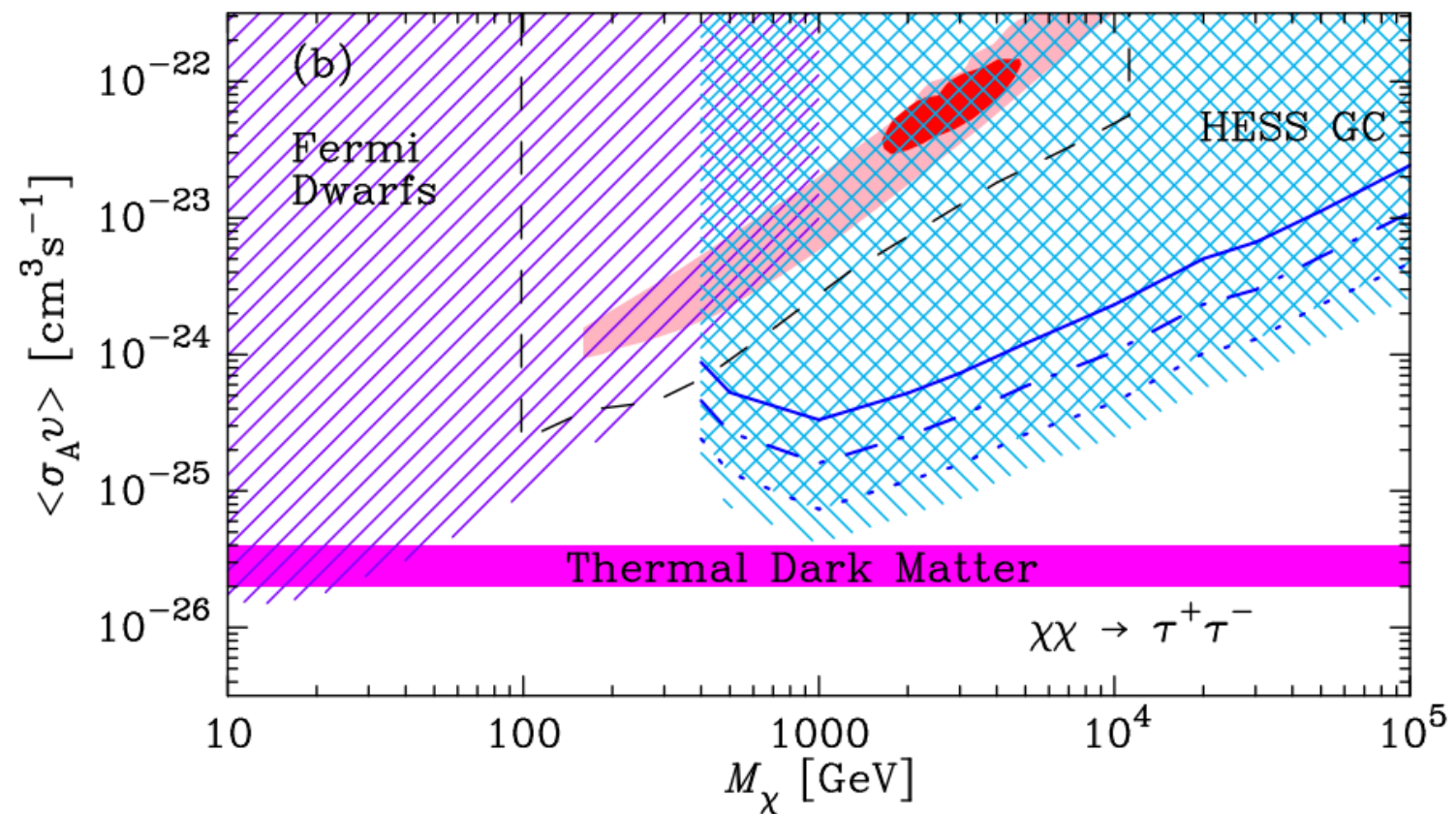
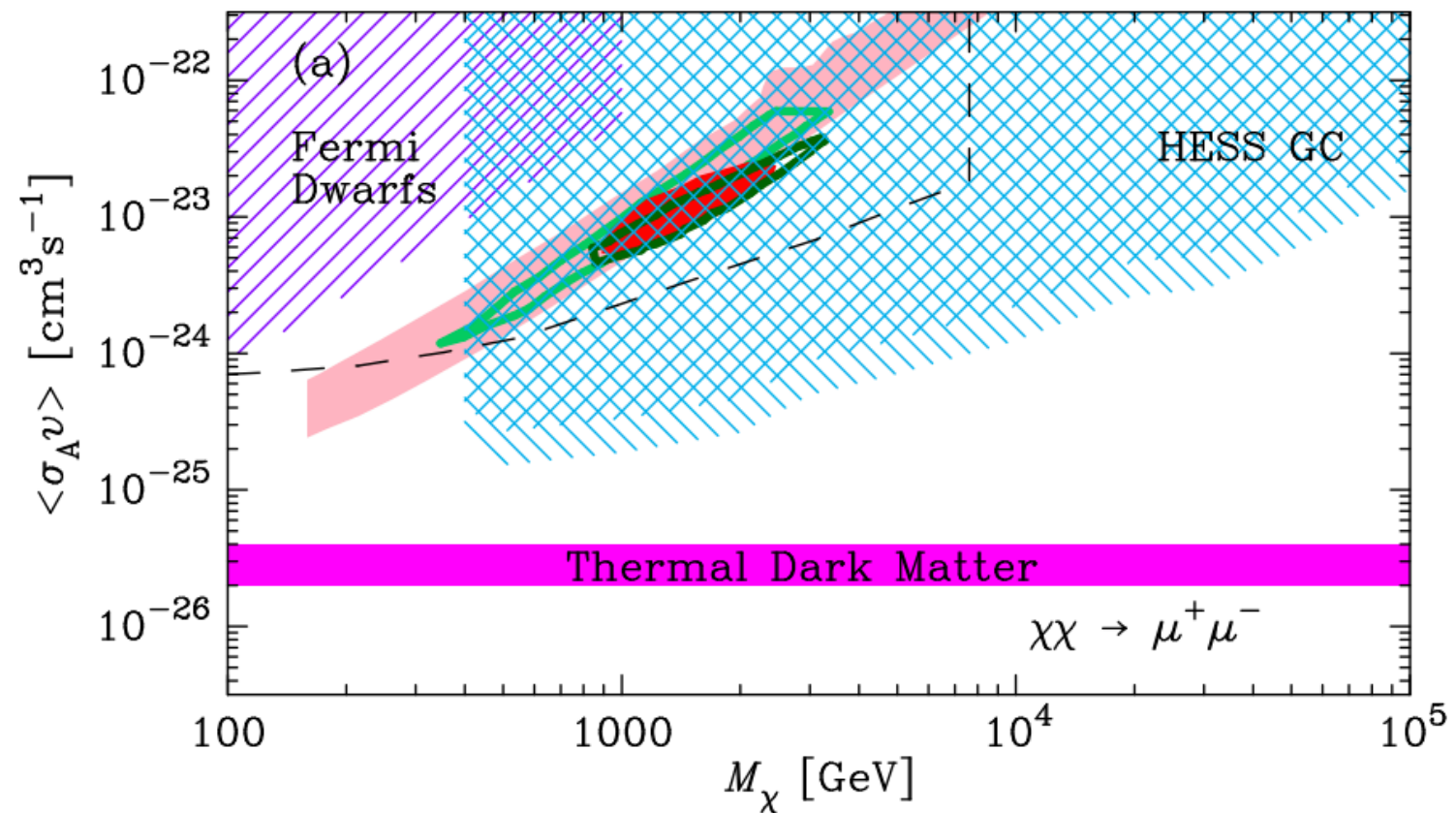
- *Efficient e^+e^- Production*
- *Hard e^+e^- Spectrum*

Spectral Features

Part 1: The Positron Excess

**Gamma-Ray Observations
are in strong tension with
dark matter models.**

Abazajian & Harding (2012; 1110.6151)



Part 1: The Positron Excess



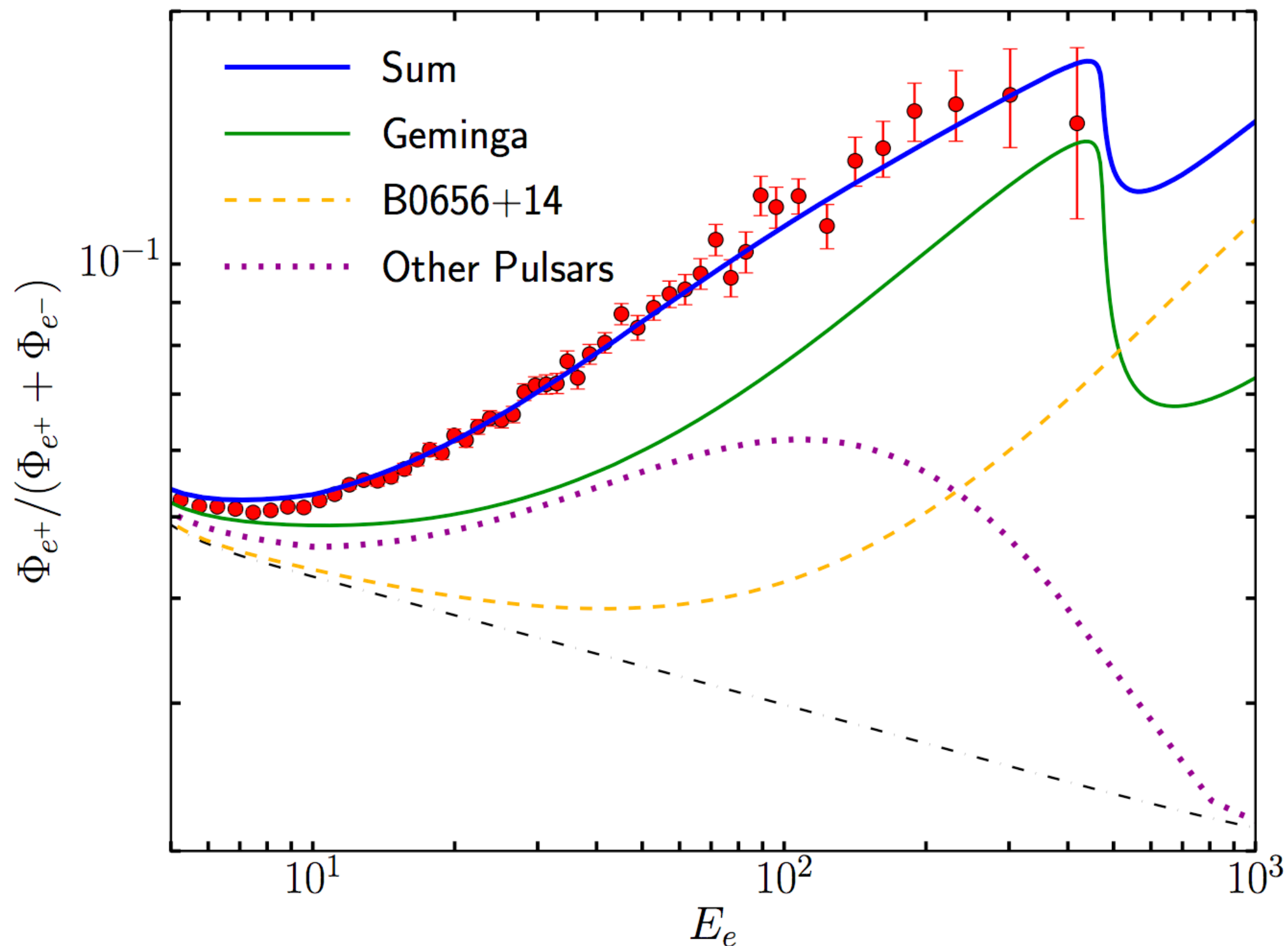
Moon (To Scale)

Geminga

PSR B0656+14

- **TeV Halos Surrounding Pulsars:**
 - **Hard e^+e^- injection spectrum**
 - **10-30% of spindown energy into e^+e^-**

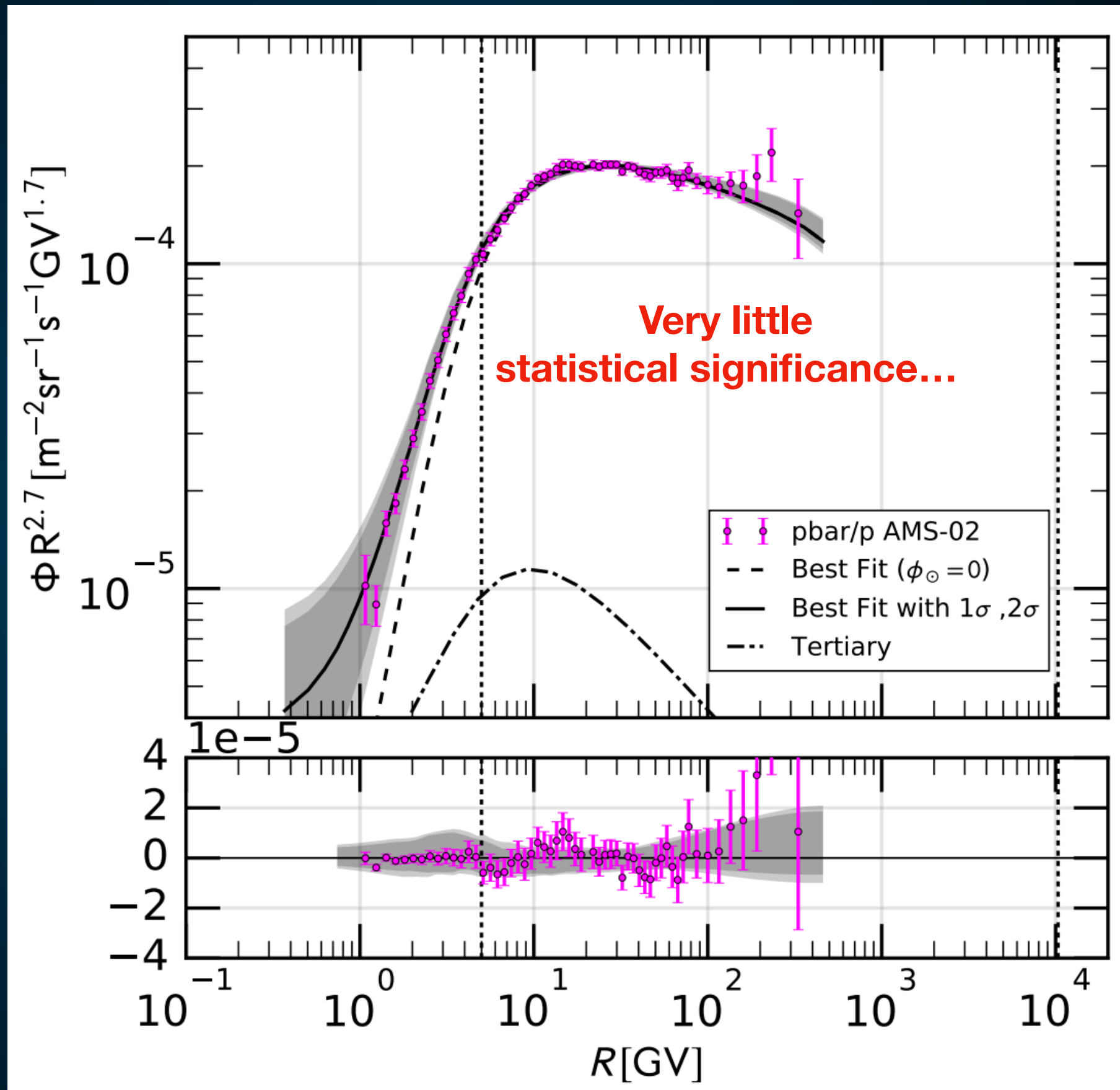
Part 1: The Positron Excess

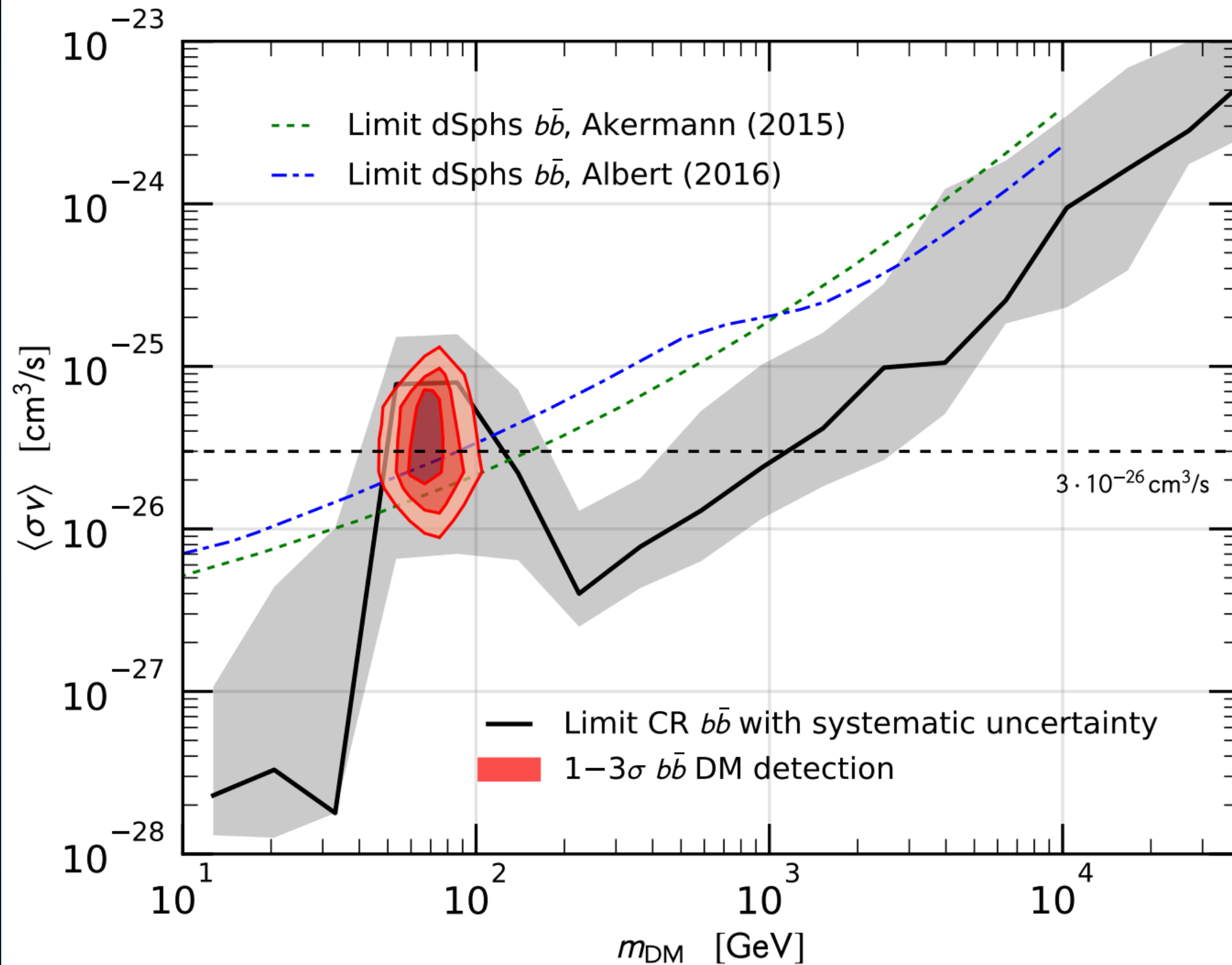


sars:

trum

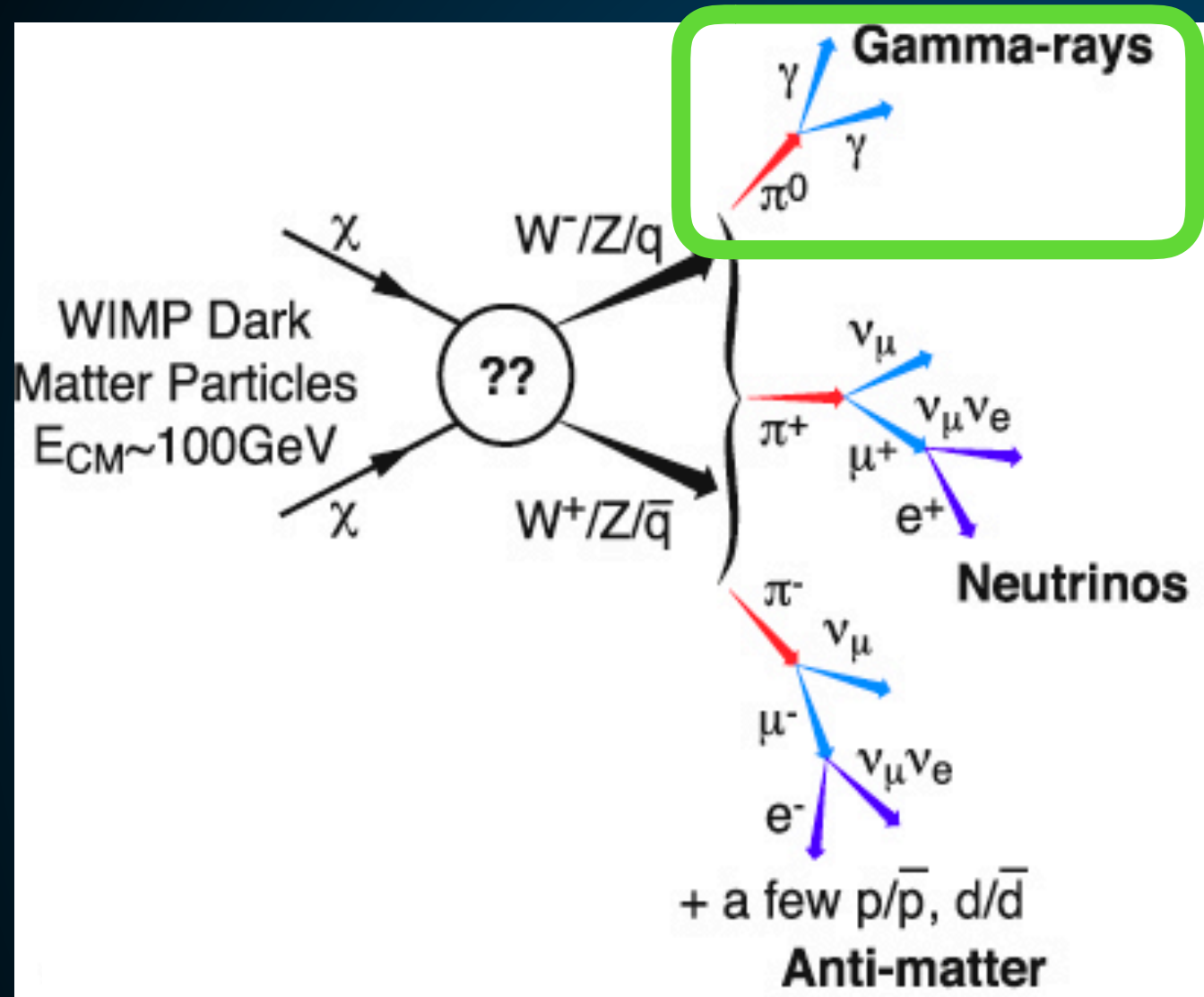
energy

Part 2: The Antiproton Excess

Part 2: The Antiproton Excess

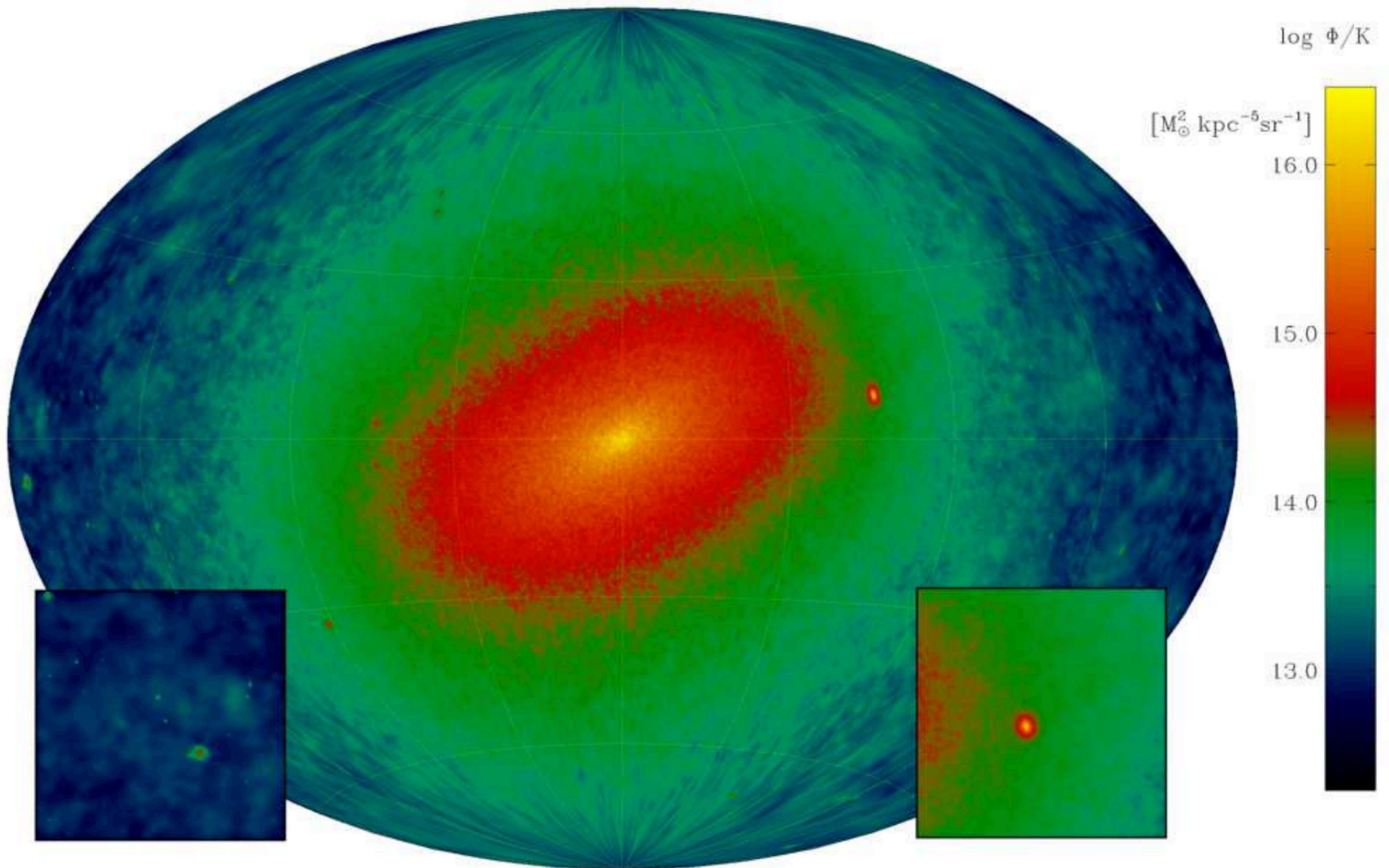
Angular Mapping

Using Gravity to Determine Where the Dark Matter Really Is



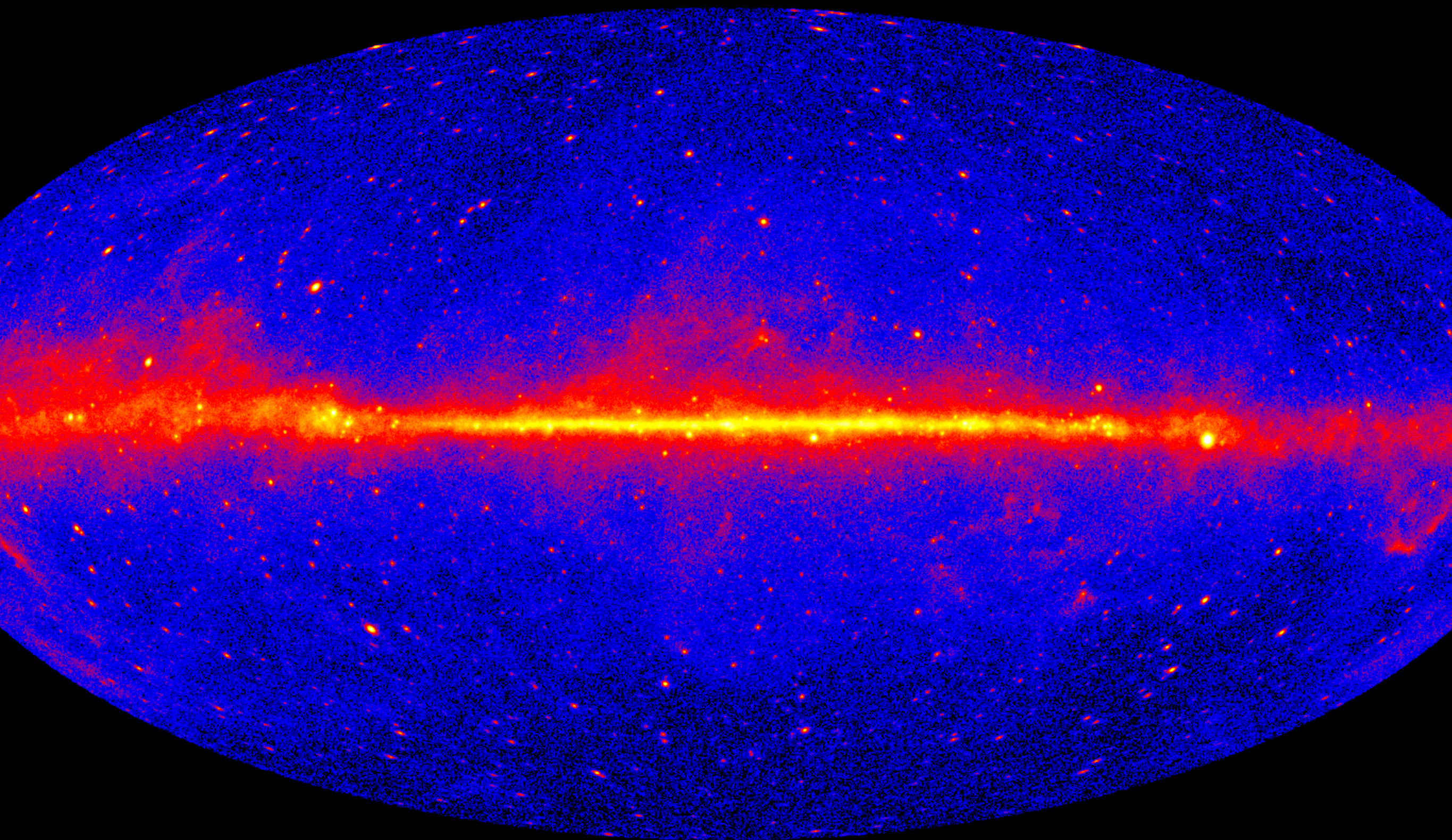
Angular Mapping

Using Gravity to Determine Where the Dark Matter Really Is



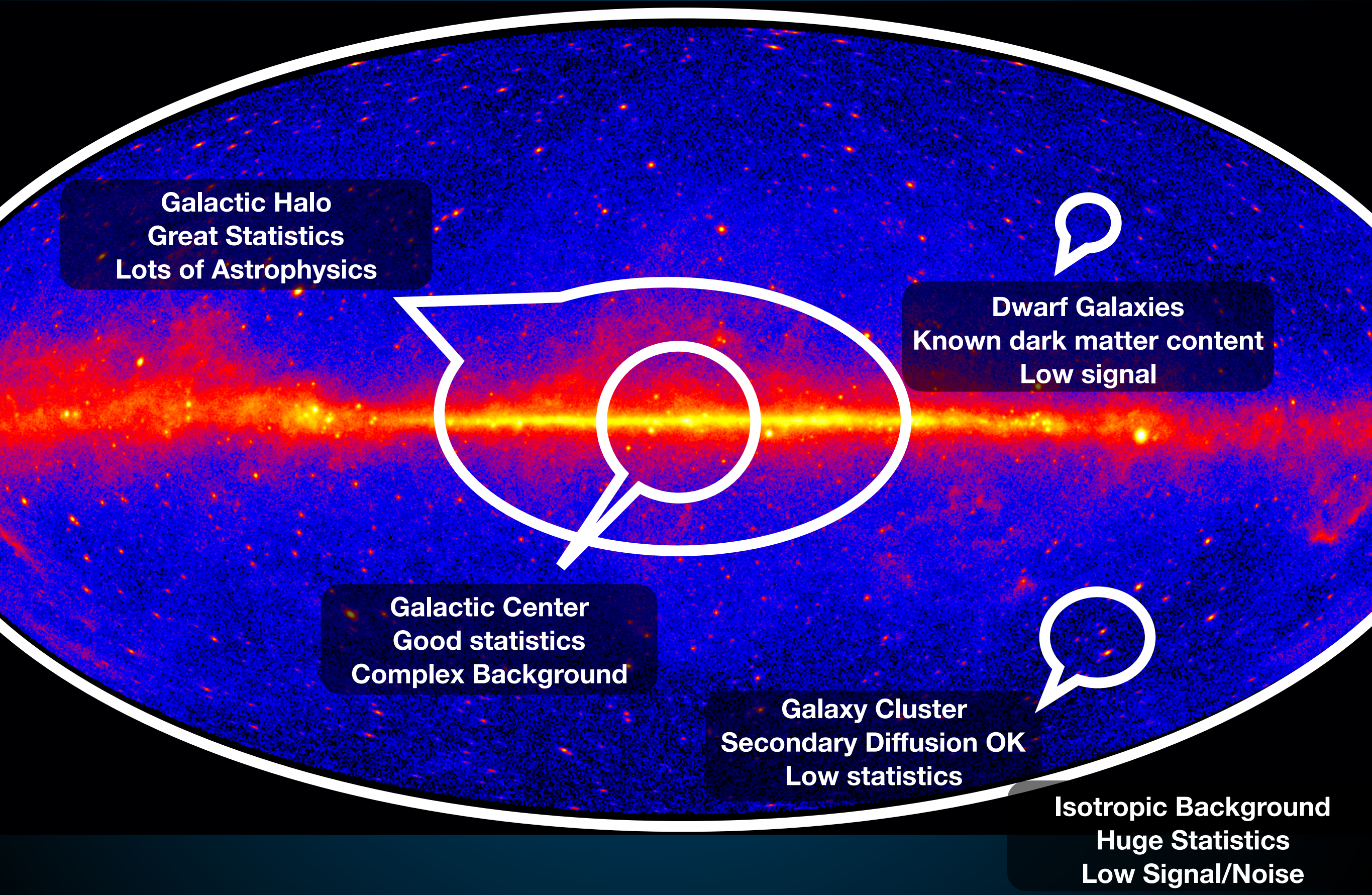
Angular Mapping

Using Gravity to Determine Where the Dark Matter Really Is



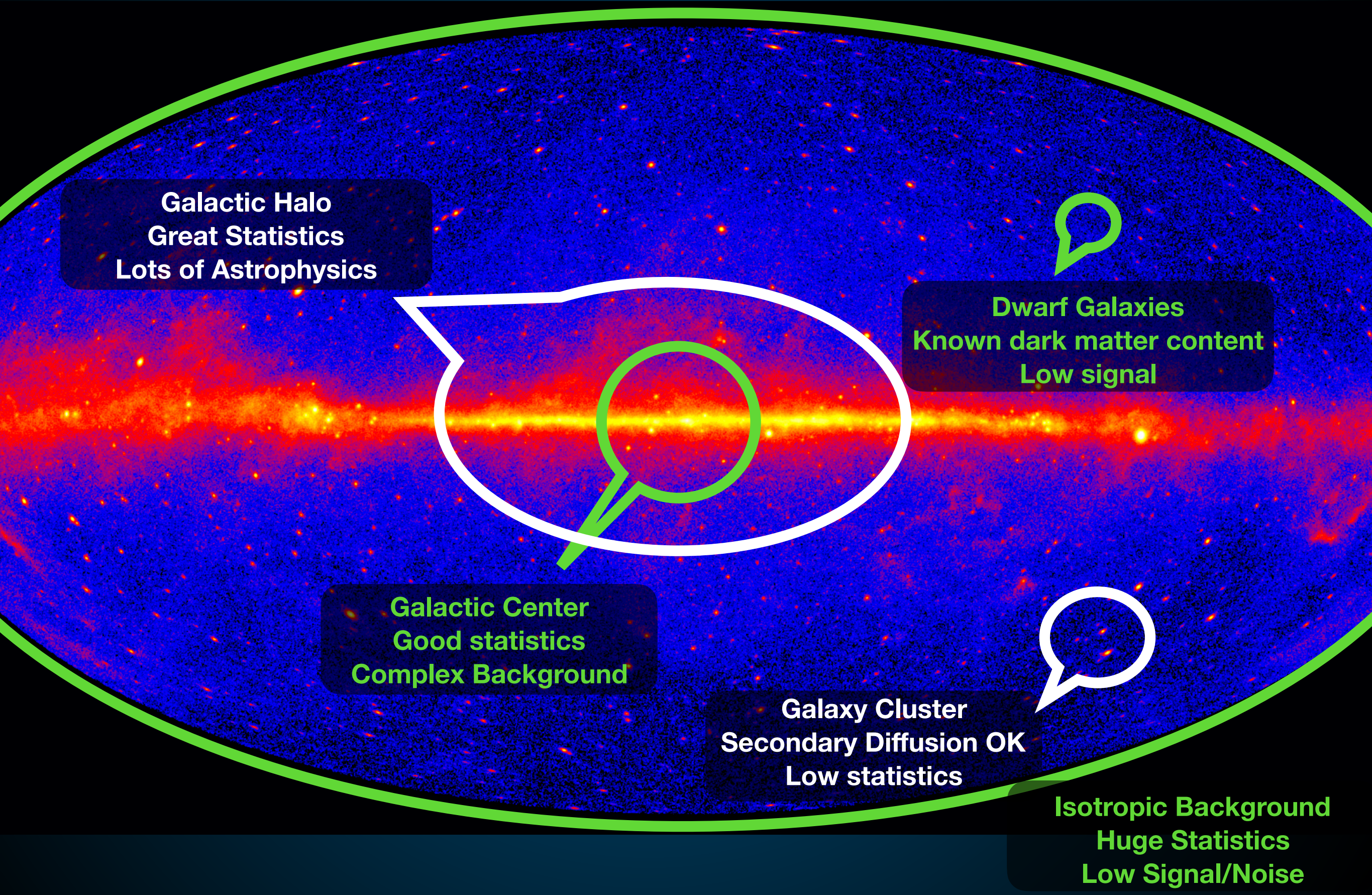
Angular Modeling

Using Gravity to Determine Where the Dark Matter Really Is



Angular Modeling

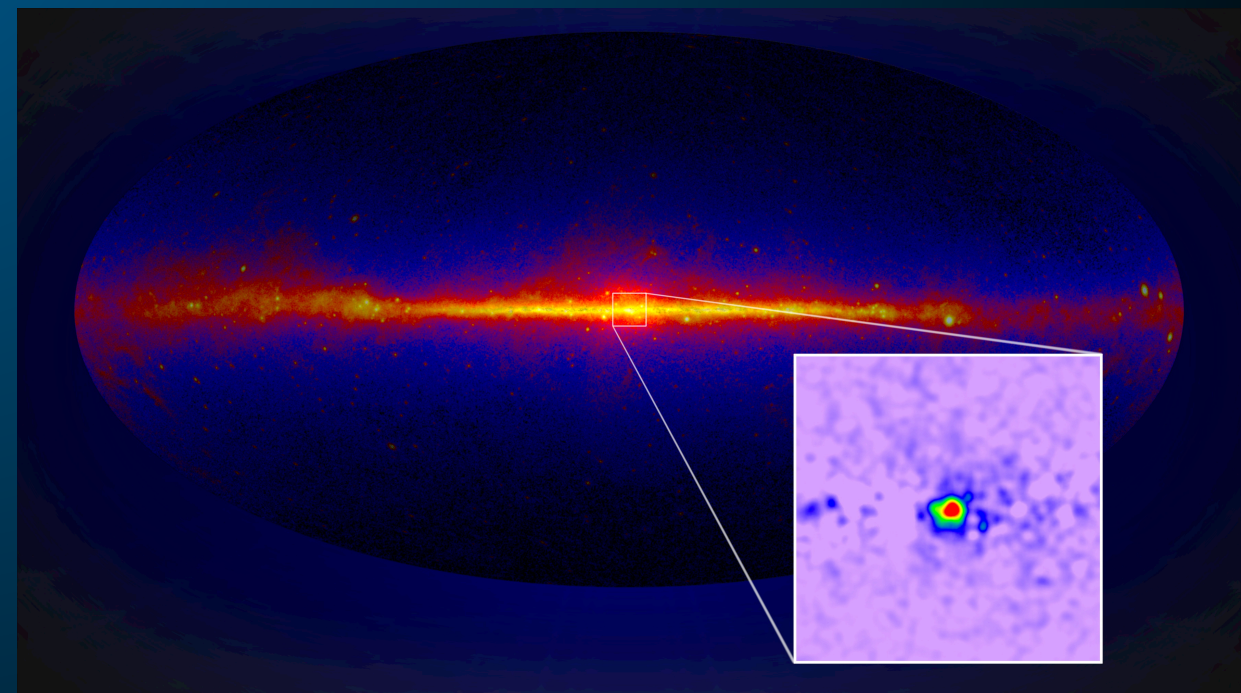
Using Gravity to Determine Where the Dark Matter Really Is



Angular Mapping

Part 1: The Galactic Center

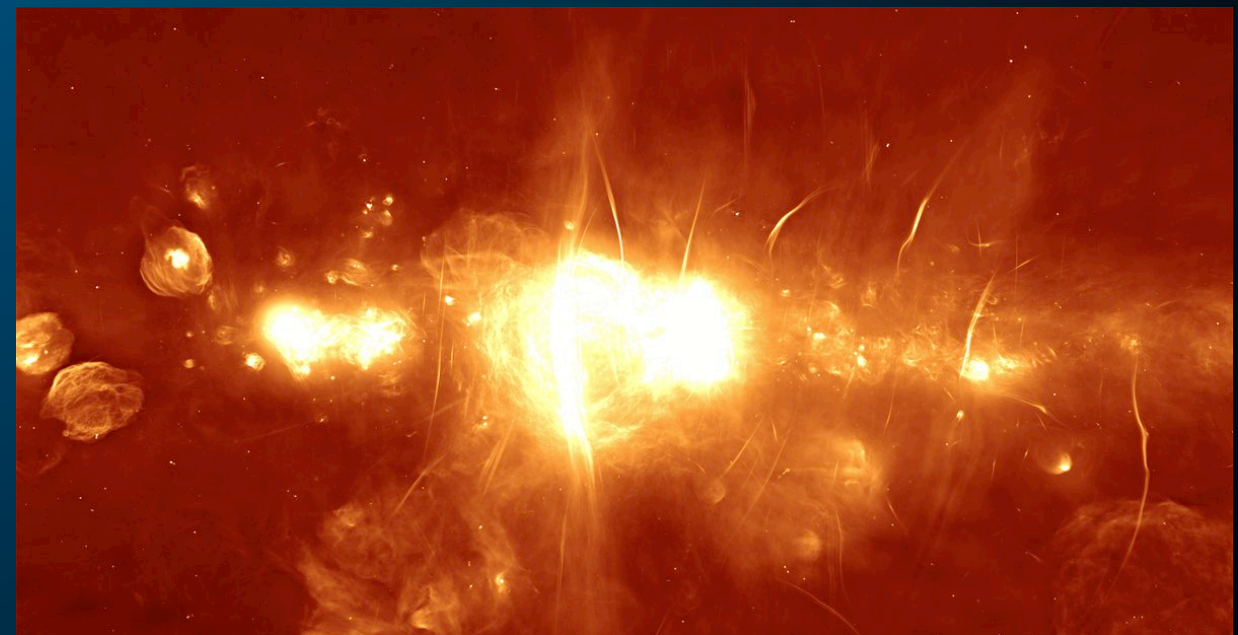
- **Model:**
 - **100 GeV dark matter particle annihilates to $b\bar{b}$**
 - **Annihilation Rate is Thermal Cross-Section**
- **Expected Galactic Center Flux (above 1 GeV):**
 - **$2 \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$**
- **Observed Flux:**
 - **$1 \times 10^{-10} \text{ erg cm}^{-2} \text{ s}^{-1}$**



Angular Mapping

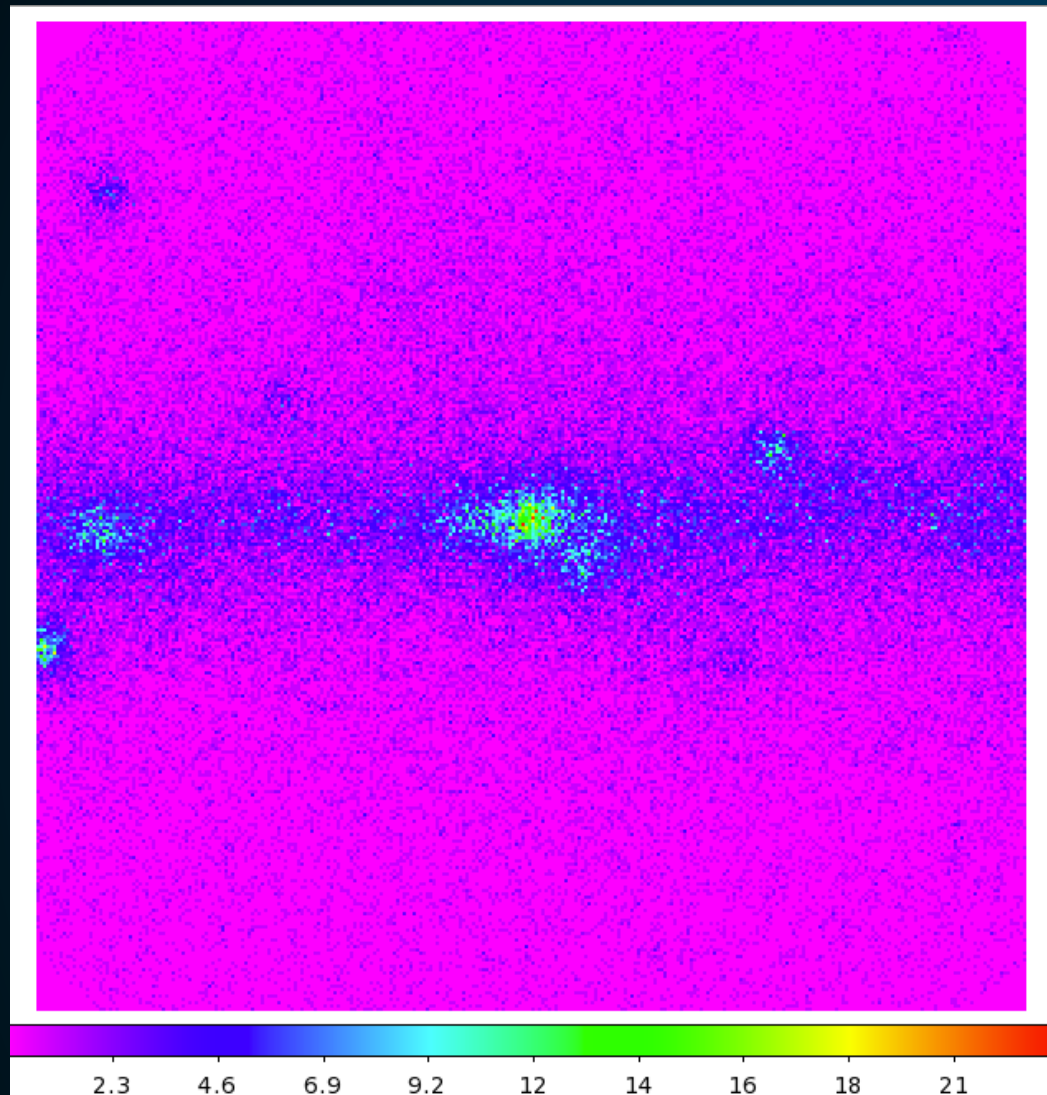
Part 1: The Galactic Center

- **Model:**
 - 100 GeV dark matter particle annihilates to $b\bar{b}$
 - Annihilation Rate is Thermal Cross-Section
- **Expected Galactic Center Radio Flux:**
 - $2 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$
- **Observed Flux:**
 - $5 \times 10^{-10} \text{ erg cm}^{-2} \text{ s}^{-1}$



Angular Mapping

Part 1: The Galactic Center



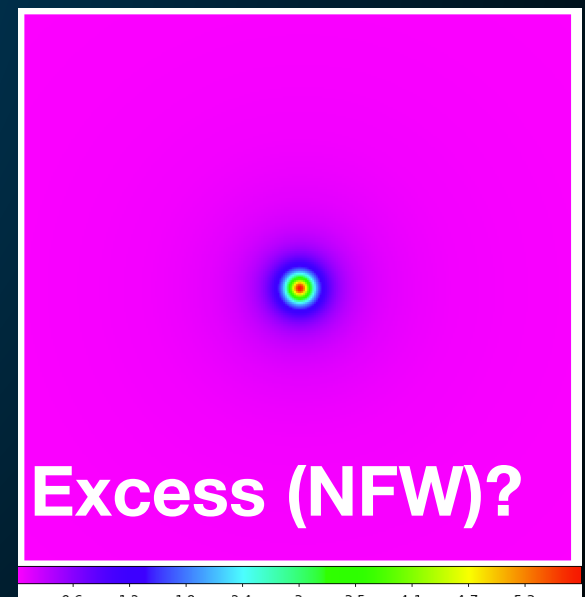
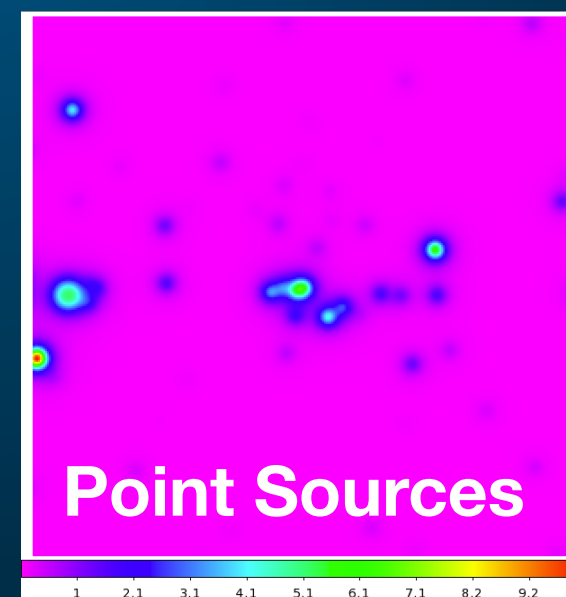
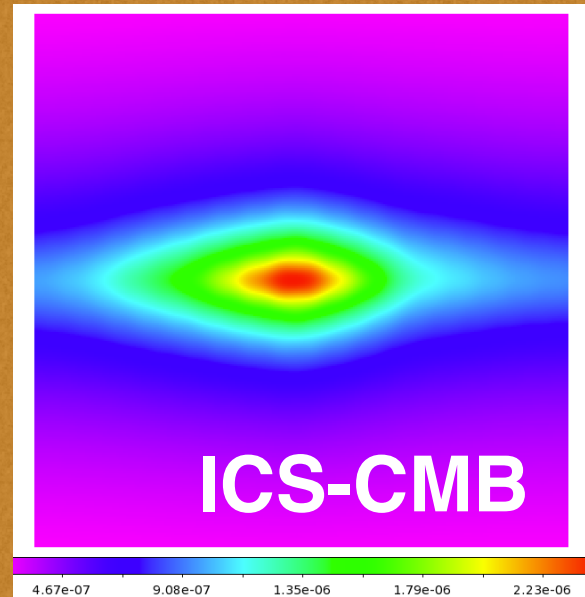
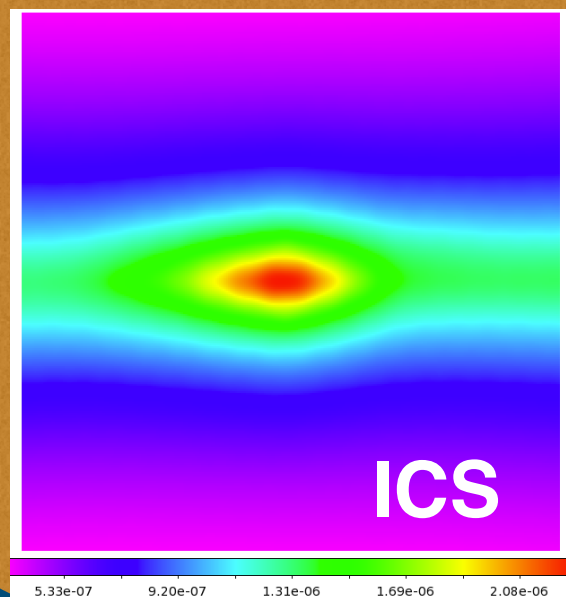
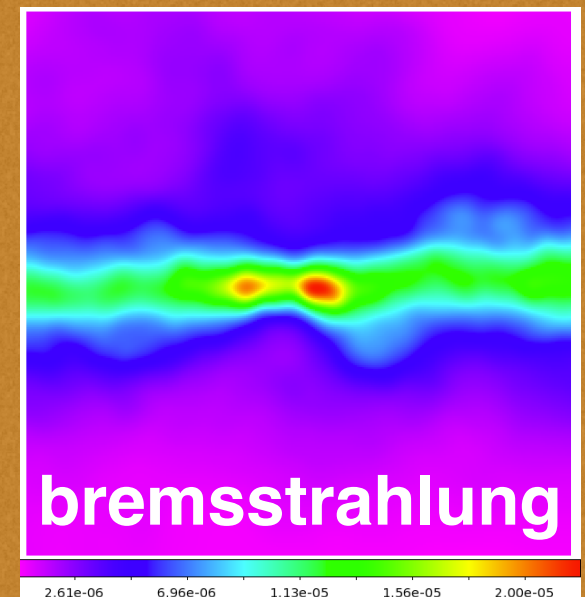
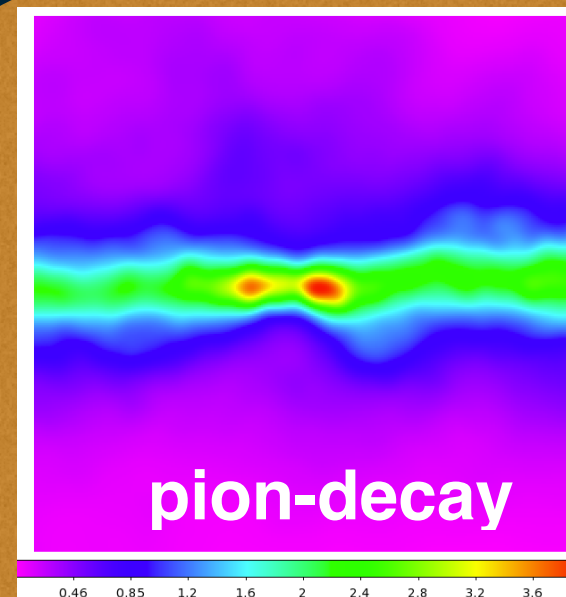
Data

750 — 950 MeV

Best Angular Resolution Cut

$10^\circ \times 10^\circ$ ROI

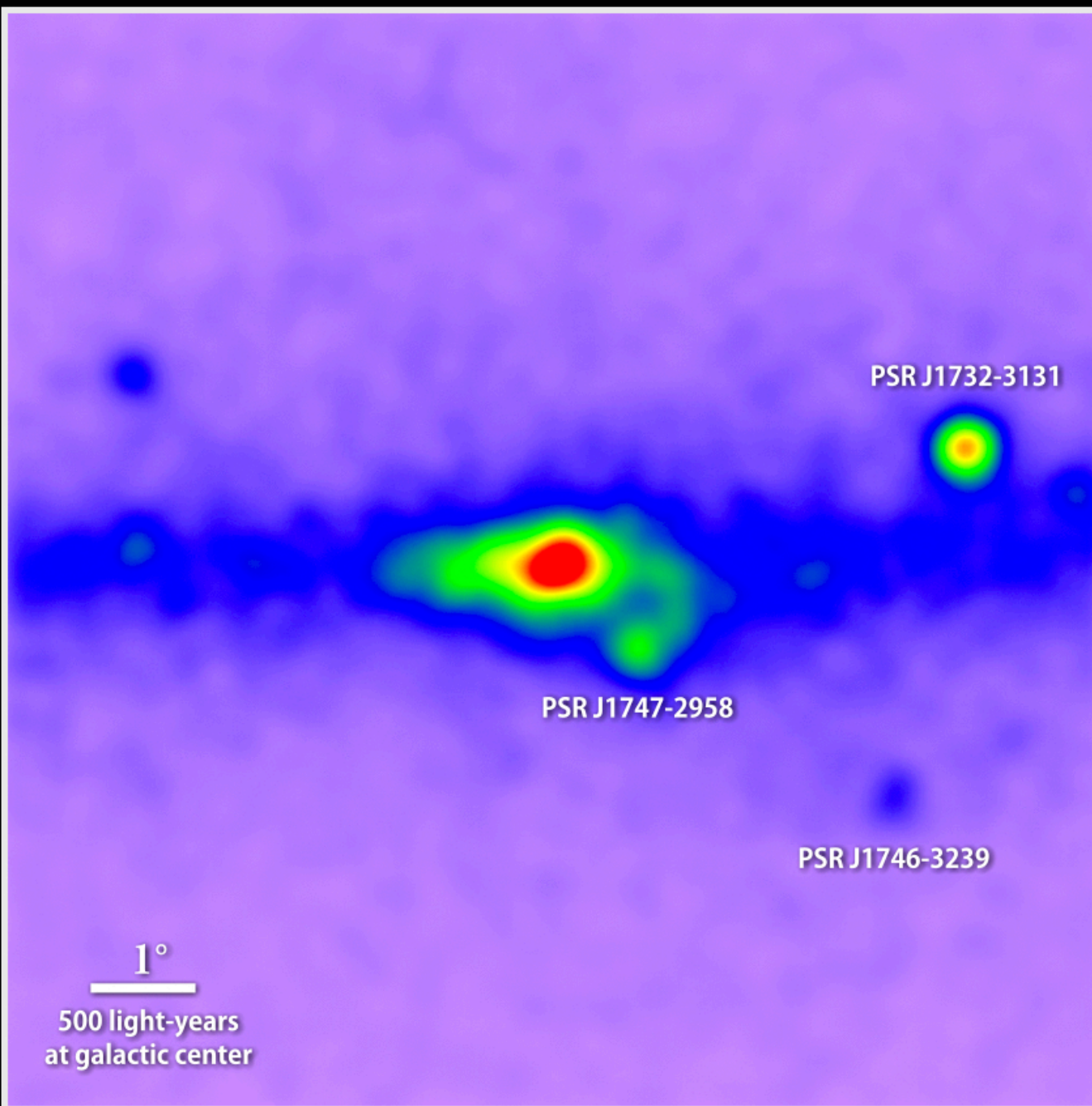
=



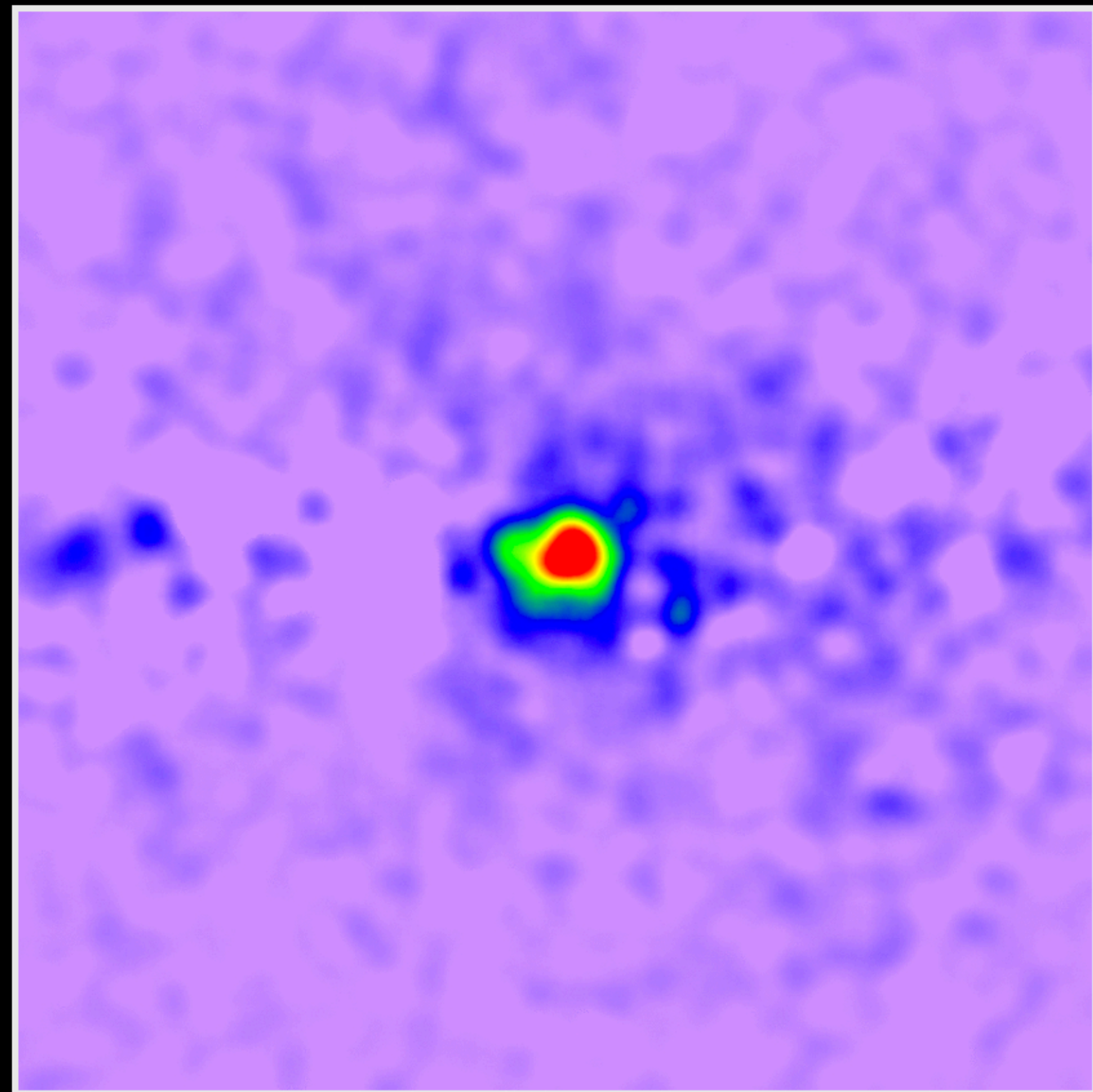
Angular Mapping

Part 1: The Galactic Center

Uncovering a gamma-ray excess at the galactic center



Unprocessed map of 1.0 to 3.16 GeV gamma rays

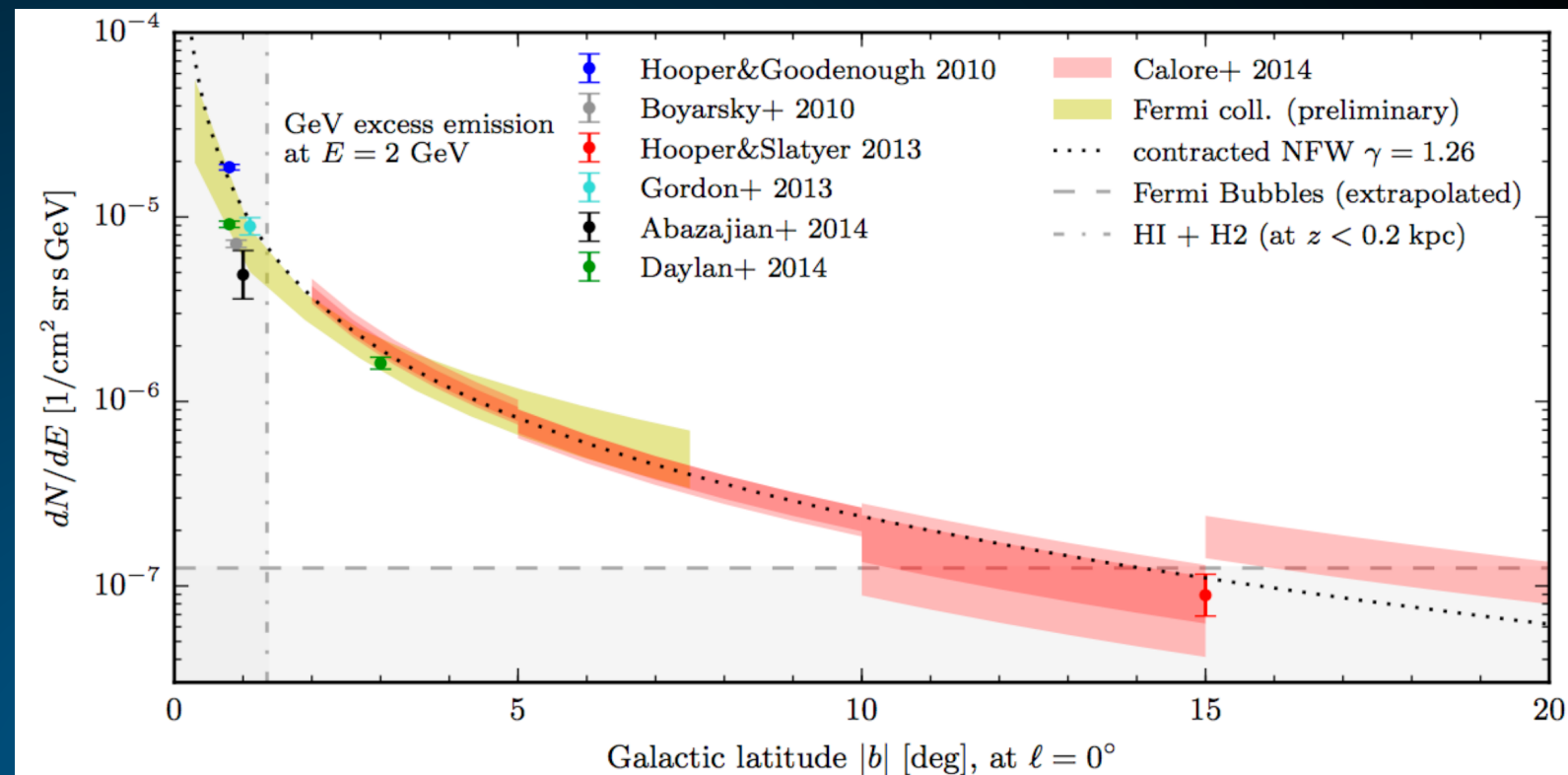
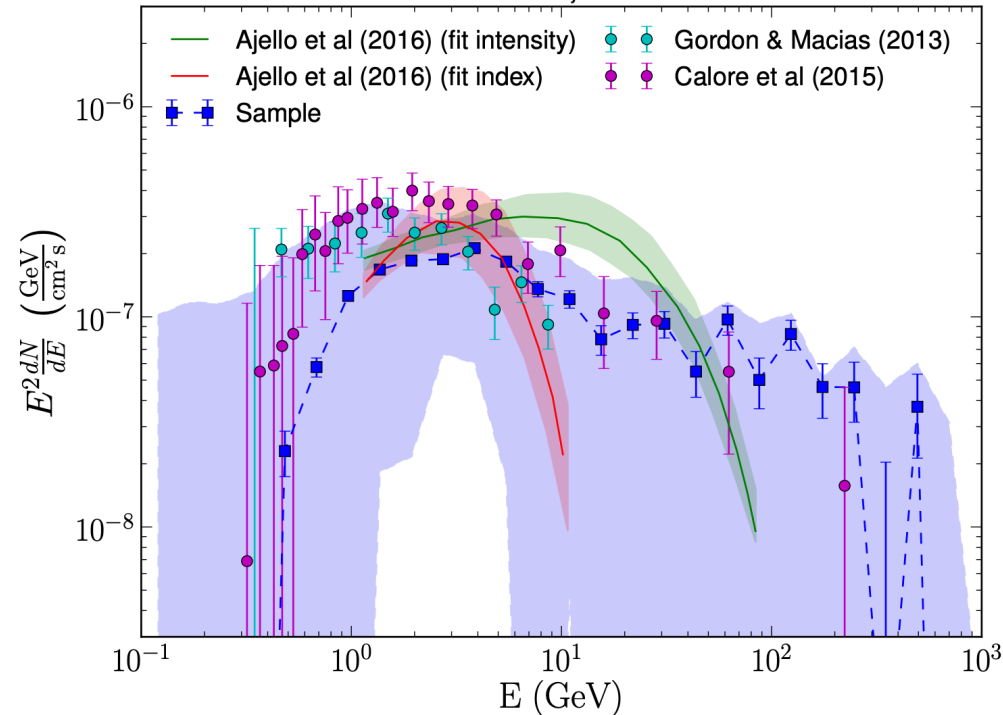


Known sources removed

Angular Mapping

Part 1: The Galactic Center

GC excess, all cases

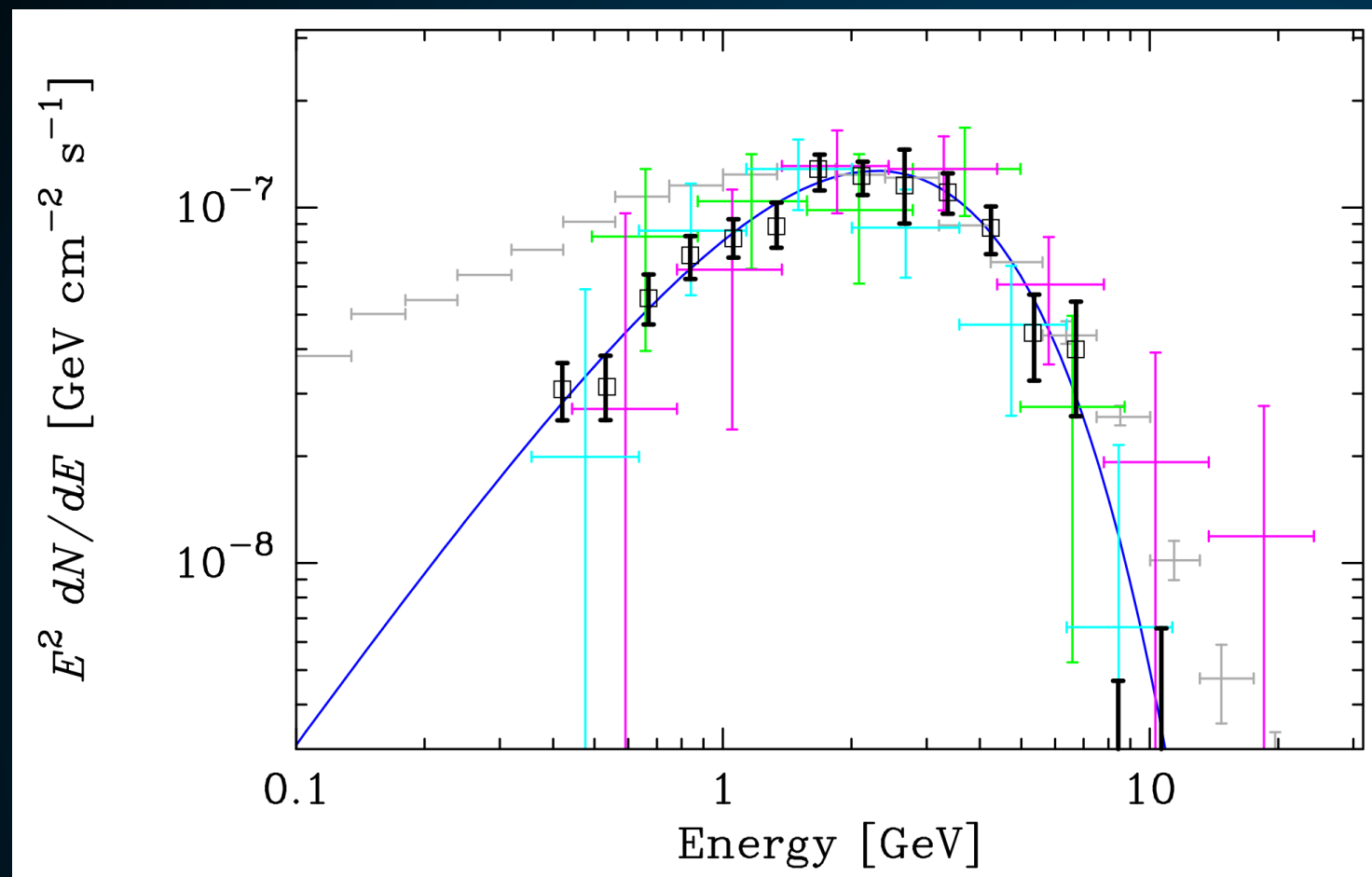


- **Result is extremely significant.**
- **Spectrum and morphology match expectations from dark matter.**

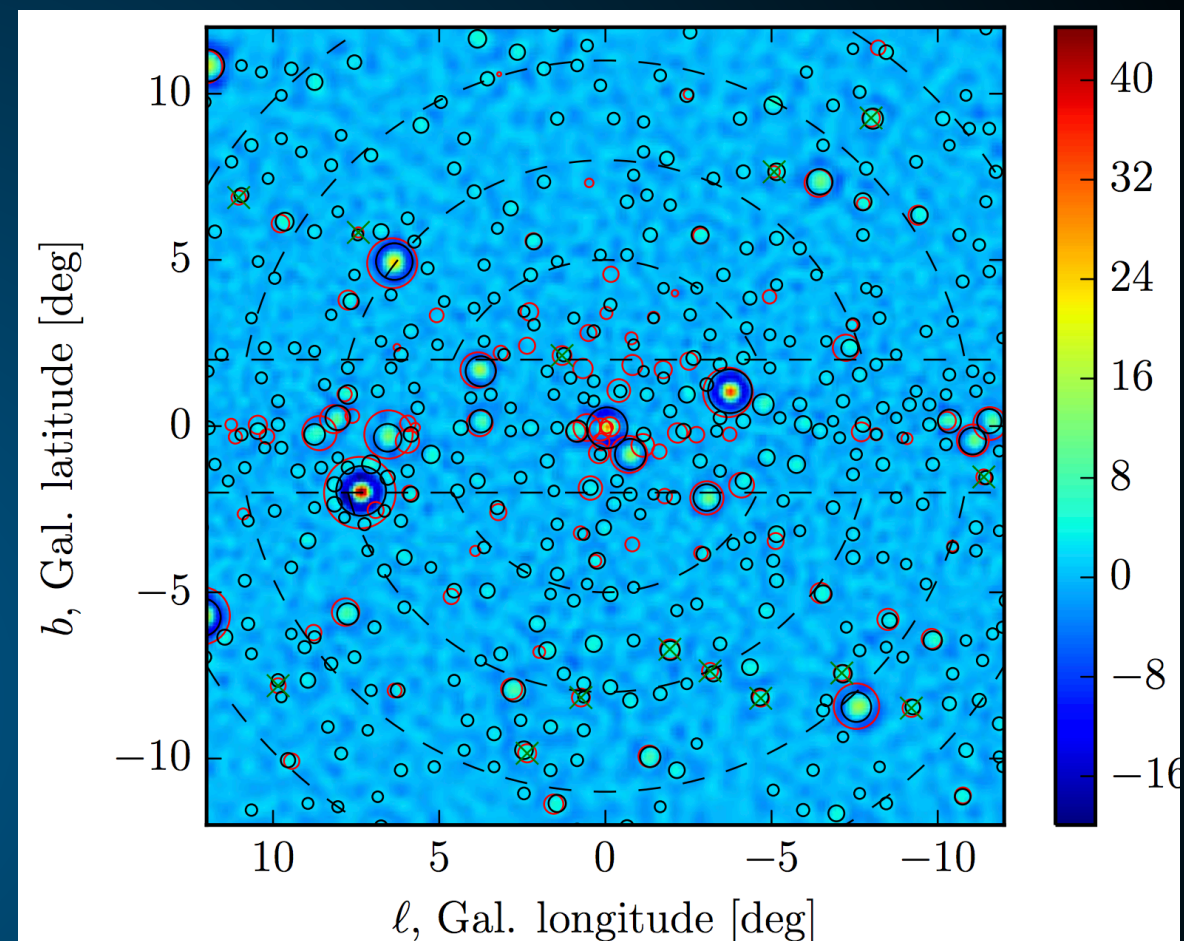
Angular Mapping

Part 1: The Galactic Center

Abazajian (1011.4275)



Bartels et al. (1506.05104)

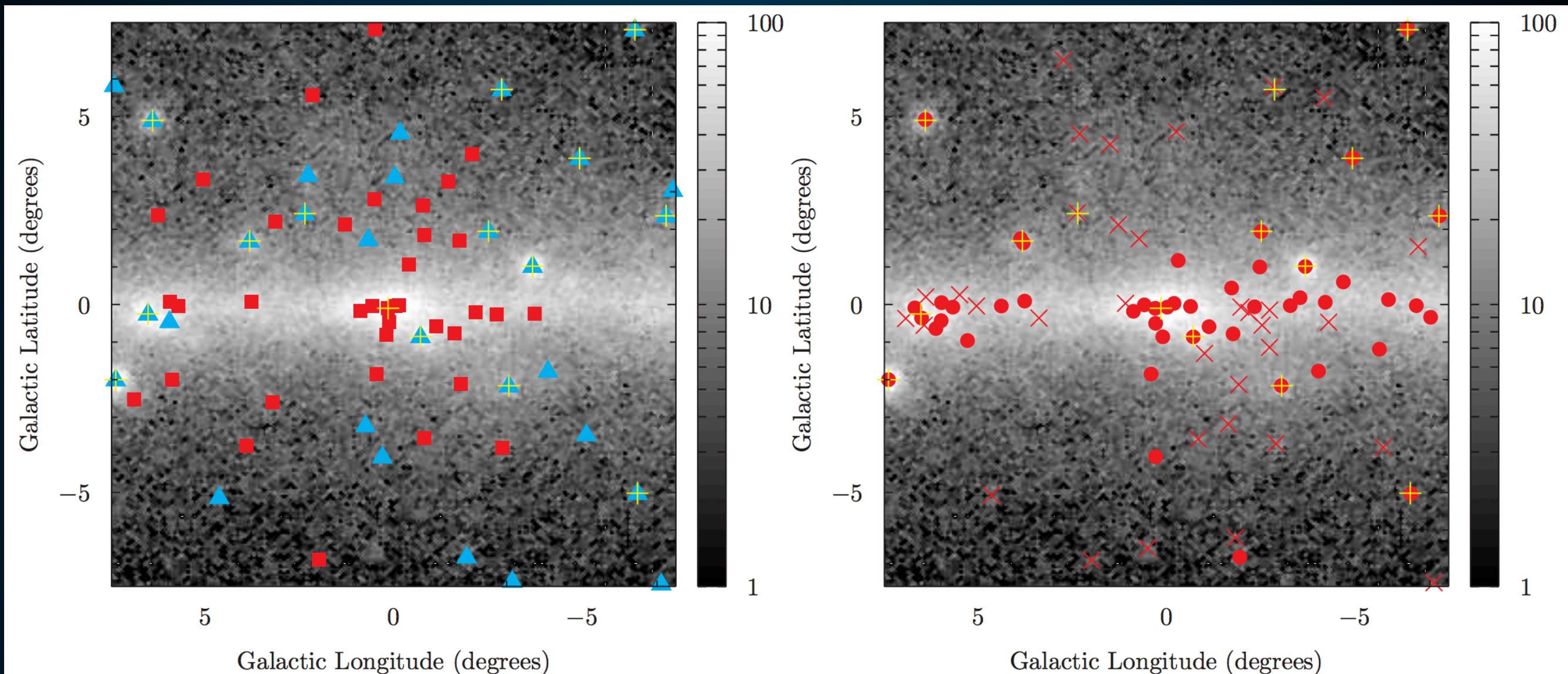


- **However, this emission may also be produced by a population of pulsars clustered in the Galactic bulge.**

Angular Mapping

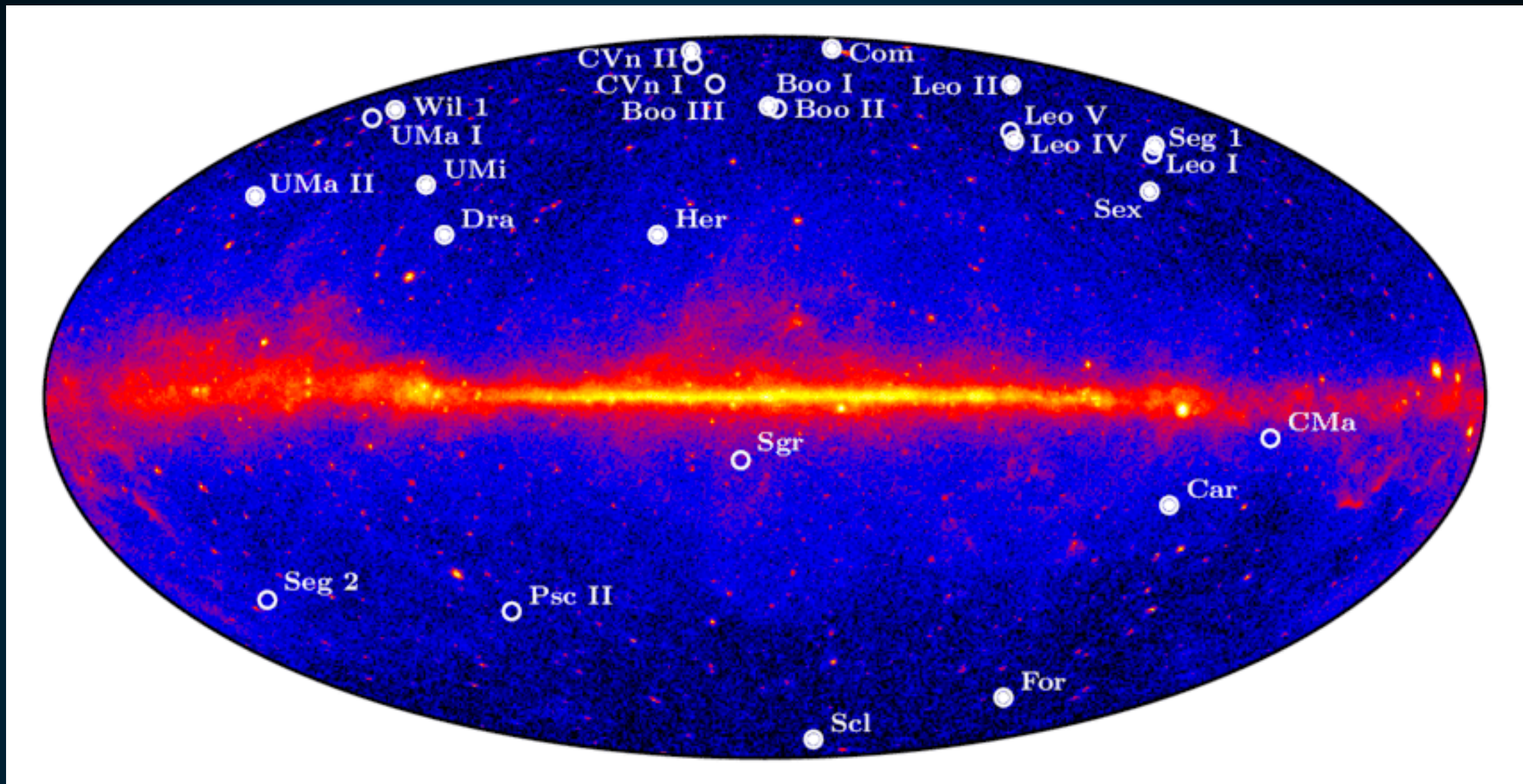
Part 1: The Galactic Center

Fermi-LAT Collaboration (1511.02938)



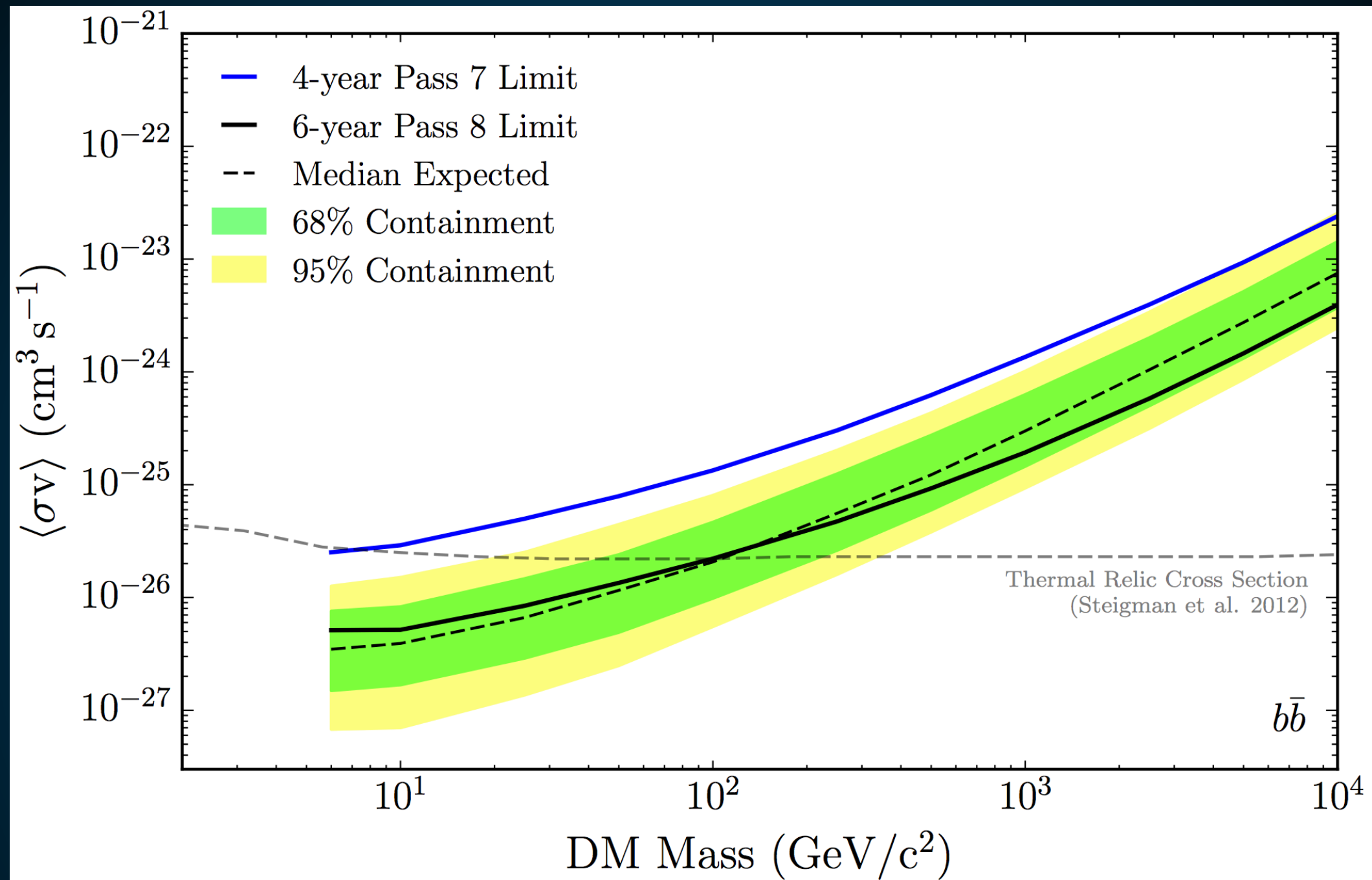
- **Models are very uncertain!**

Part 2: Dwarf Spheroidal Galaxies



- **Dwarf Spheroidal Galaxies**
 - Much dimmer dark matter signal
 - **Dark Matter density is measured by rotation curves**
 - **Very little background**

Part 2: Dwarf Spheroidal Galaxies

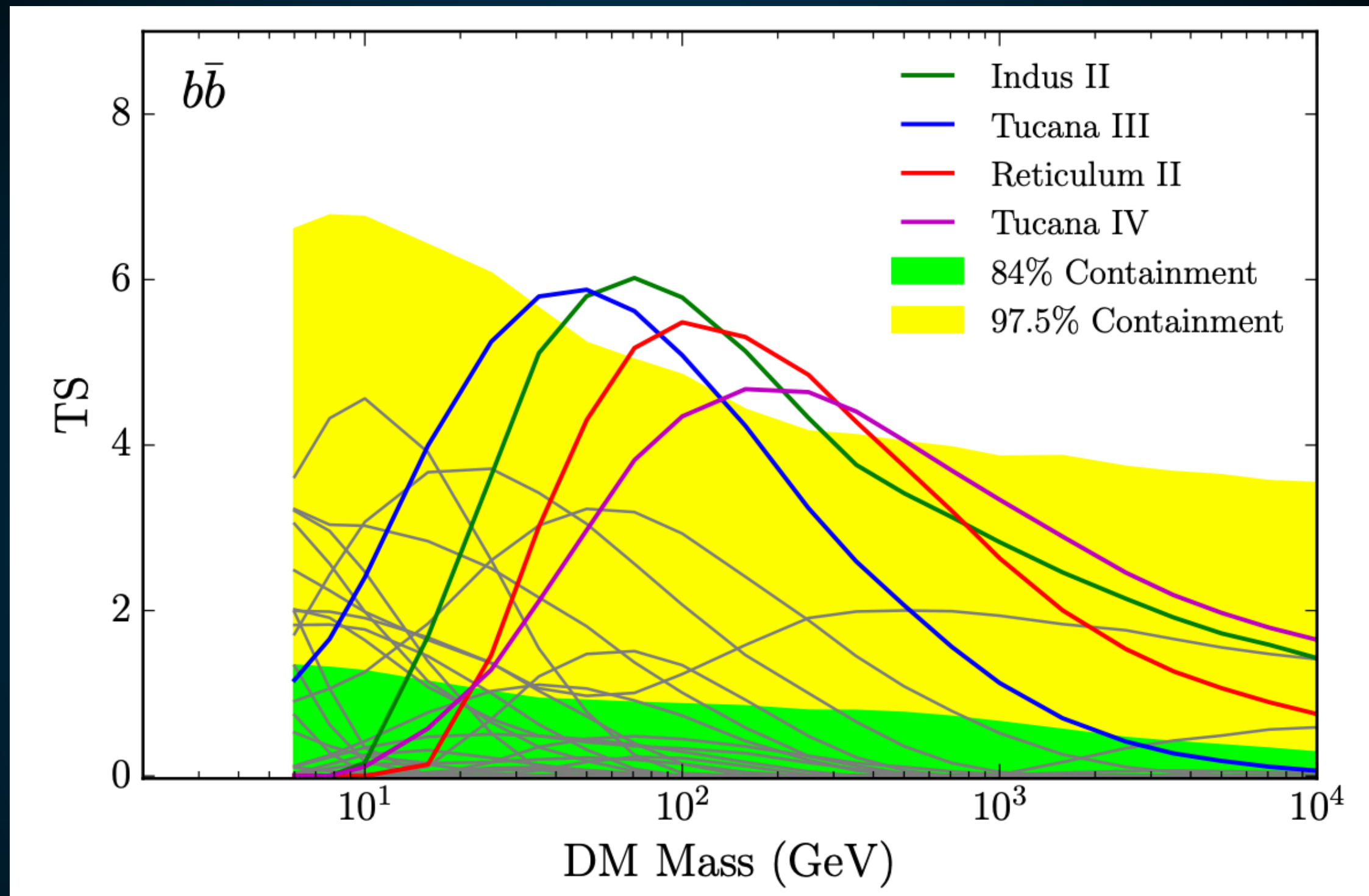


- **Constraints hit the thermal cross-section!**

Angular Mapping

Fermi-LAT Collaboration (2016; 1611.03184)

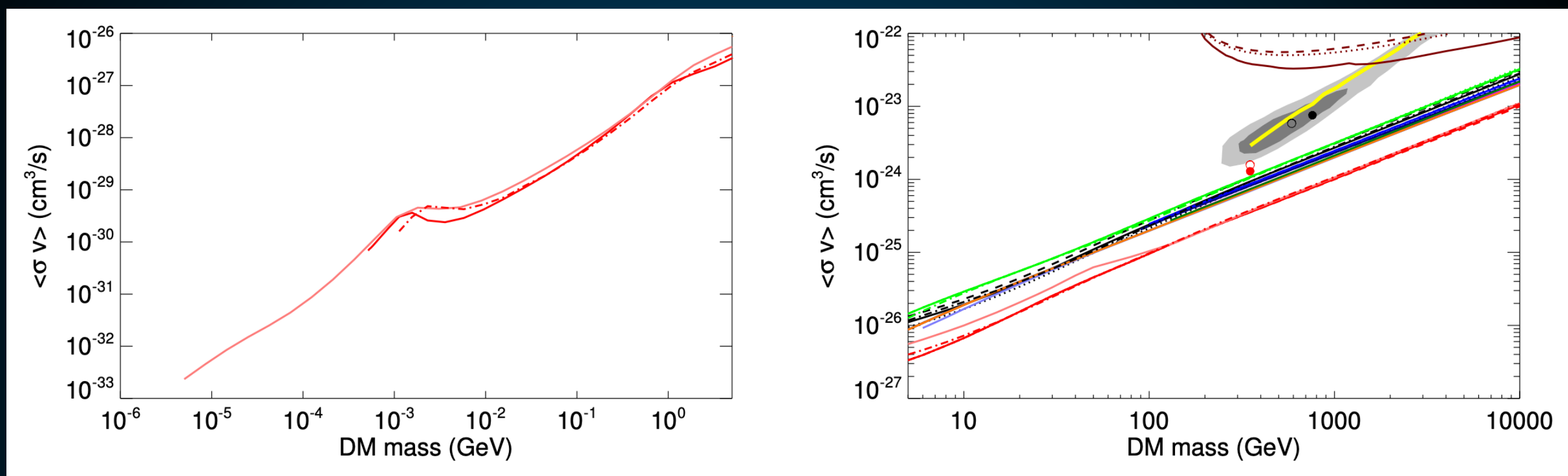
Part 2: Dwarf Spheroidal Galaxies



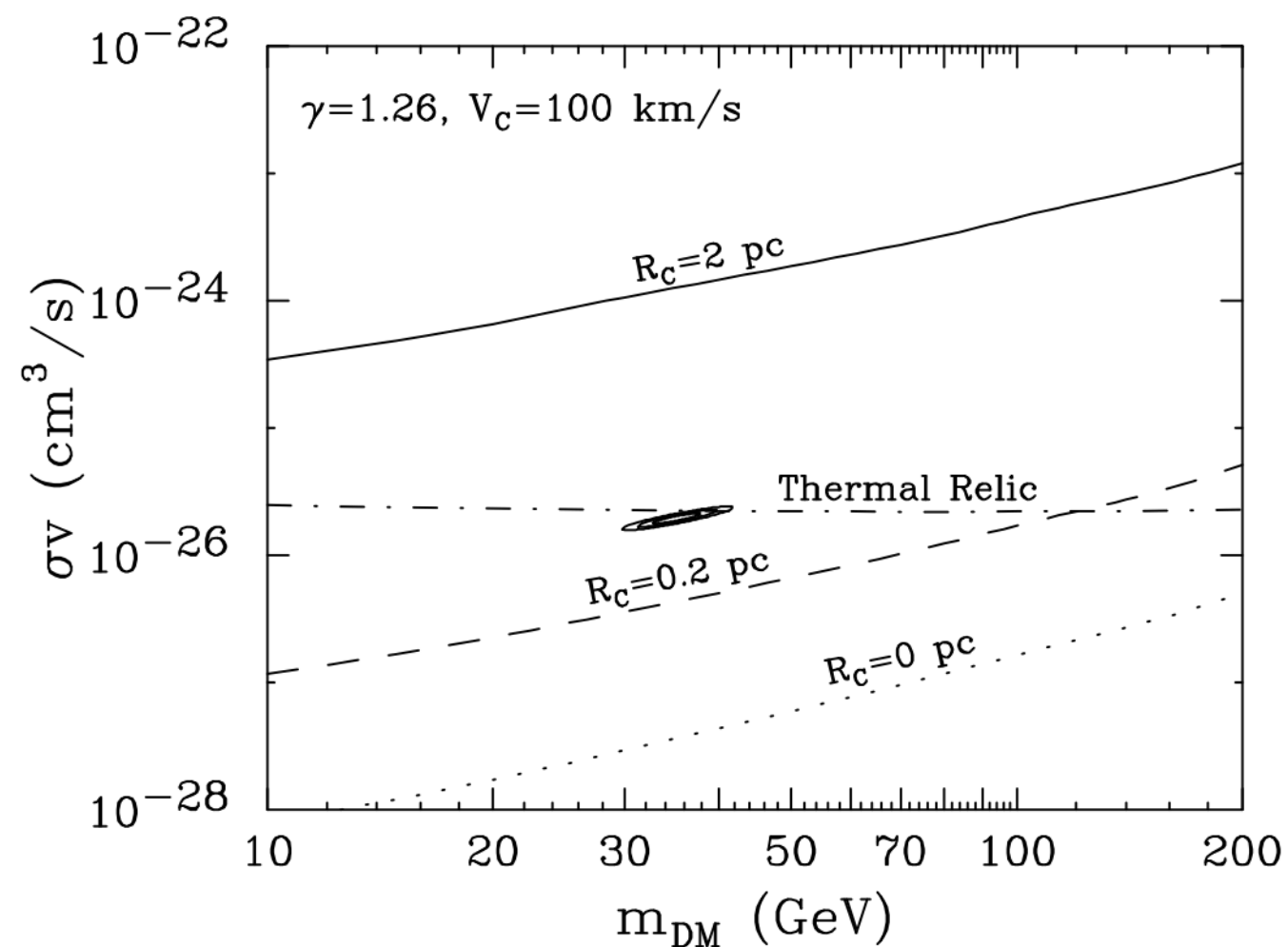
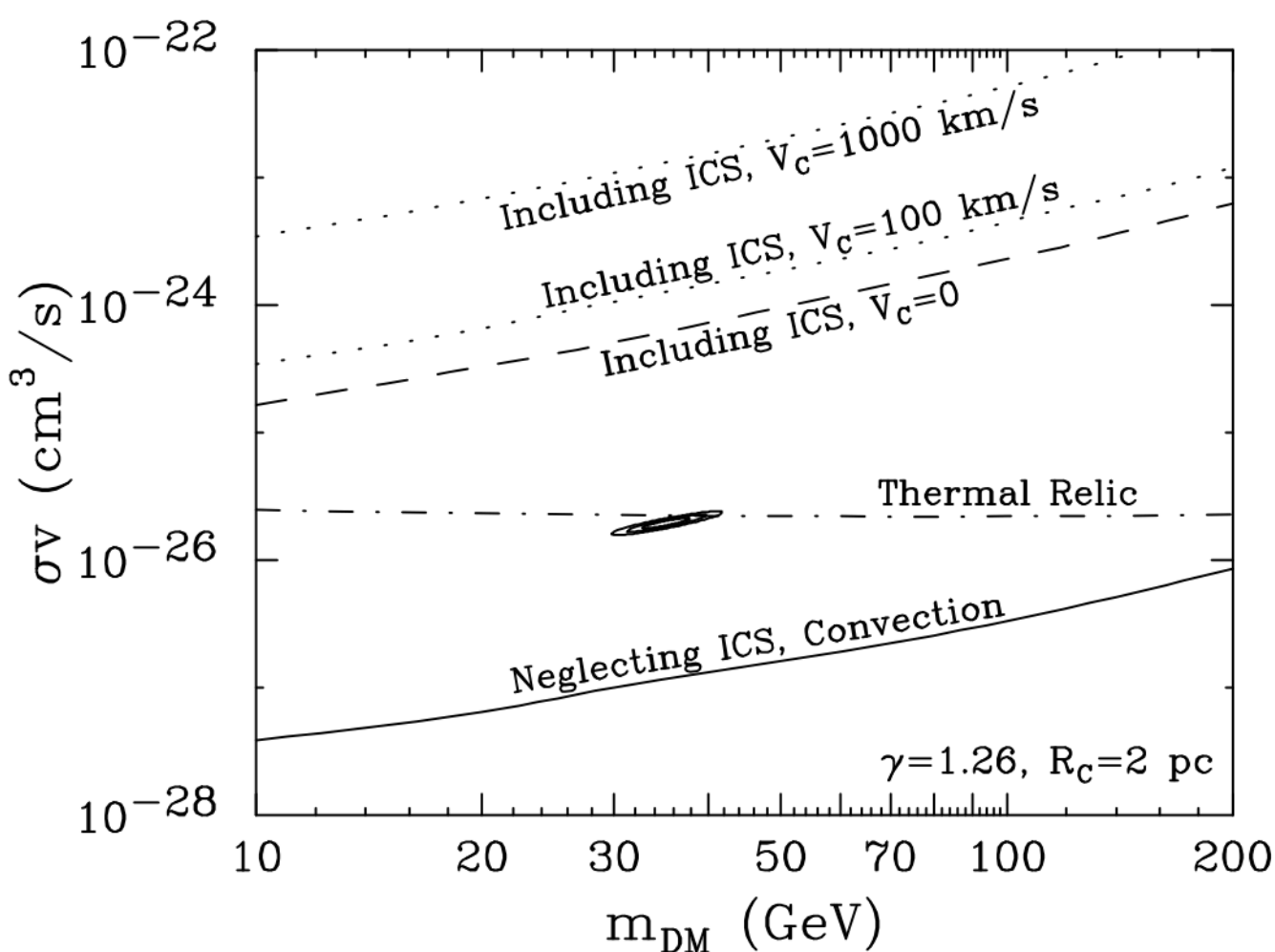
- Even in easy systems, mismodeling is a significant issue.

The Two Types of Dark Matter

- **Dark Matter signals that have low statistical significance.**
- **Dark Matter signals that might also be pulsars.**

Cosmic Microwave Background Limits

- **Very strong constraints on light dark matter.**
- **Cosmic-variance limited. Constraints will not improve much.**

Radio Detection Limits

- **Radio telescopes are very sensitive.**
- **Propagation of electrons near sources is highly uncertain.**

REVIEW

<https://doi.org/10.1038/s41586-018-0542-z>

A new era in the search for dark matter

Gianfranco Bertone^{1*} & Tim M. P. Tait^{1,2*}

There is a growing sense of ‘crisis’ in the dark-matter particle community, which arises from the absence of evidence for the most popular candidates for dark-matter particles—such as weakly interacting massive particles, axions and sterile neutrinos—despite the enormous effort that has gone into searching for these particles. Here we discuss what we have learned about the nature of dark matter from past experiments and the implications for planned dark-matter searches in the next decade. We argue that diversifying the experimental effort and incorporating astronomical surveys and gravitational-wave observations is our best hope of making progress on the dark-matter problem.

The fall of natural weakly interacting massive particles

The existence of dark matter has been discussed for more than a century^{1,2}. In the 1970s, astronomers and cosmologists began to build what is today a compelling body of evidence for this elusive component of the Universe, based on a variety of observations, including temperature anisotropies of the cosmic microwave background, baryonic acoustic oscillations, type Ia supernovae, gravitational lensing of galaxy clusters and rotation curves of galaxies^{3,4}. The standard model of particle physics contains no suitable particle to explain these observations, and

the observed Higgs mass at the weak scale appears highly unnatural, requiring an incredibly fine-tuned cancellation between the individually much larger intrinsic contribution and the correction terms, such that their sum is the value observed at the Large Hadron Collider (LHC). Natural theories introduce additional particles and symmetries, which are arranged so that these large corrections cancel each other out, protecting the Higgs mass from the influence of heavy mass scales.

The prototypical natural theory is the minimal supersymmetric (SUSY) standard model, which introduces an additional partner for

GeV-Scale Thermal WIMPs: Not Even Slightly Dead

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²*Center for Cosmology and AstroParticle Physics (CCAPP),*

Ohio State University, Columbus, OH 43210, USA

³*Department of Physics, Ohio State University, Columbus, OH 43210, USA*

⁴*Department of Astronomy, Ohio State University, Columbus, OH 43210, USA*

⁵*Department of Particle Physics and Astrophysics,*

Weizmann Institute of Science, Rehovot 76100, Israel

(Dated: July 13, 2018)

Weakly Interacting Massive Particles (WIMPs) have long reigned as one of the leading classes of dark matter candidates. The observed dark matter abundance can be naturally obtained by freeze-out of weak-scale dark matter annihilations in the early universe. This “thermal WIMP” scenario makes direct predictions for the total annihilation cross section that can be tested in present-day experiments. While the dark matter mass constraint can be as high as $m_\chi \gtrsim 100$ GeV for particular annihilation channels, the constraint on the *total* cross section has not been determined. We construct the first model-independent limit on the WIMP total annihilation cross section, showing that allowed combinations of the annihilation-channel branching ratios considerably weaken the sensitivity. For thermal WIMPs with s -wave $2 \rightarrow 2$ annihilation to visible final states, we find the dark matter mass is only known to be $m_\chi \gtrsim 20$ GeV. This is the strongest largely model-independent lower limit on the mass of thermal-relic WIMPs; together with the upper limit on the mass from the unitarity bound ($m_\chi \lesssim 100$ TeV), it defines what we call the “WIMP window”. To probe the remaining mass range, we outline ways forward.

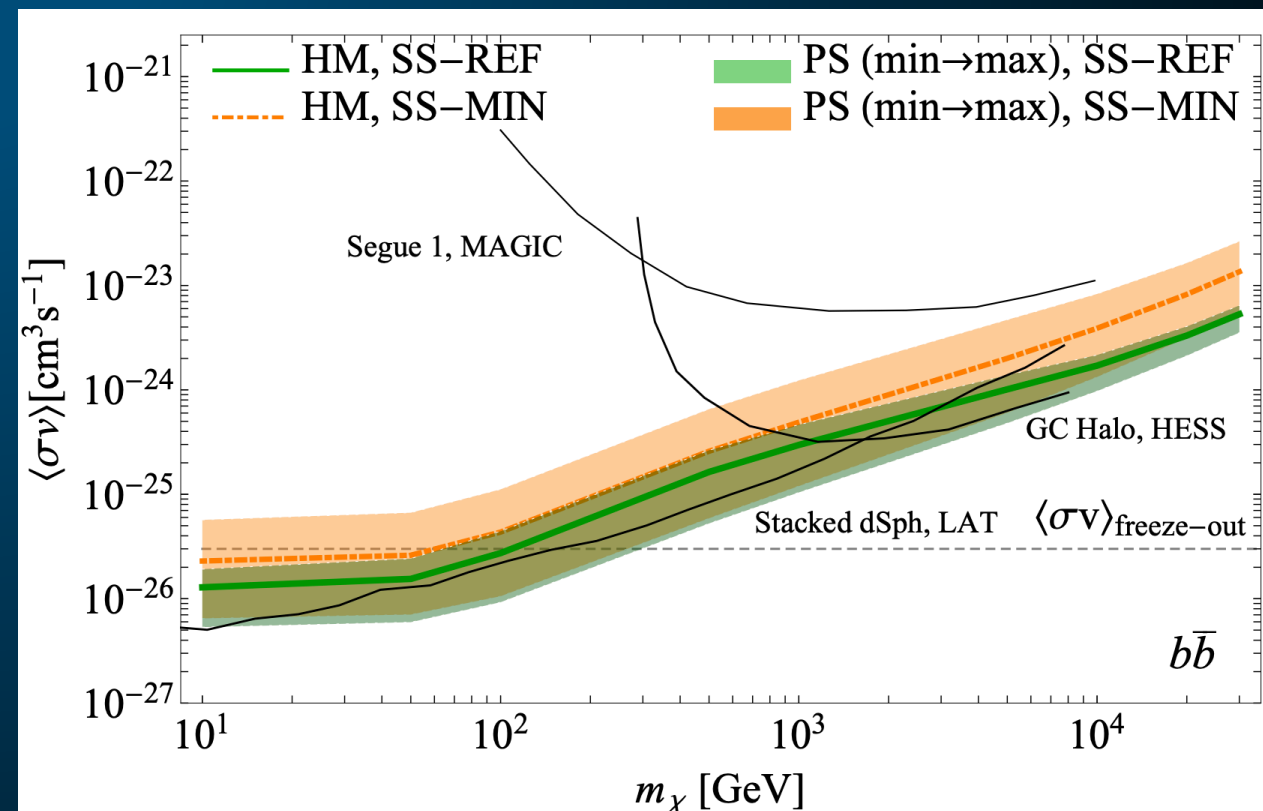
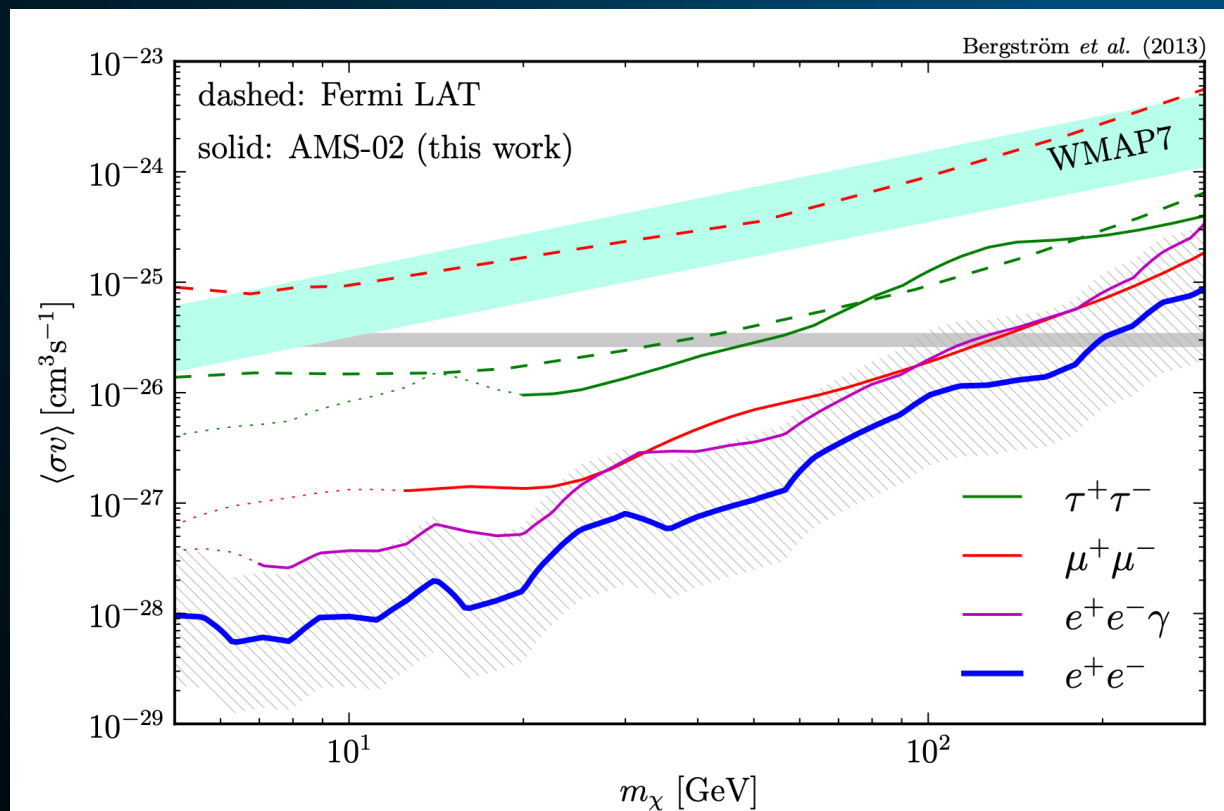
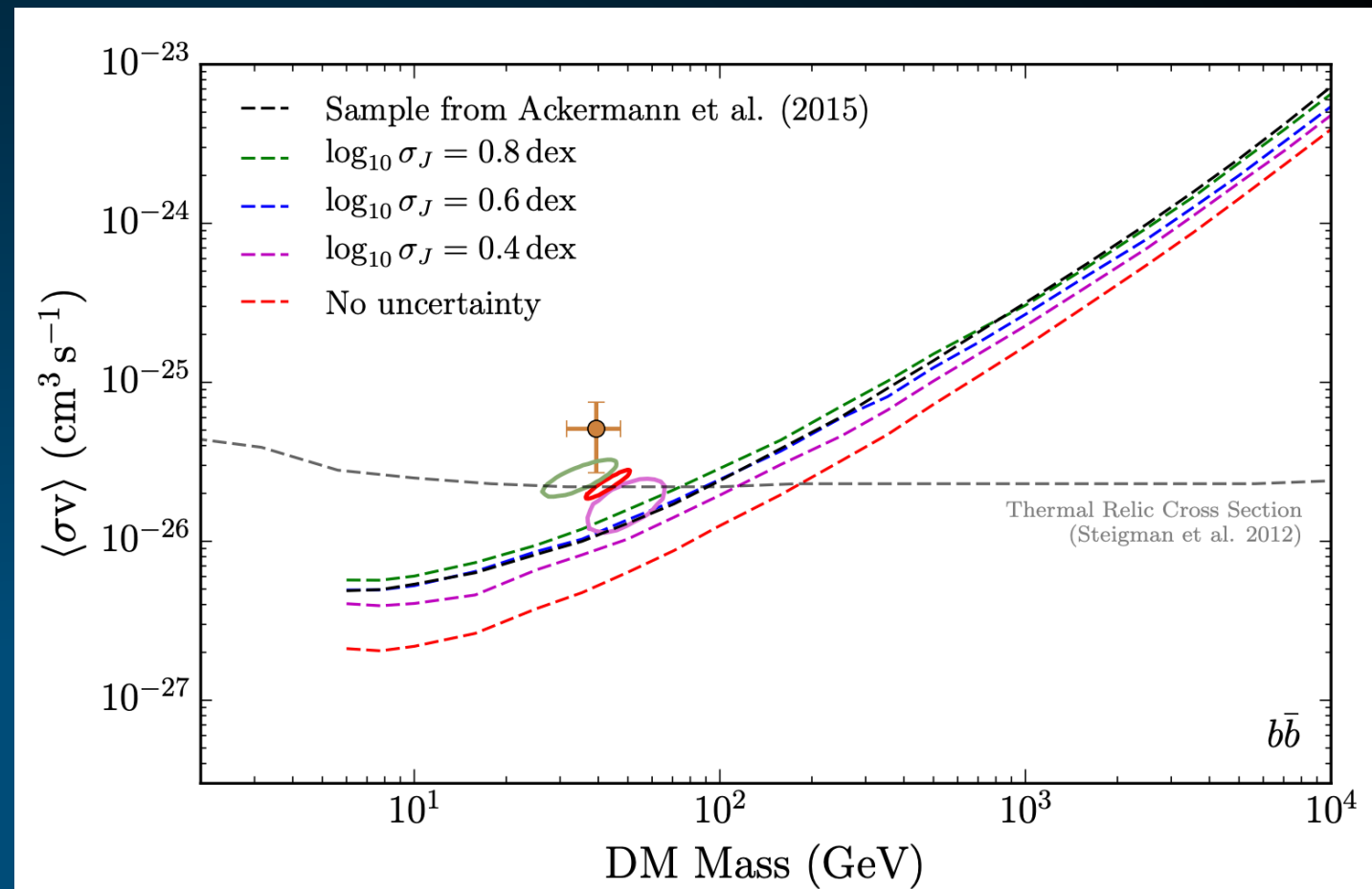
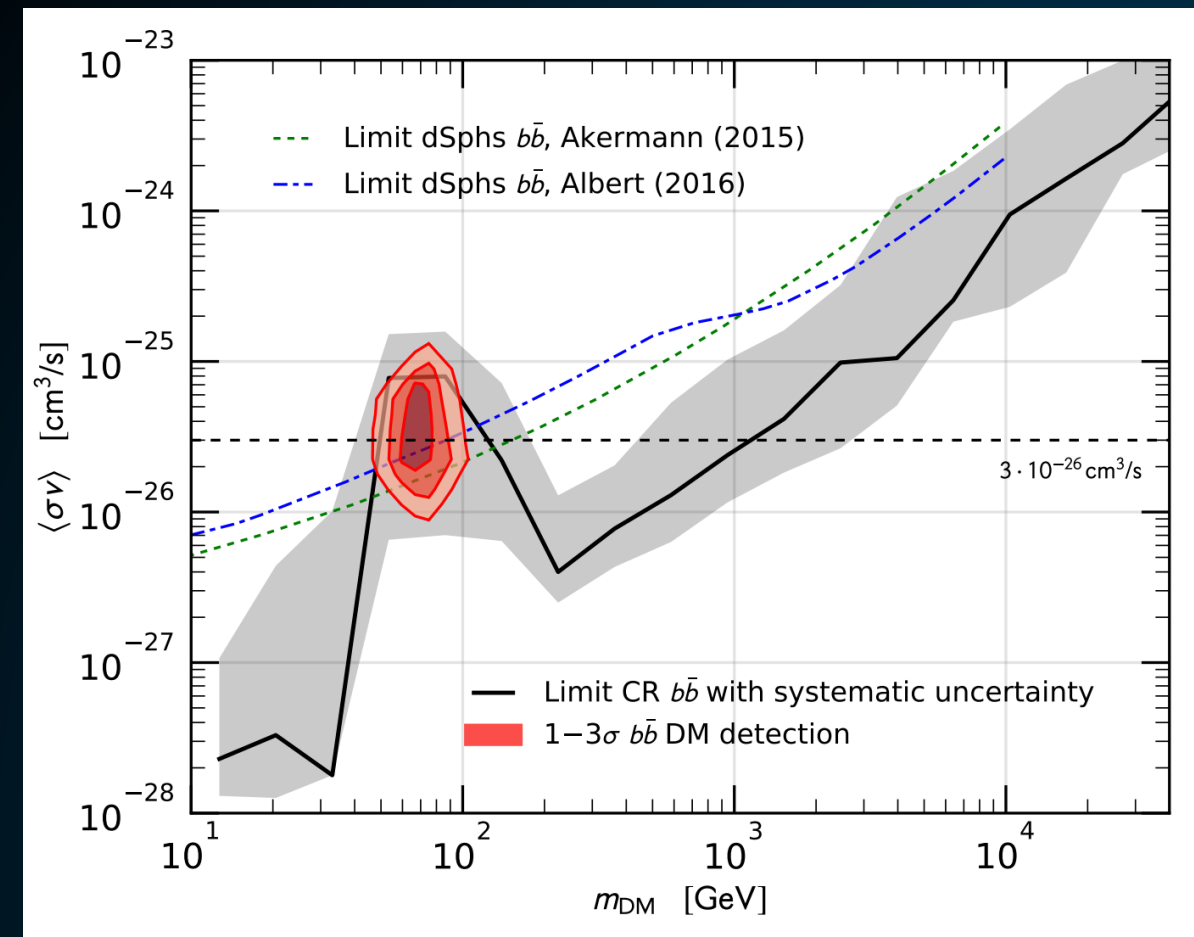
I. INTRODUCTION

A leading candidate for dark matter (DM) is a Weakly Interacting Massive Particle (WIMP) that is a thermal

scenarios. The branching ratios, coupling types and signals are model-dependent, and so the lack of observations may just be due to such features. For example, there can be interference effects, momentum suppression, or

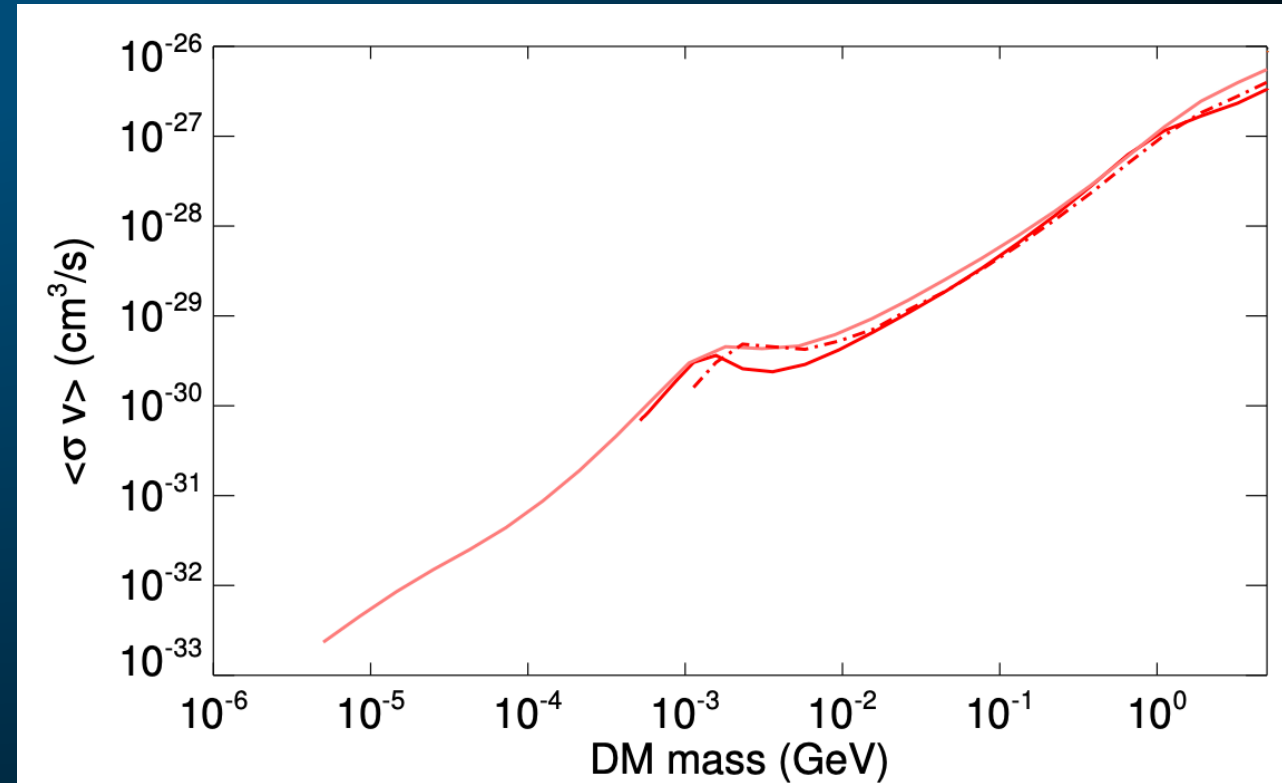
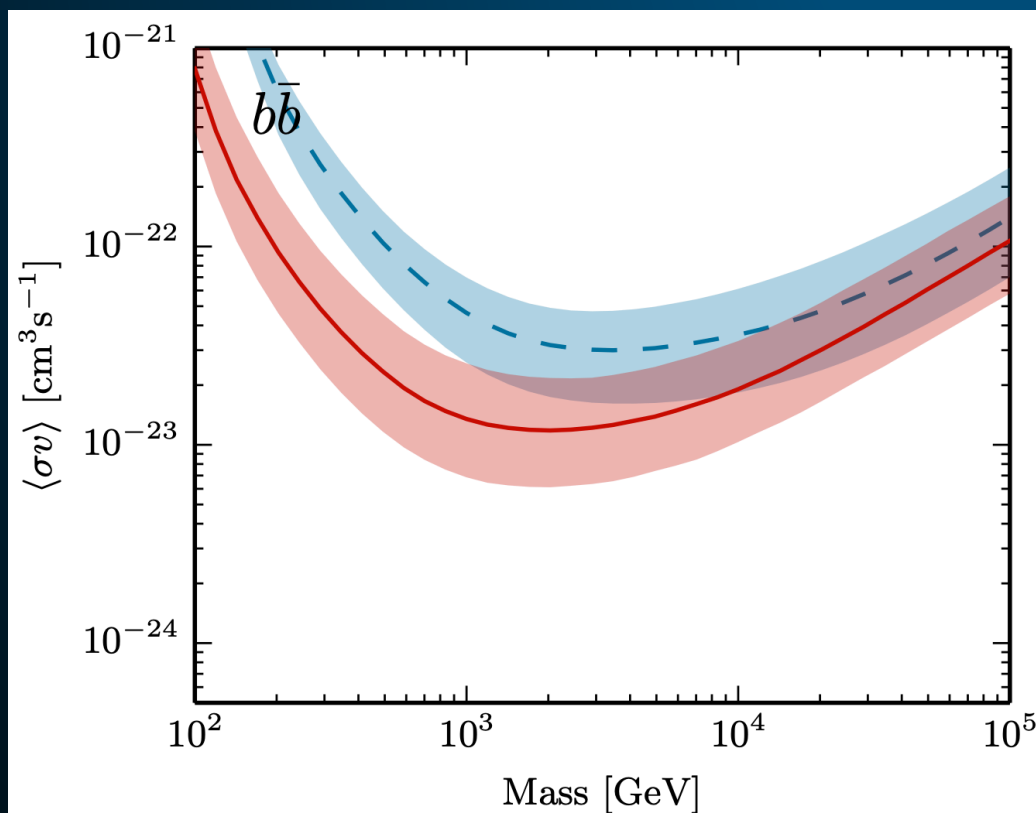
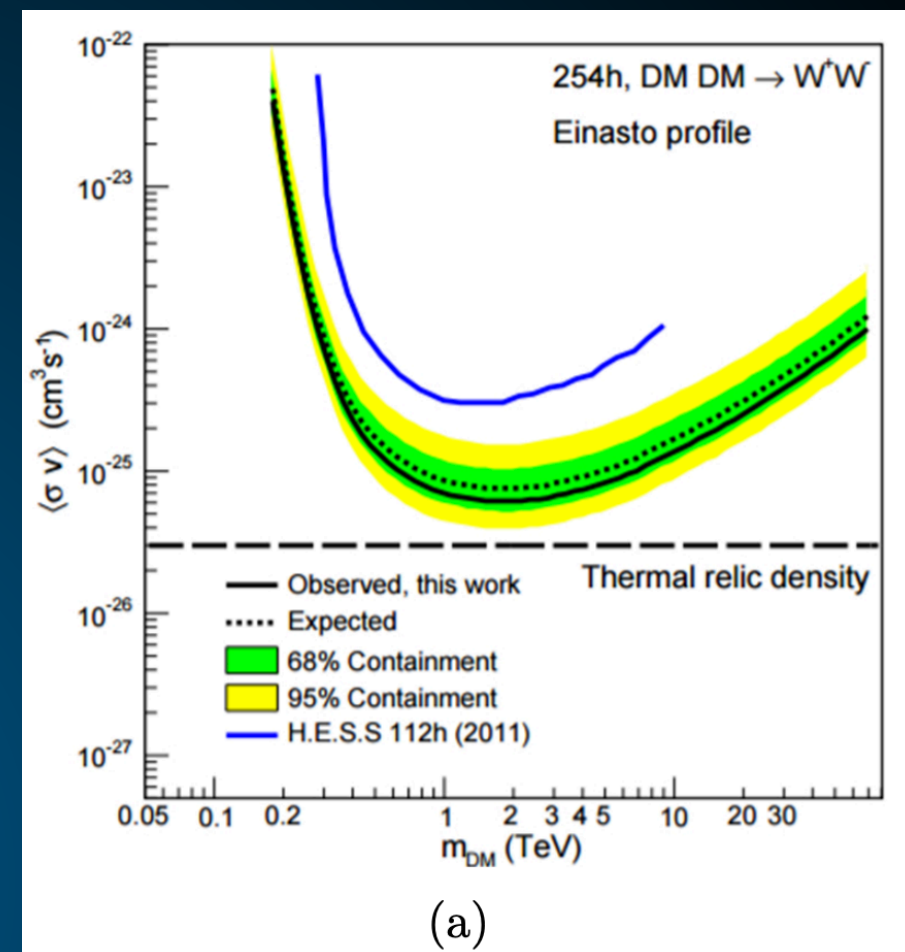
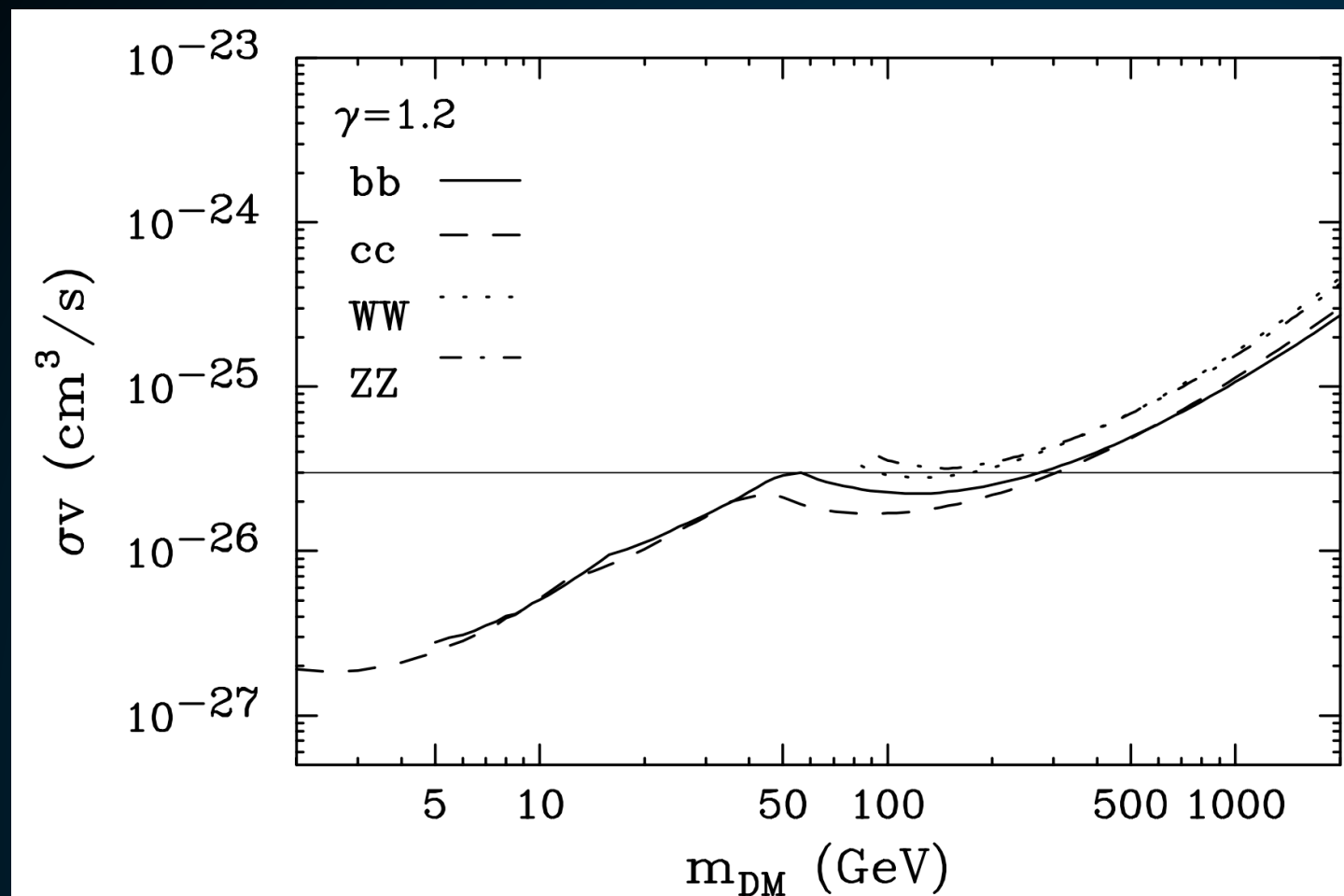
Looking Forward

How to End a Talk When You Don't Have a Conclusion



Looking Forward

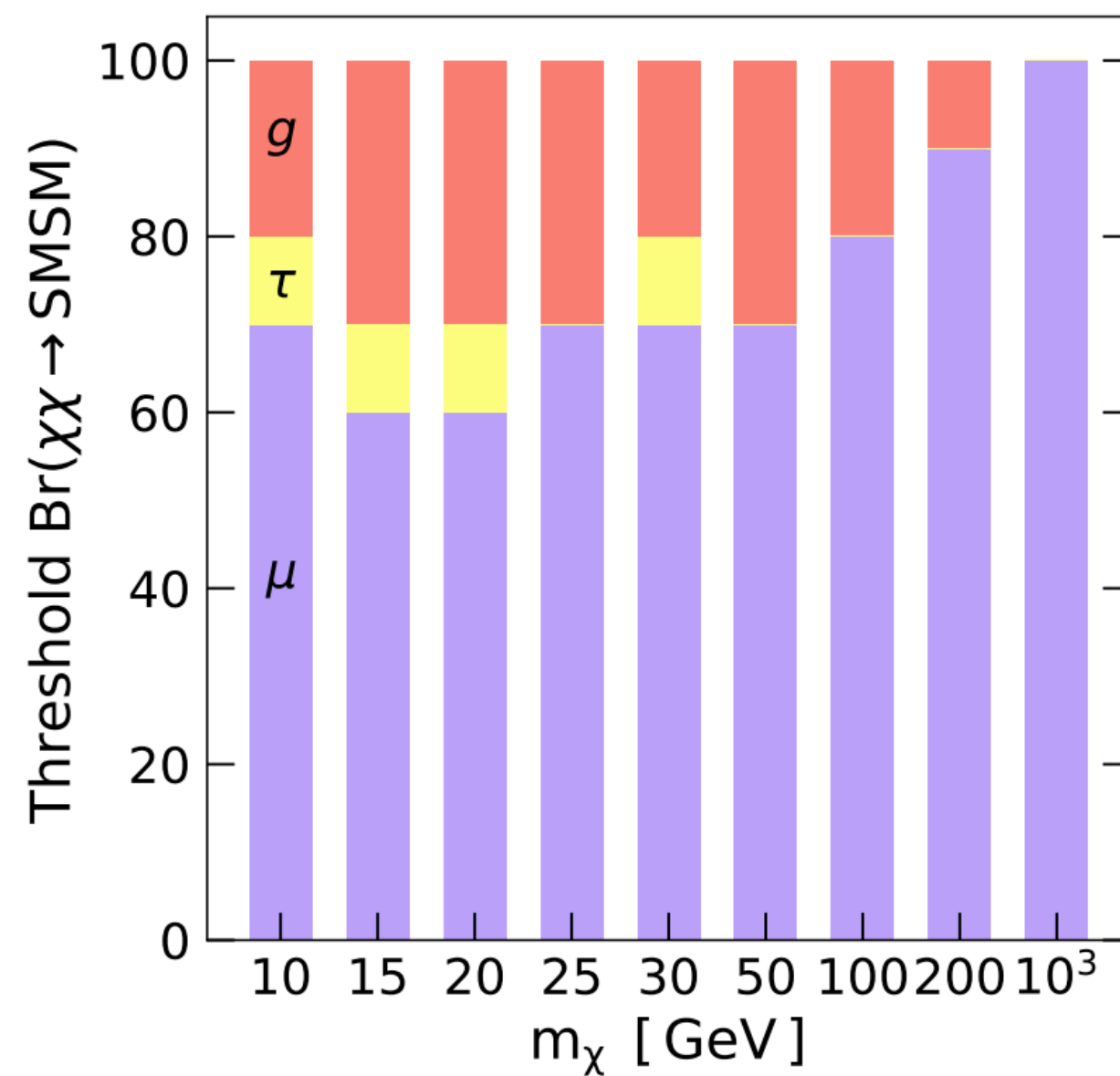
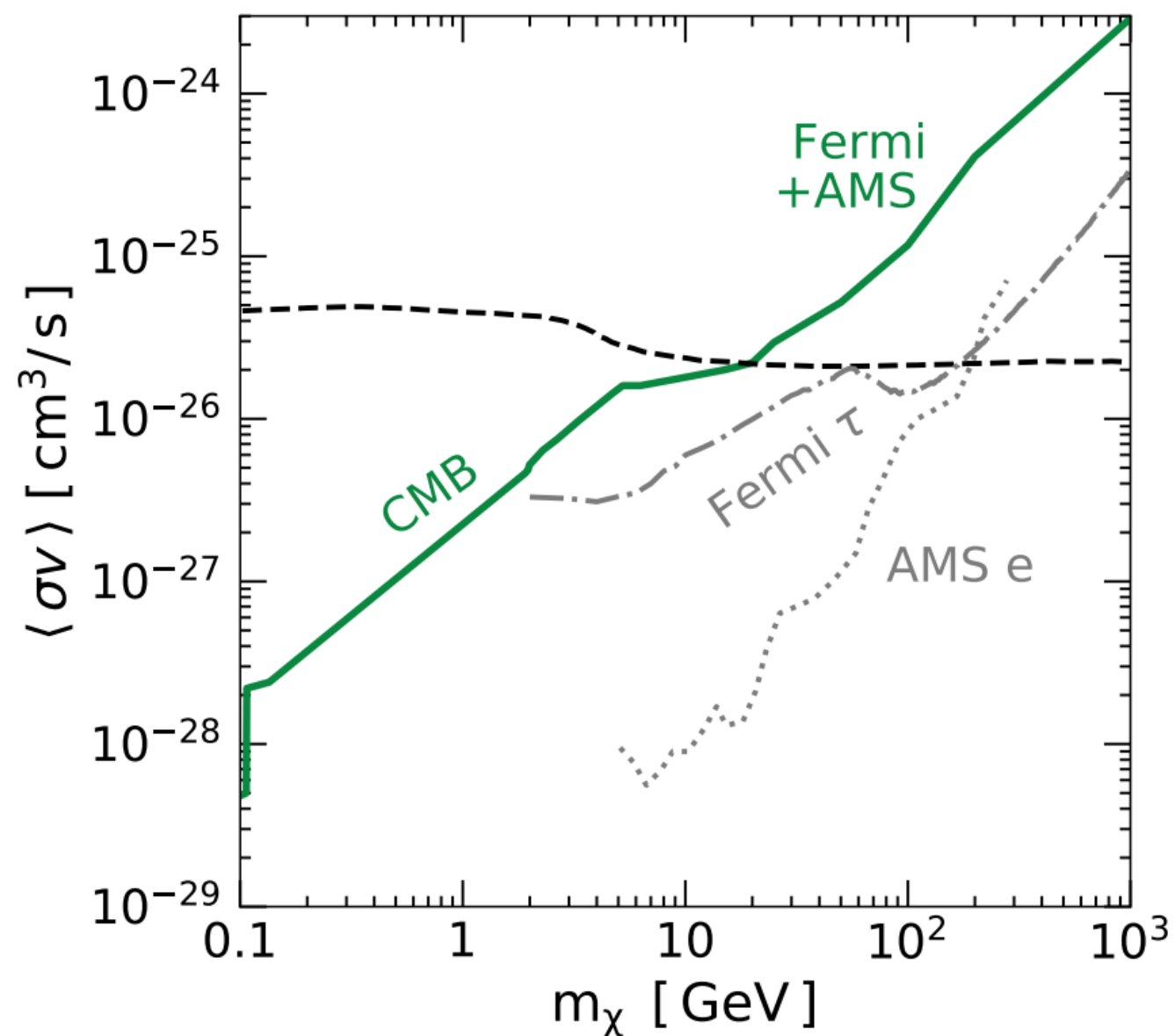
How to End a Talk When You Don't Have a Conclusion



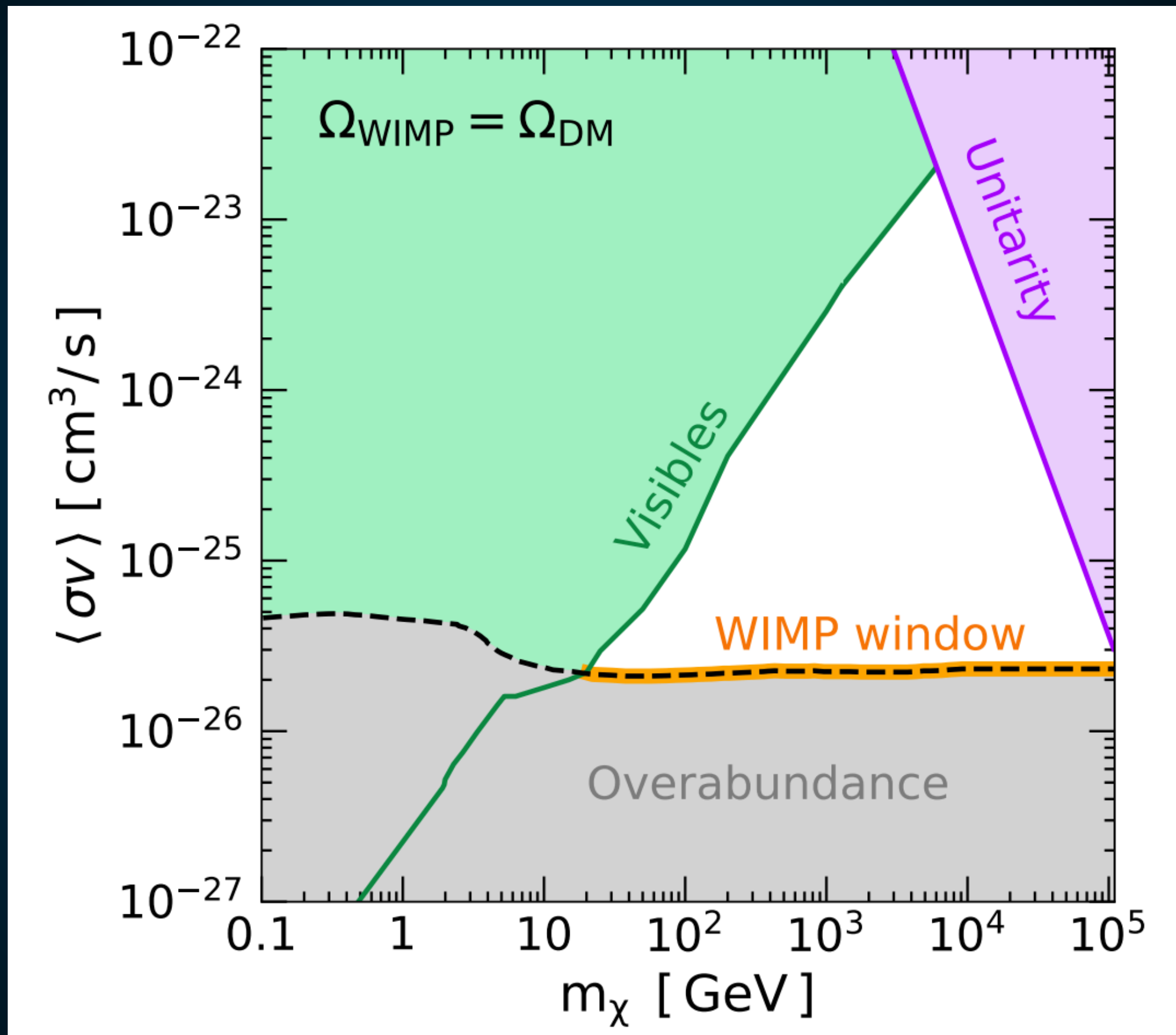
Looking Forward

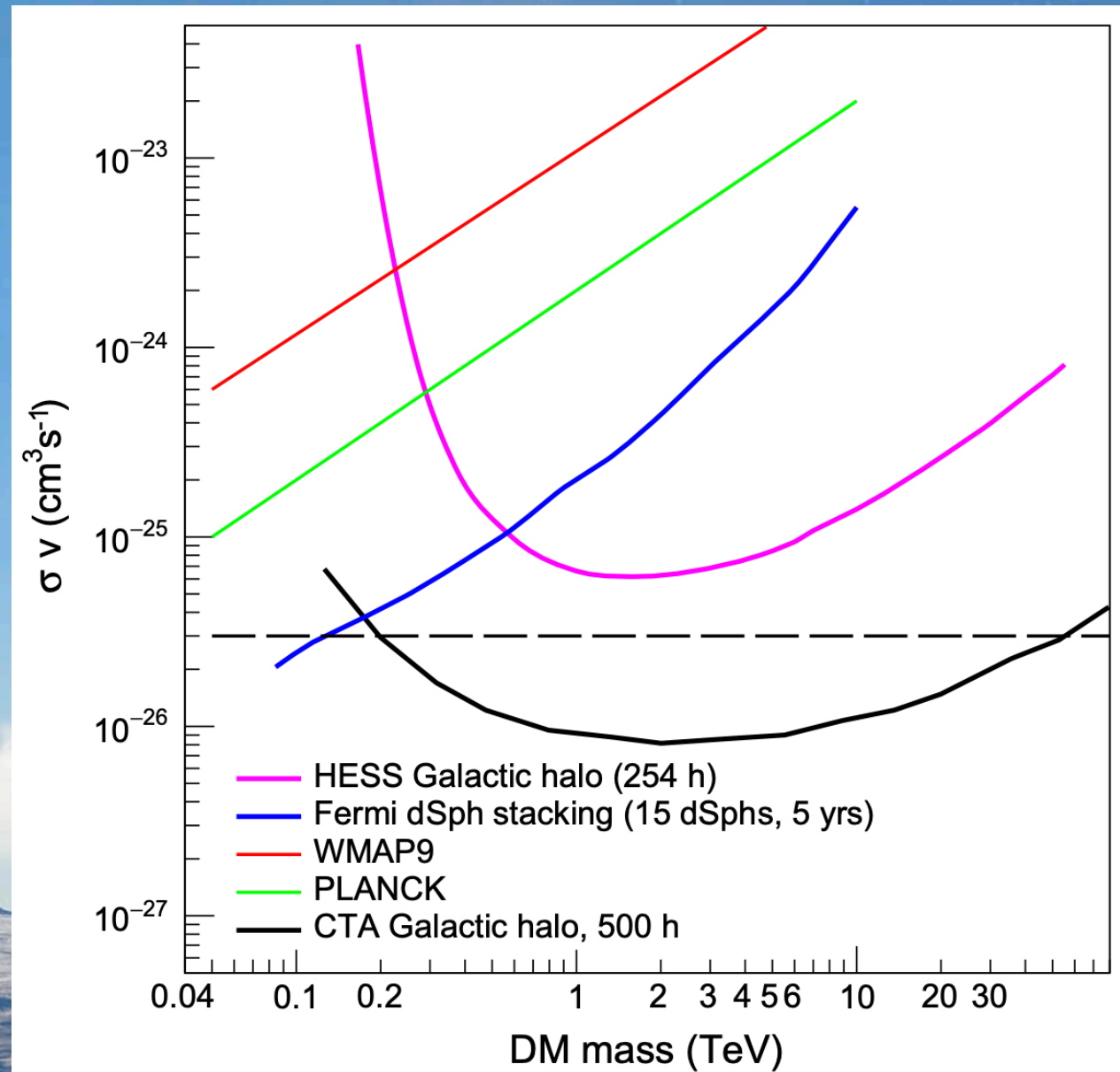
Leane et al. (2018; 1805.10305)

How to End a Talk When You Don't Have a Conclusion



How to End a Talk When You Don't Have a Conclusion





AMEGO

ALL-SKY MEDIUM ENERGY GAMMA-RAY OBSERVATORY



Polarization

Spectroscopy

511 keV
Galactic Map

Supernovae

Nearby
Galaxies

Galactic Diffuse
Emission

Supernova Remnants

Novae

Compact
Object
Binaries

Gravitational
Wave
Counterparts

Dark Matter

Pulsars

Pulsar Wind
Nebulae

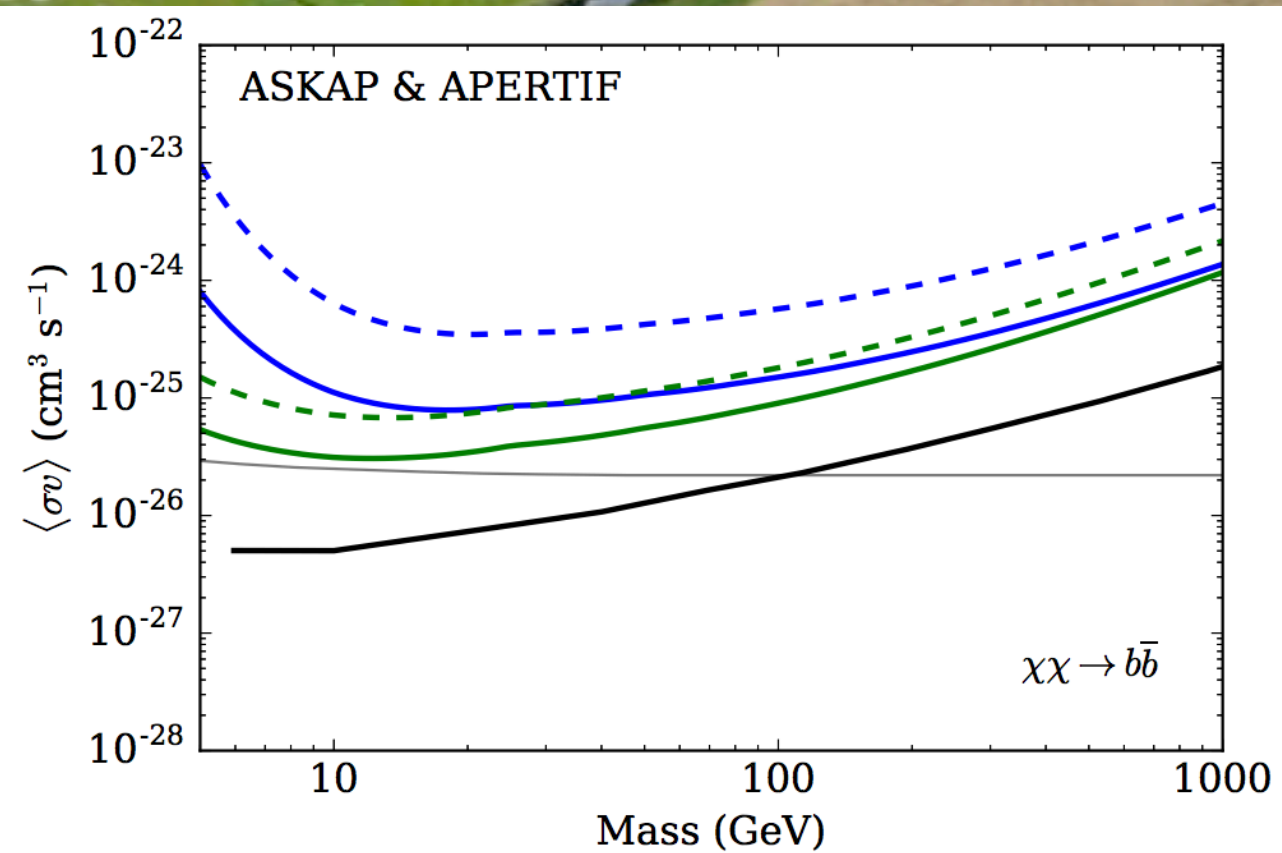
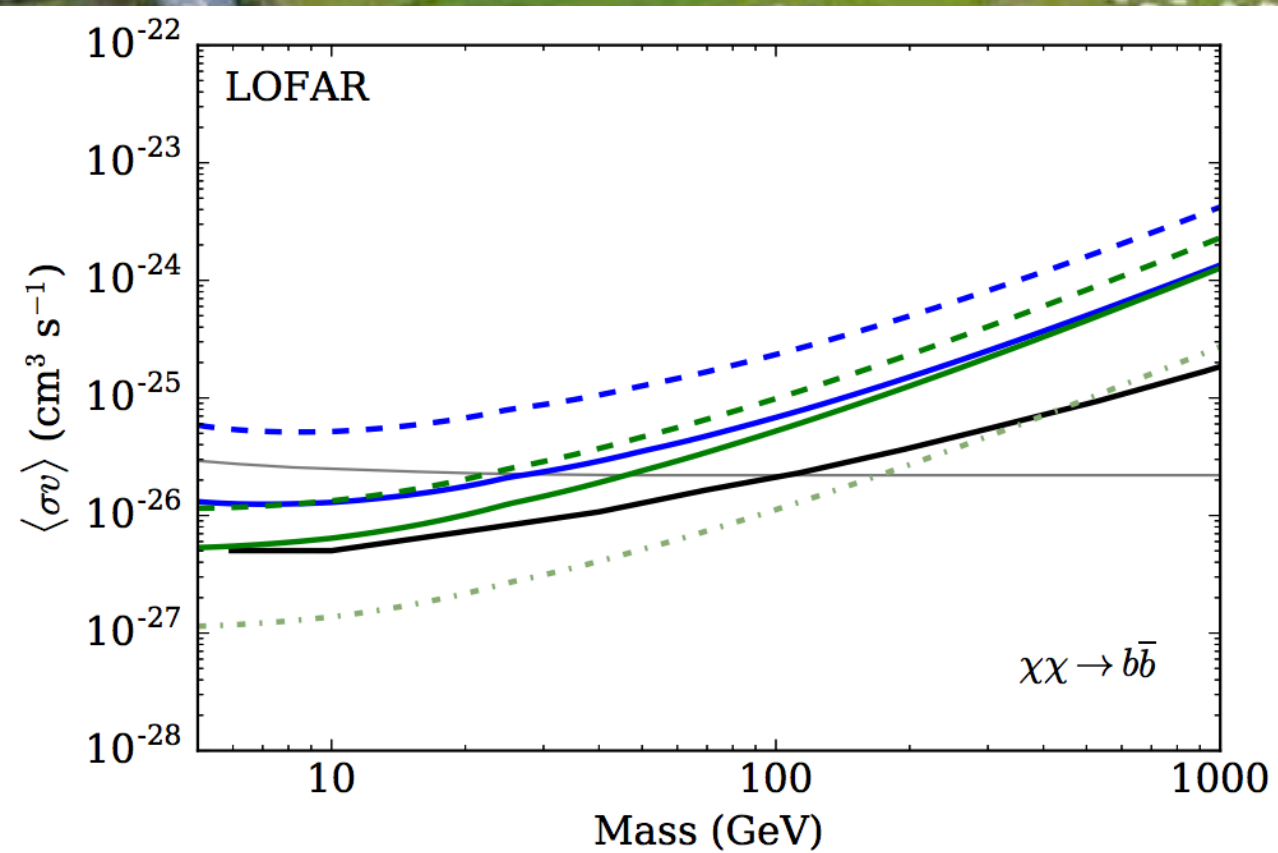
Magnetars

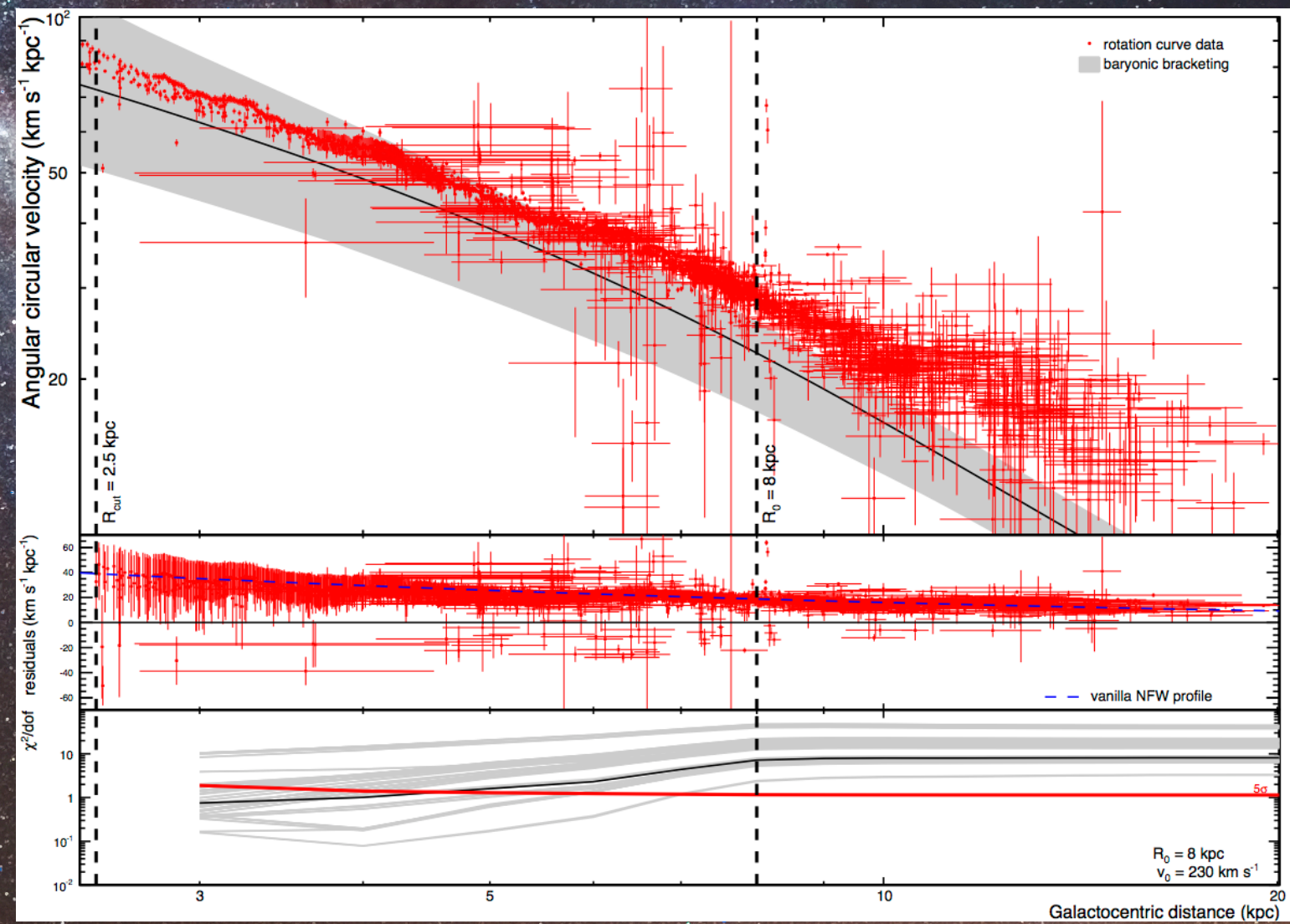
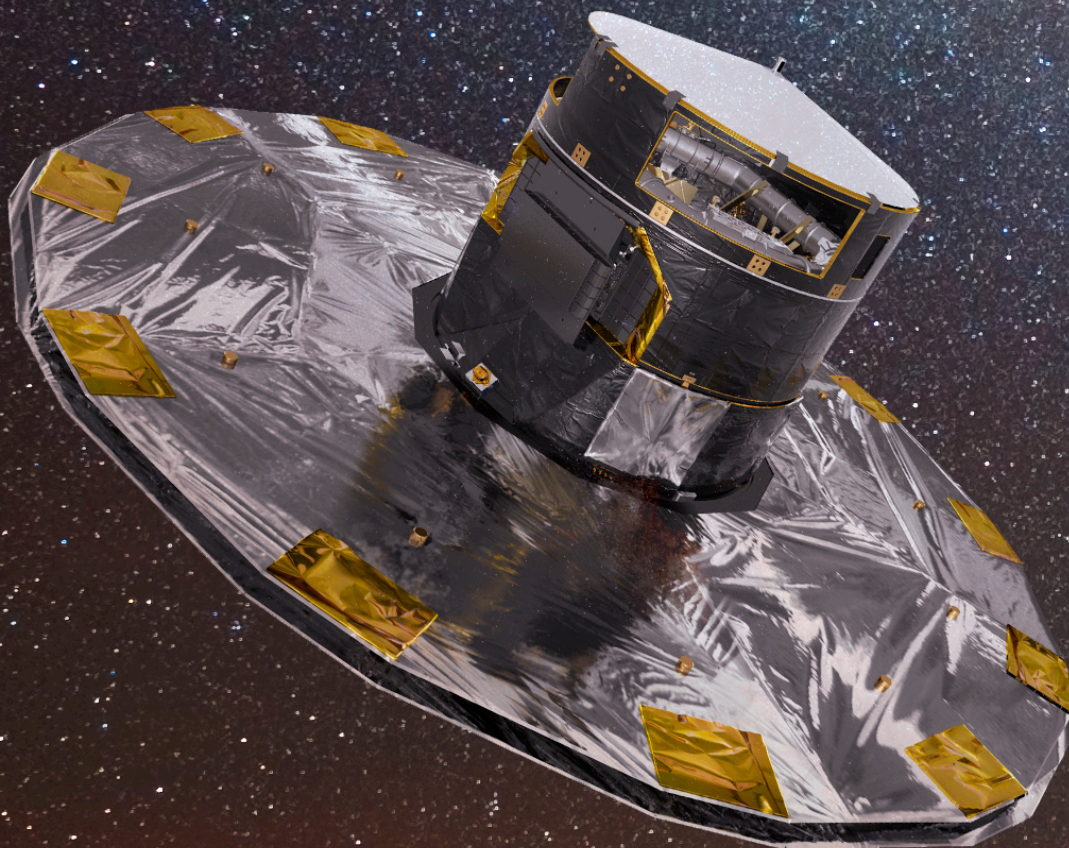
Gamma-ray
Bursts

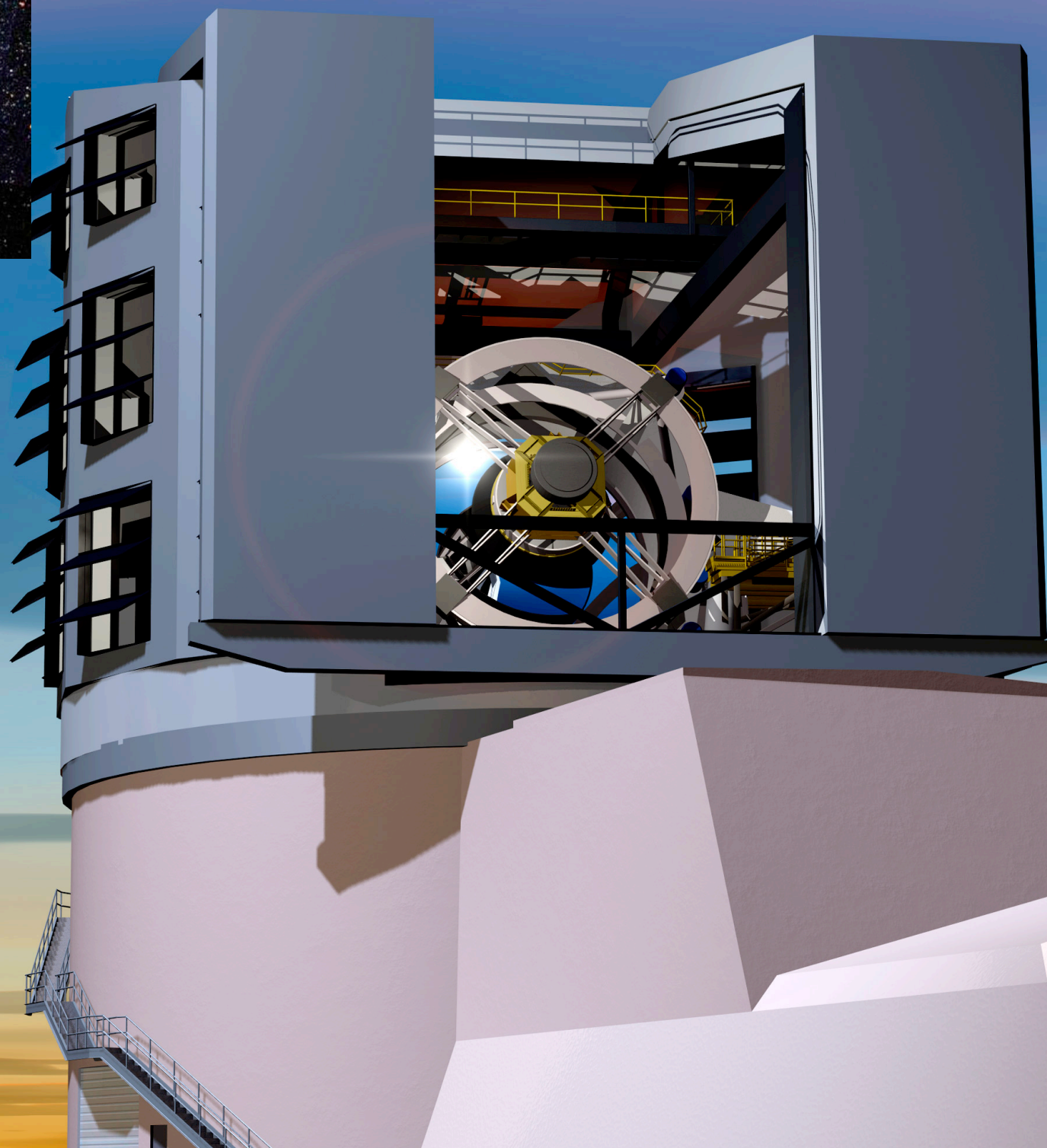
Active Galactic
Nuclei

Neutrino
Counterparts

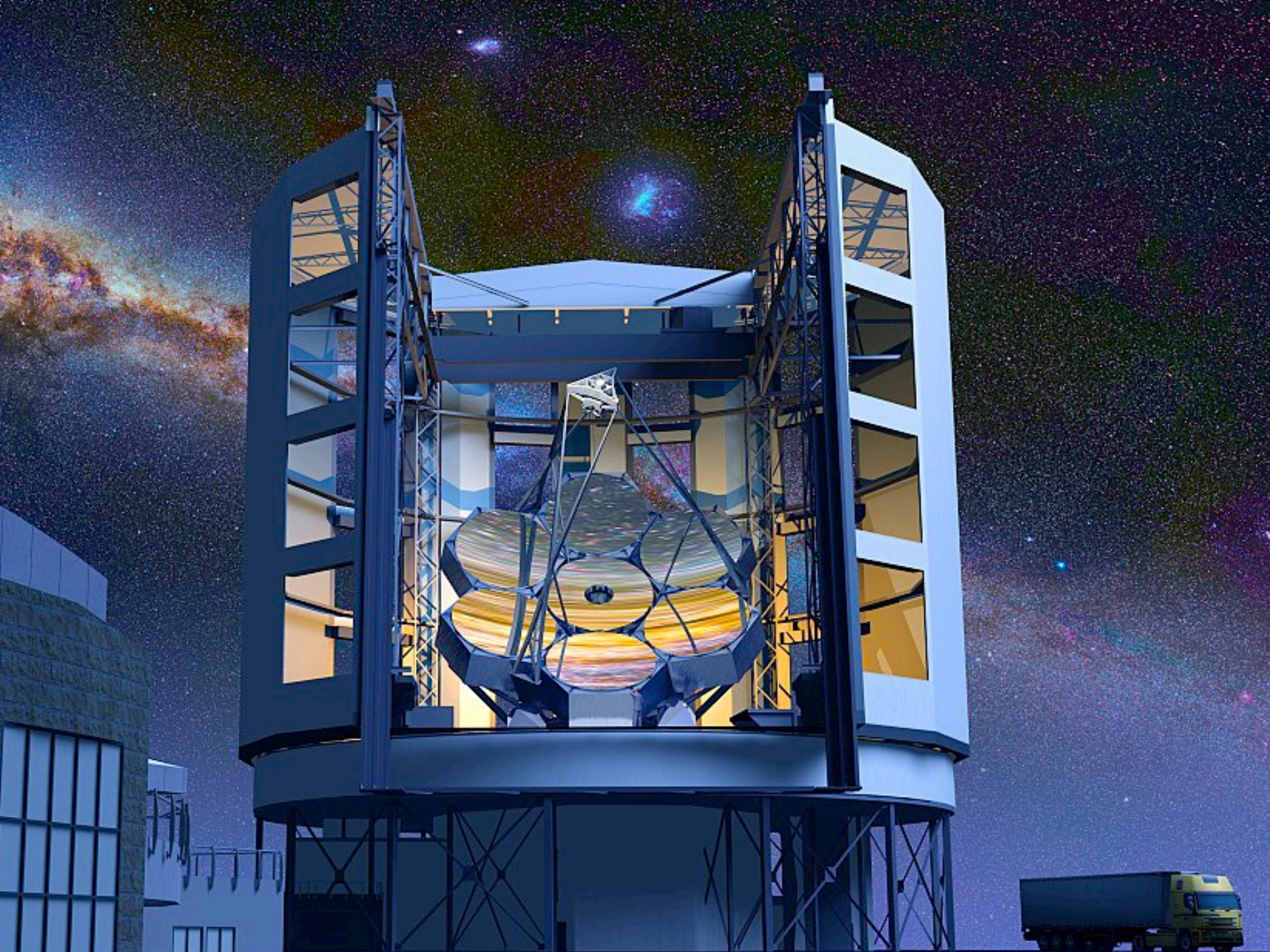
Continuum

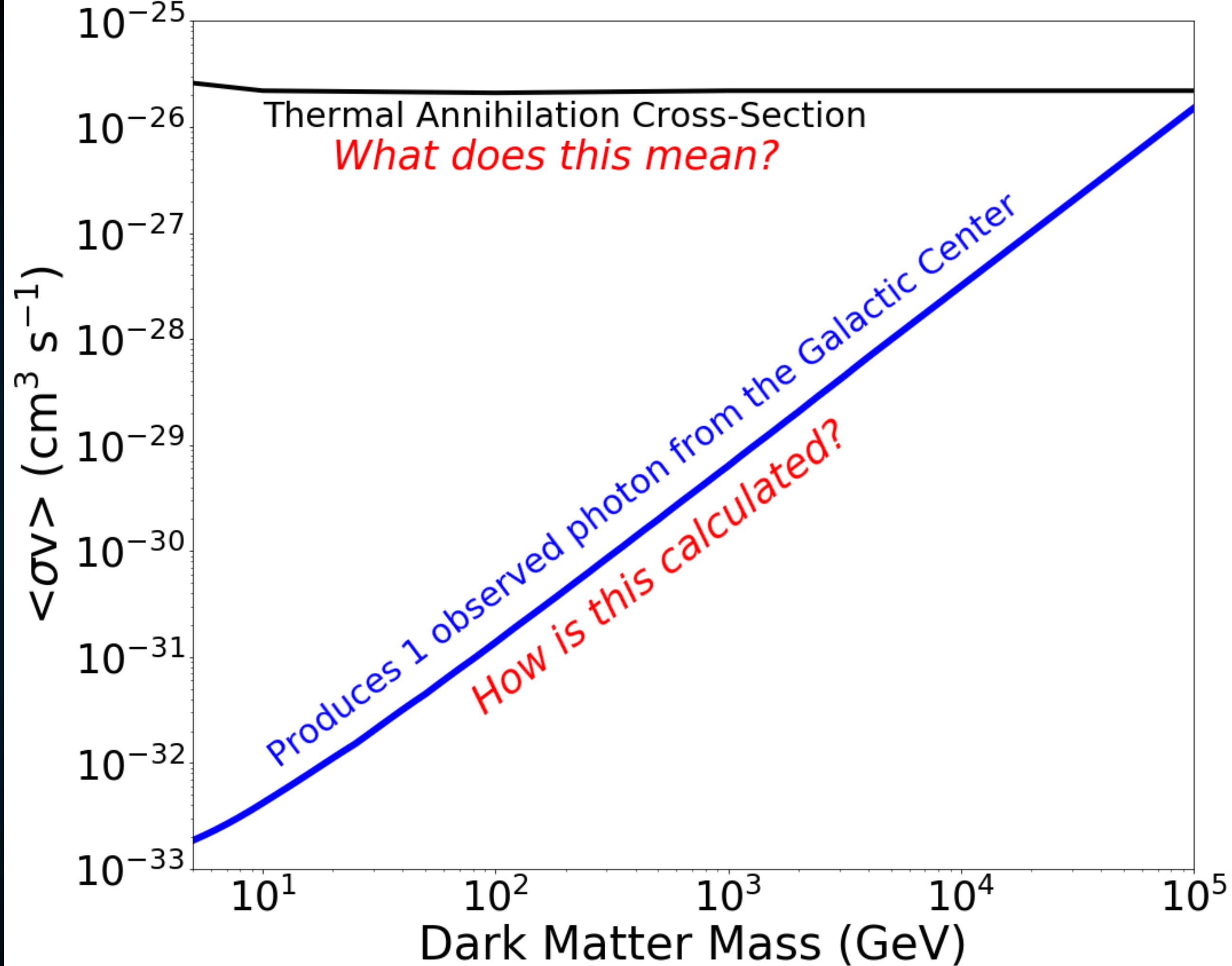






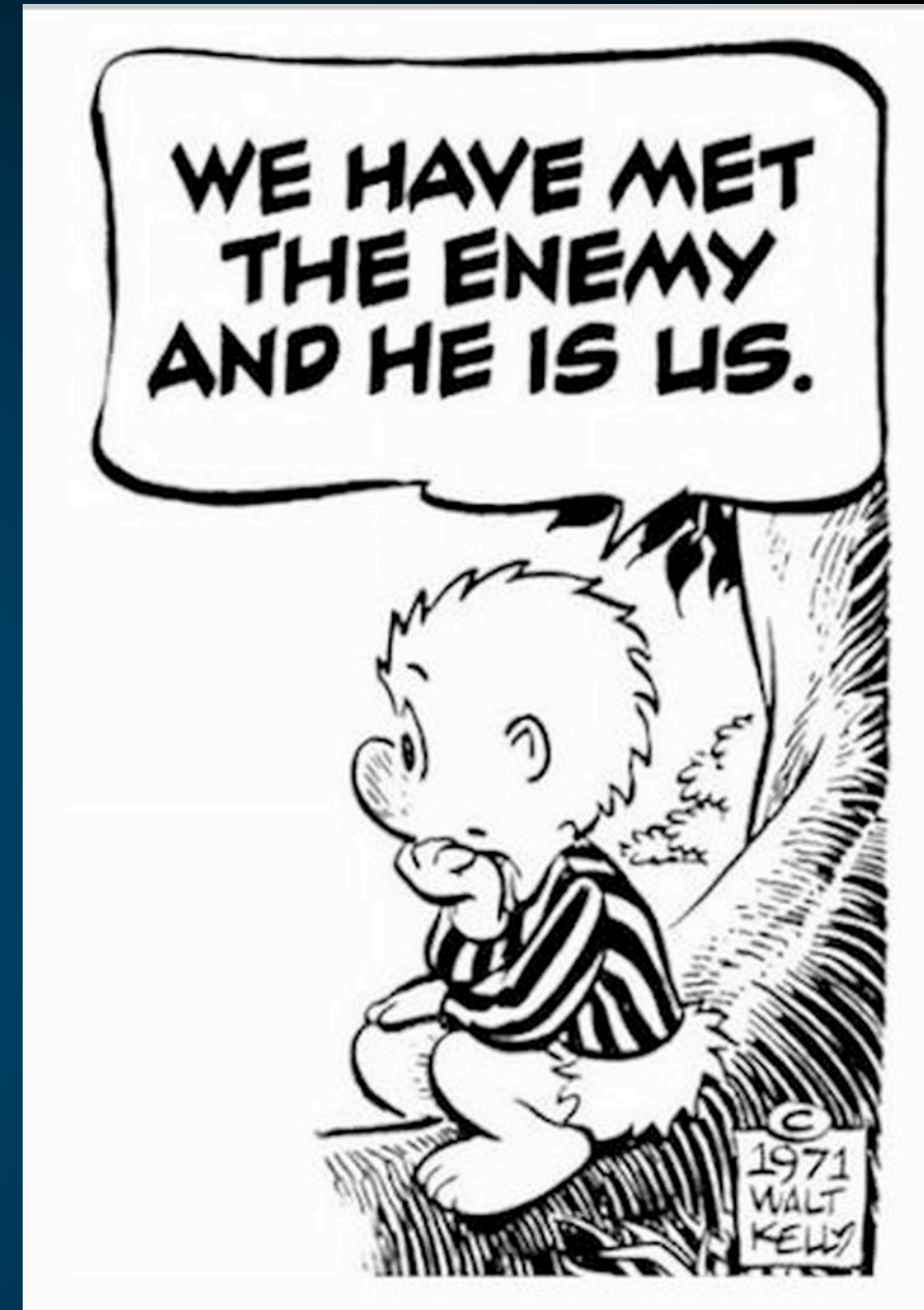
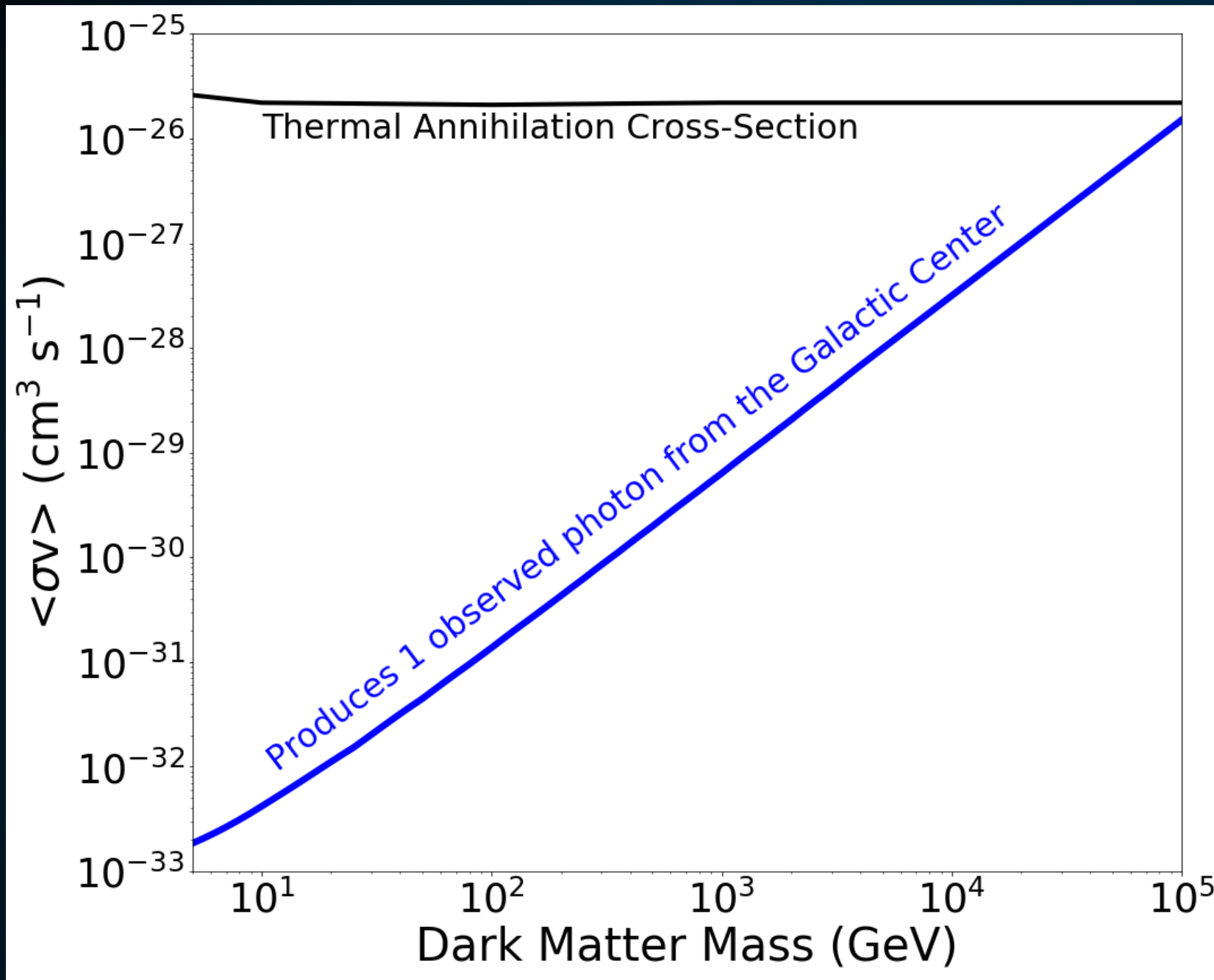






Looking Forward

How to End a Talk When You Don't Have a Conclusion



Origin of the Cosmic Ray Galactic Halo Driven by Advected Turbulence and Self-Generated Waves

Carmelo Evoli,^{1,2,*} Pasquale Blasi,^{1,2,3,†} Giovanni Morlino,^{1,2,3,‡} and Roberto Aloisio^{1,2,§}

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²*INFN/Laboratori Nazionali del Gran Sasso, Via G. Acitelli 22, Assergi (AQ), Italy*

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The diffusive paradigm for the transport of Galactic cosmic rays is central to our understanding of the origin of these high energy particles. However it is worth recalling that the normalization, energy dependence and spatial extent of the diffusion coefficient in the interstellar medium are fitted to the data and typically are not derived from more basic principles. Here we discuss a scenario in which the diffusion properties of cosmic rays are derived from a combination of wave self-generation and advection from the Galactic disc, where the sources of cosmic rays are assumed to be located. We show for the first time that a halo naturally arises from these phenomena, with a size of a few kpc, compatible with the value that typically best fits observations in simple parametric approaches to cosmic ray diffusion. We also show that transport in such a halo results in a hardening in the spectra of primary cosmic rays at ~ 300 GV.

Introduction – Understanding cosmic-ray (CR) propagation in the Galaxy and its implications for observations at different energies and with different messengers is one of the challenges of modern astroparticle physics.

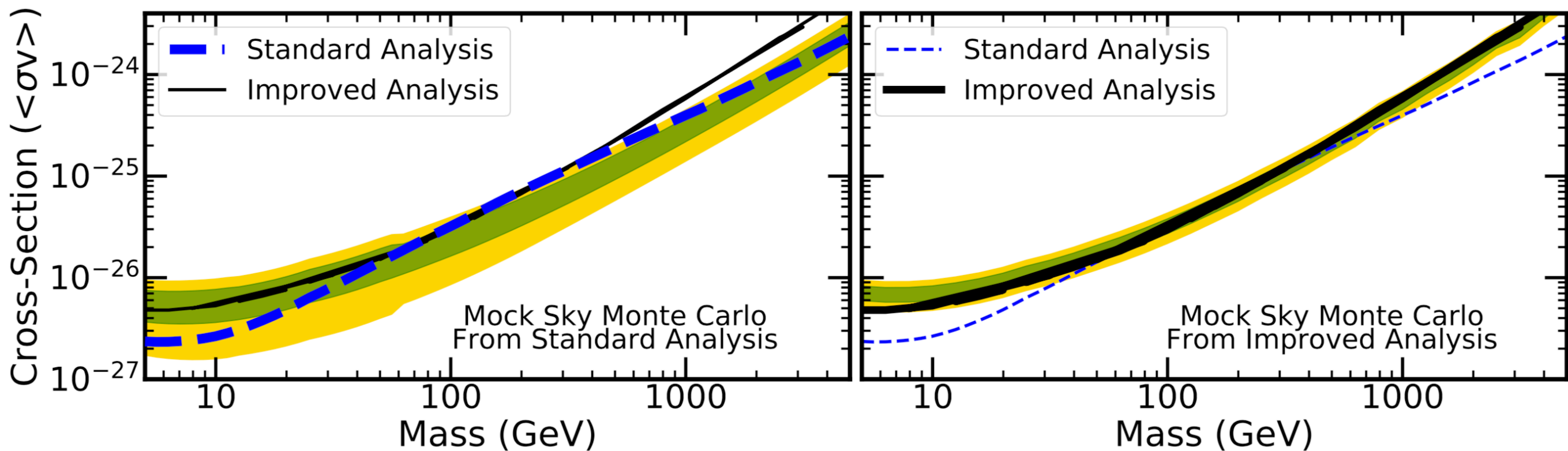
The standard scenario adopted to describe Galactic propagation in terms of properties of the interstellar turbulence is the so called *galactic halo model* proposed by Ginzburg and Syrovatskii in 1964 [1] and described in detail in [2]. The halo model is usually implemented assuming that CRs are produced by sources located in the thin Galactic disc and then diffuse by scattering off random magnetic fluctuations in a low-density confinement region (“halo”) extending well be-

grammage that scales with rigidity as $R^{-1/3}$, that is claimed to be consistent with the diffusion coefficient expected from transport in a turbulence with Kolmogorov spectrum, $D(R) \sim 10^{28} (R/\text{GV})^{1/3} \text{ cm}^2 \text{ s}^{-1}$.

An independent piece of evidence of the existence of a magnetized halo comes from observations in the radio band of diffuse synchrotron emission, revealing the presence of electrons and magnetic fields above and below the Galactic plane [6]. The existence of a halo of several kpc size can be inferred from a comparison between numerical models for the CR electron distribution and the morphology of the radio emission [7, 8]. It is worth mentioning that radio halos with a

Looking Forward

How to End a Talk When You Don't Have a Conclusion



Accounting for Astrophysical Mismodeling Significantly Affects Dark Matter Annihilation Constraints from Dwarf Spheroidal Galaxies

Tim Linden*

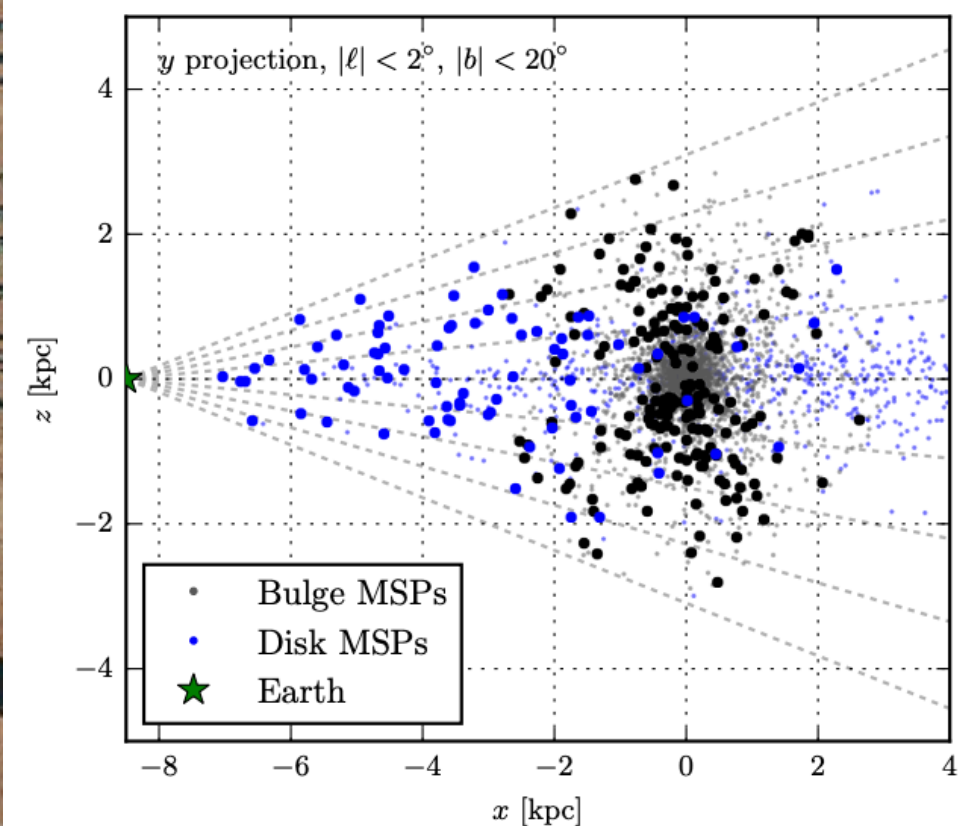
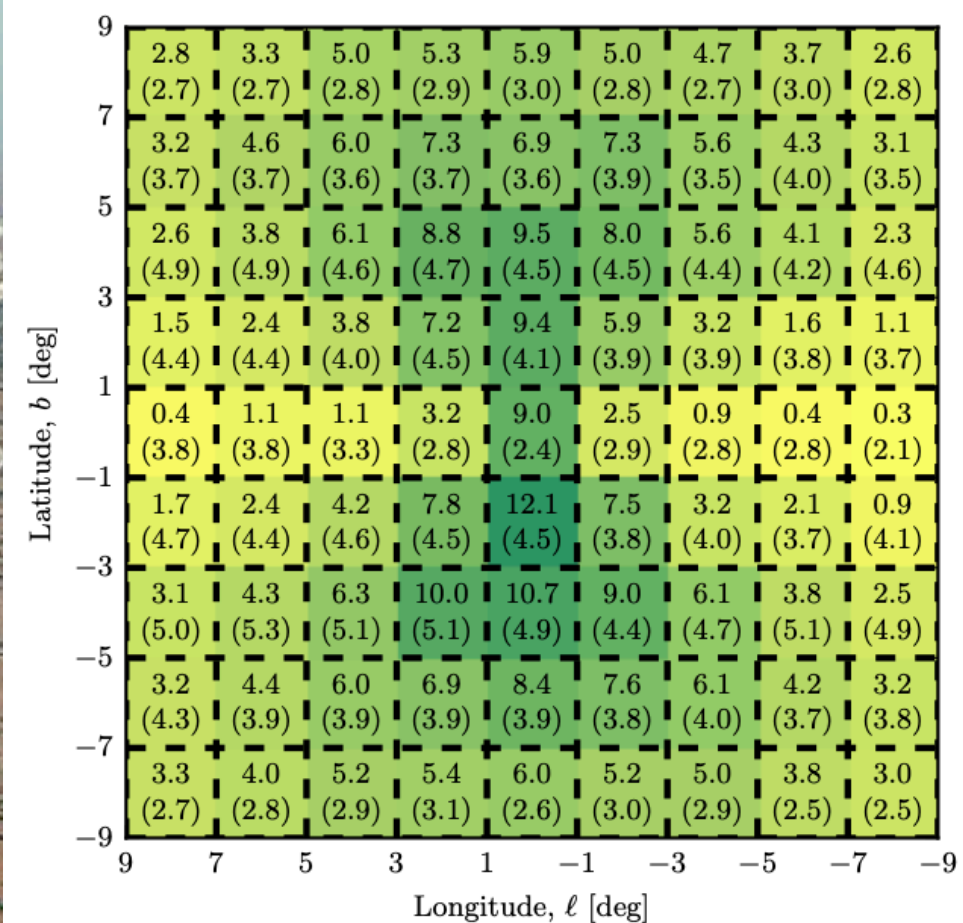
*Center for Cosmology and AstroParticle Physics (CCAPP), and
Department of Physics, The Ohio State University Columbus, OH, 43210*

Fermi-LAT observations have strongly constrained dark matter annihilation through the joint-likelihood analysis of dwarf spheroidal galaxies (dSphs). These constraints are expected to be robust because dSphs have measurable dark matter content and produce negligible astrophysical emission. However, each dSph is dim, with a typical flux that falls below the accuracy of the background model. We show that this significantly diminishes the reliability of previous joint-likelihood algorithms, and develop an improved analysis that directly accounts for the effect of background mismodeling. This method produces more robust limits and detections of dark matter in both real and mock data. We calculate improved limits on the dark matter annihilation cross-section,

Looking Forward

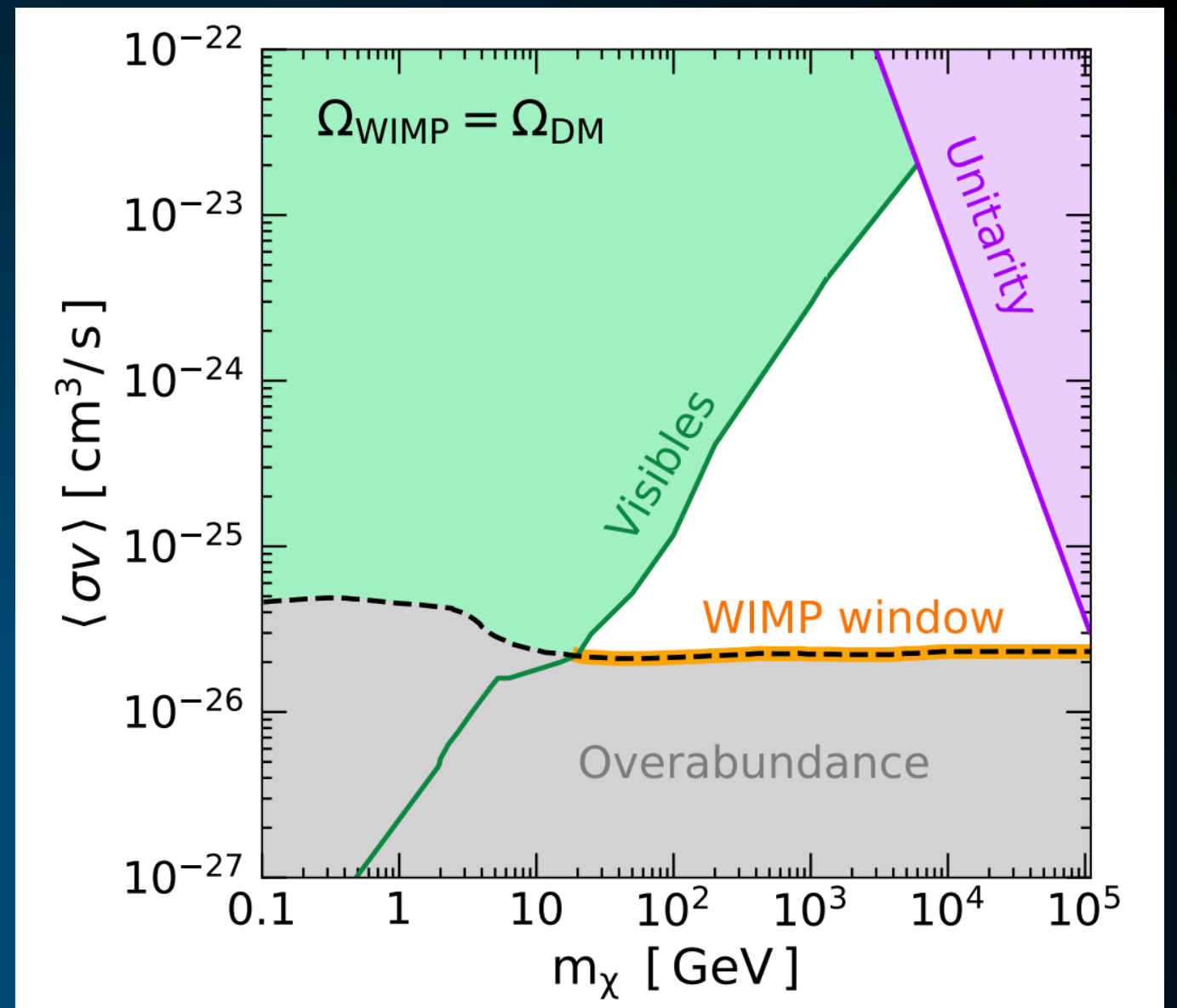
Calore et al. (2016; 1512.06825)

How to End a Talk When You Don't Have a Conclusion



Null Results are Interesting!

- **non-Thermal Dark Matter**
 - **Axions**
 - **Sterile Neutrinos**
- **Dark Sectors**
- **Modifications to Gravity**
- **Primordial Black Holes**
- **Q-balls**
- **Asymmetric Dark Matter**
- **Coannihilations, Resonances**
- **Self-Interacting Dark Matter, Light Mediators**



10^{-25} GeV

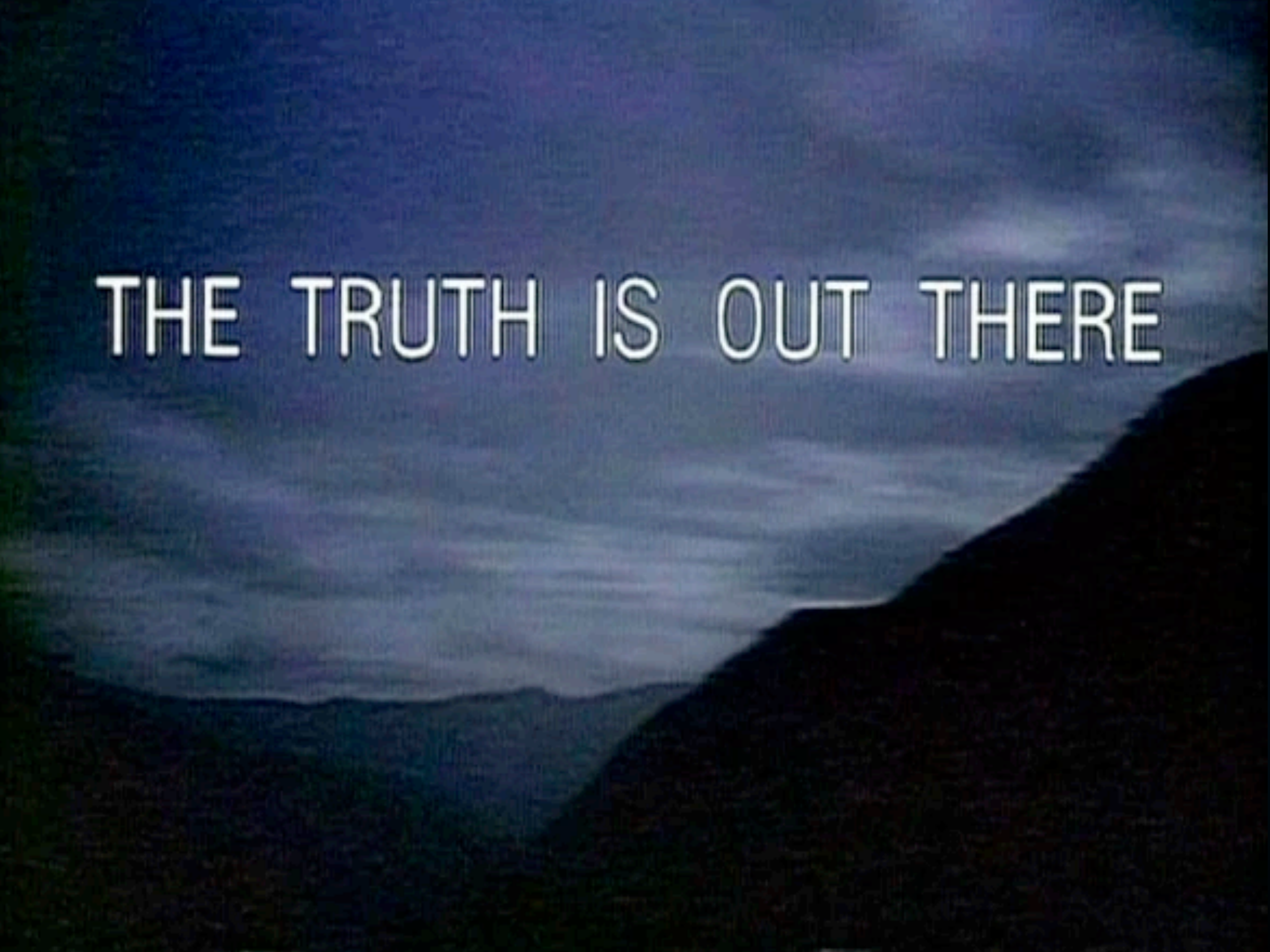
$\sigma_\chi > R_{\text{UFD}}$

BORING

10^{62} GeV

$m_\chi > M_{\text{UFD}}$

THE TRUTH IS OUT THERE

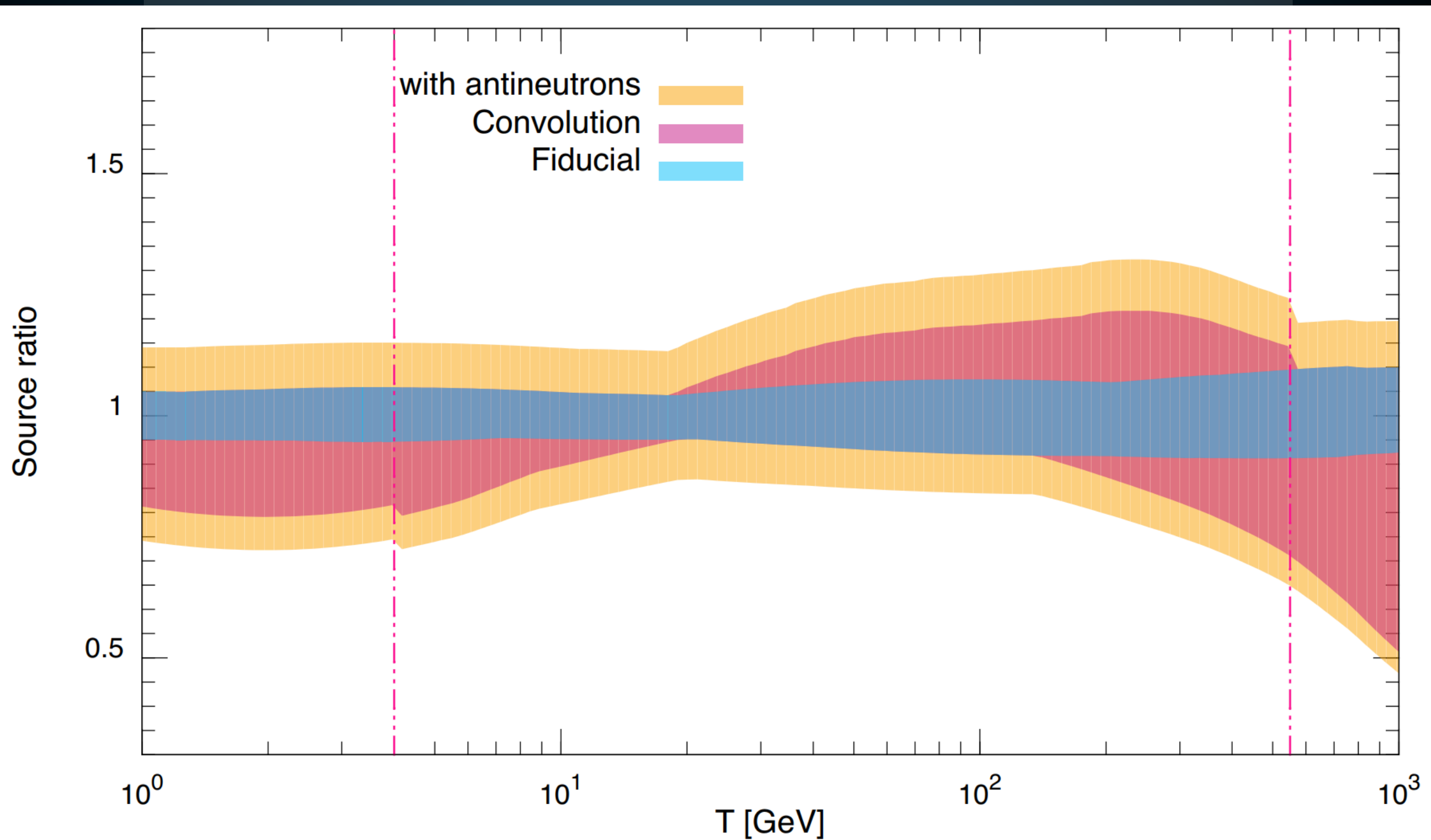


Conclusions



Extra Slides

Part 2: The Antiproton Excess



m_{DM} [GeV]

Spectral Features

Pretending We Understand How Diffusion Works

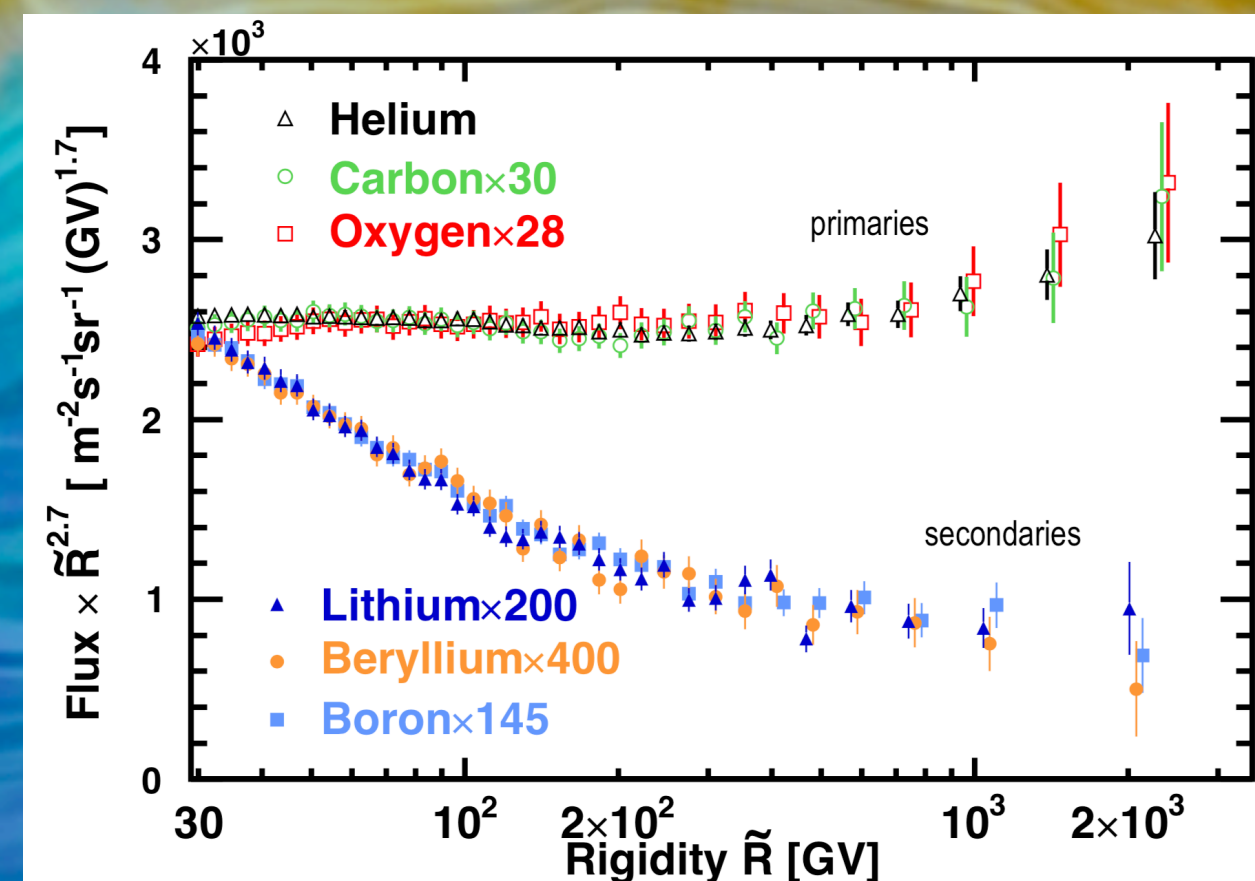
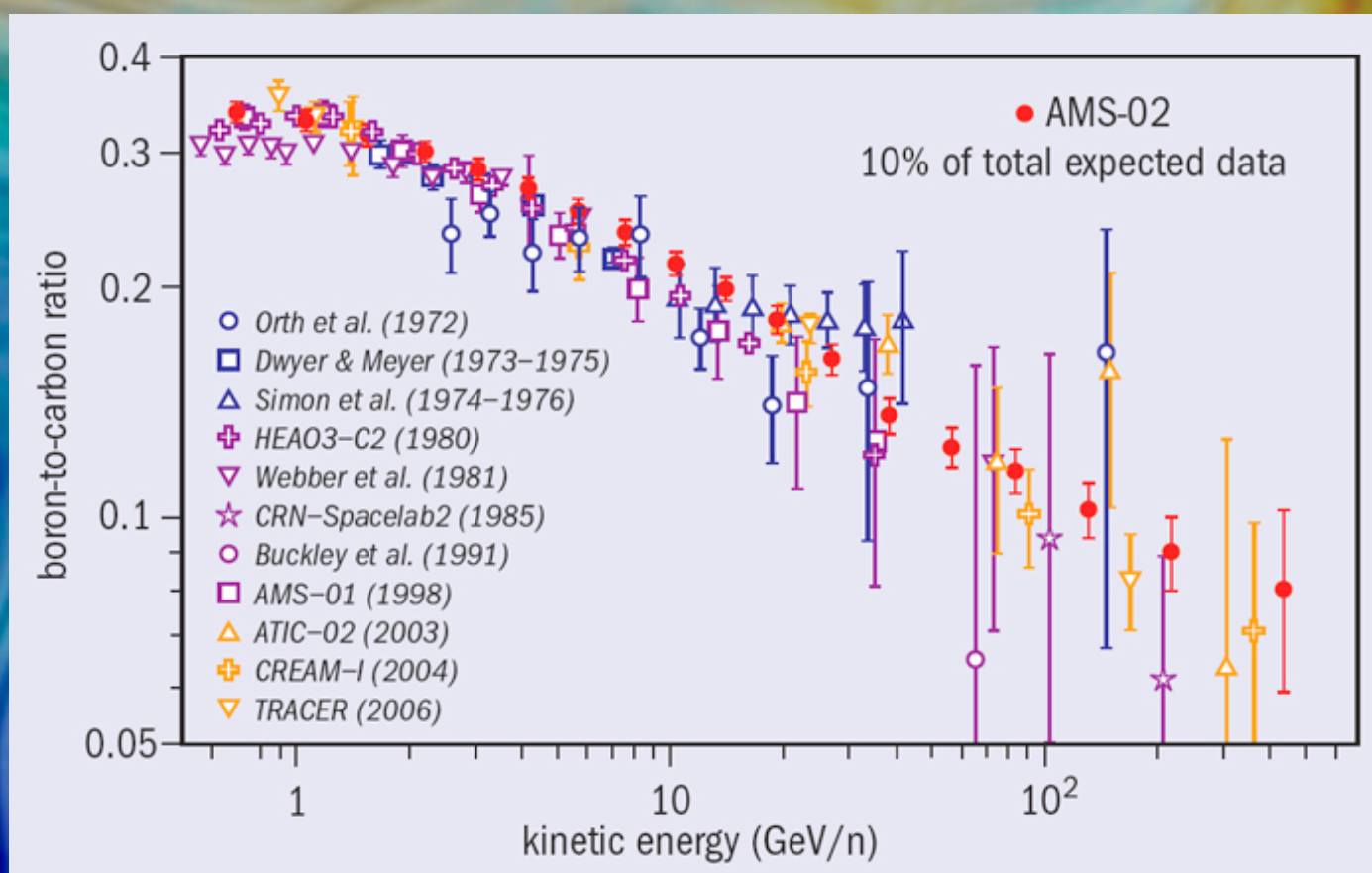
$$\frac{\partial f}{\partial t} = \underbrace{-V_{w,i} \frac{\partial f}{\partial x_i}}_{\text{advection}} + \underbrace{\frac{\partial}{\partial x_i} \kappa_{ij} \frac{\partial f}{\partial x_j}}_{\text{diffusion}} - \underbrace{V_{D,i} \frac{\partial f}{\partial x_i}}_{\text{drift}} + \underbrace{\frac{1}{3} \frac{\partial V_{w,i}}{\partial x_i} \frac{\partial f}{\partial \ln p}}_{\text{energy change}} + Q$$

advection

diffusion

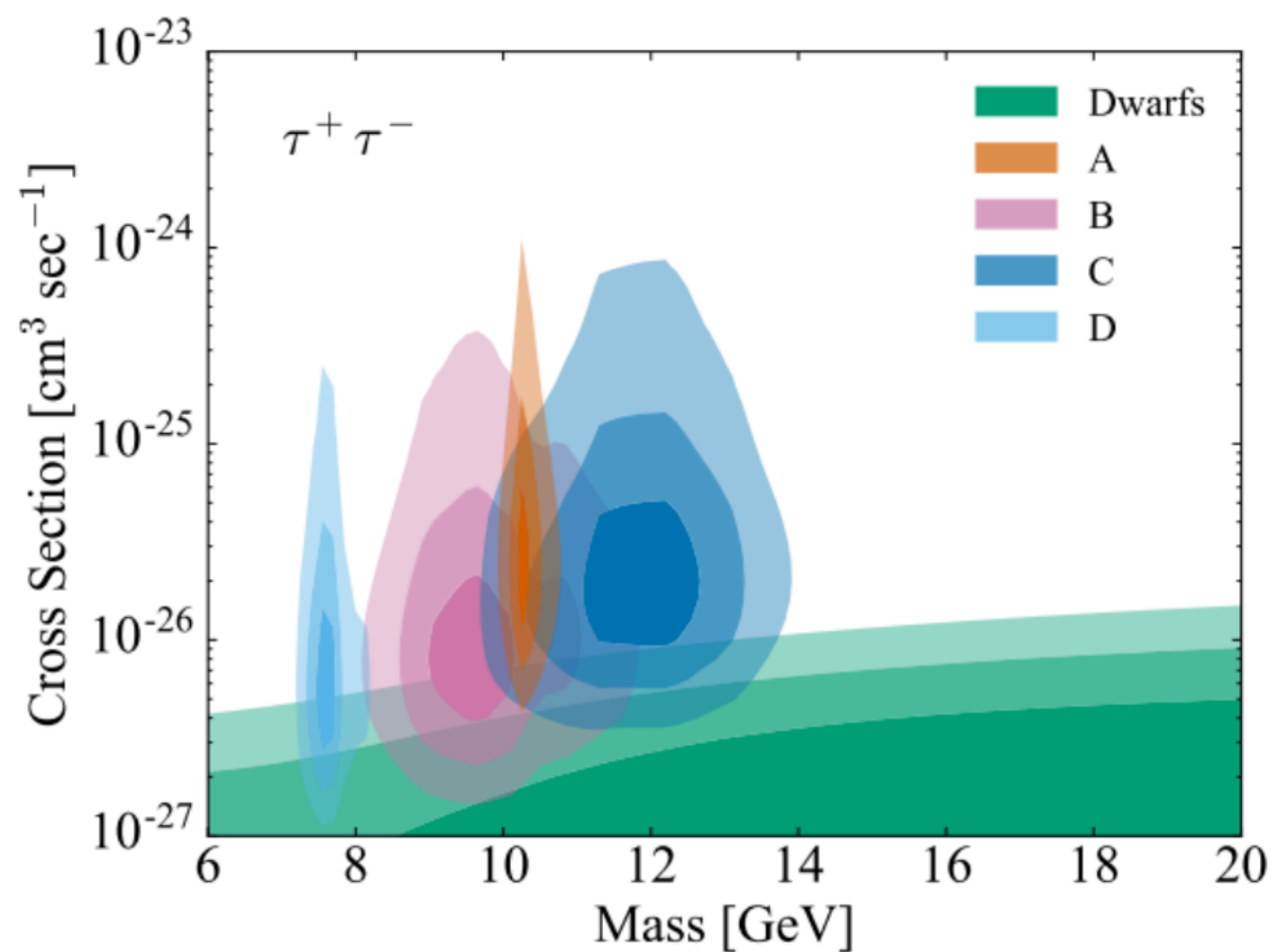
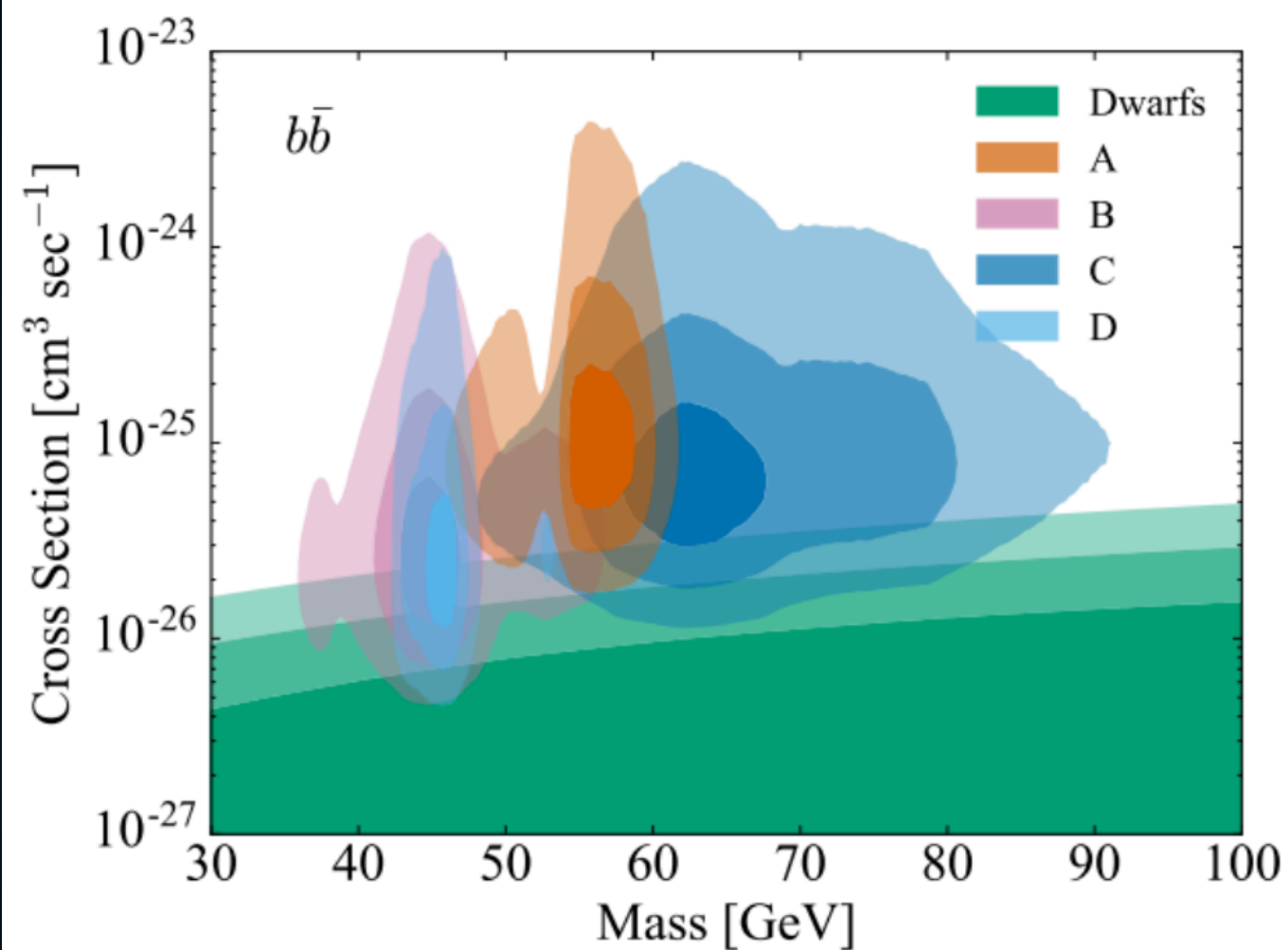
drift

energy change



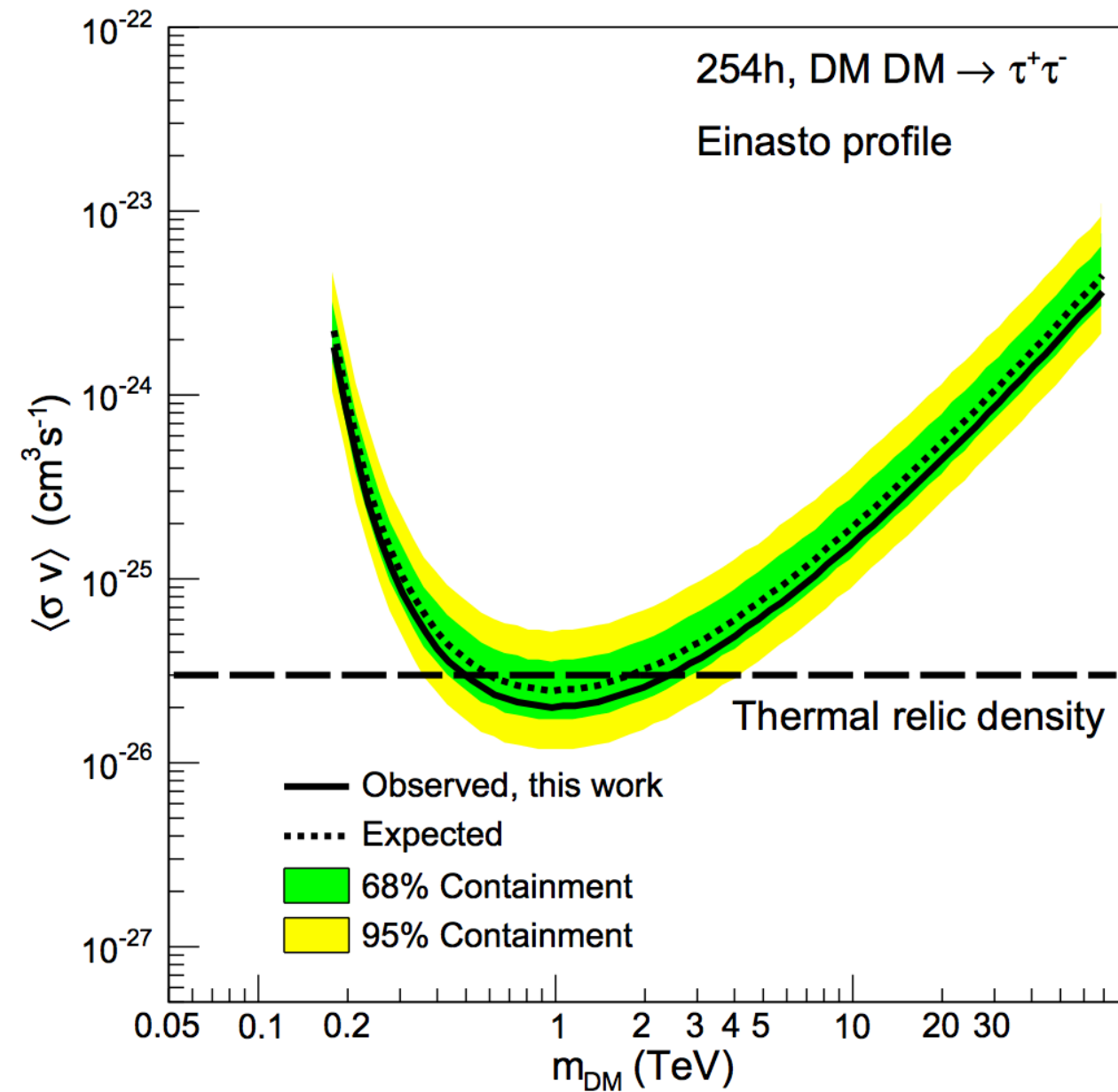
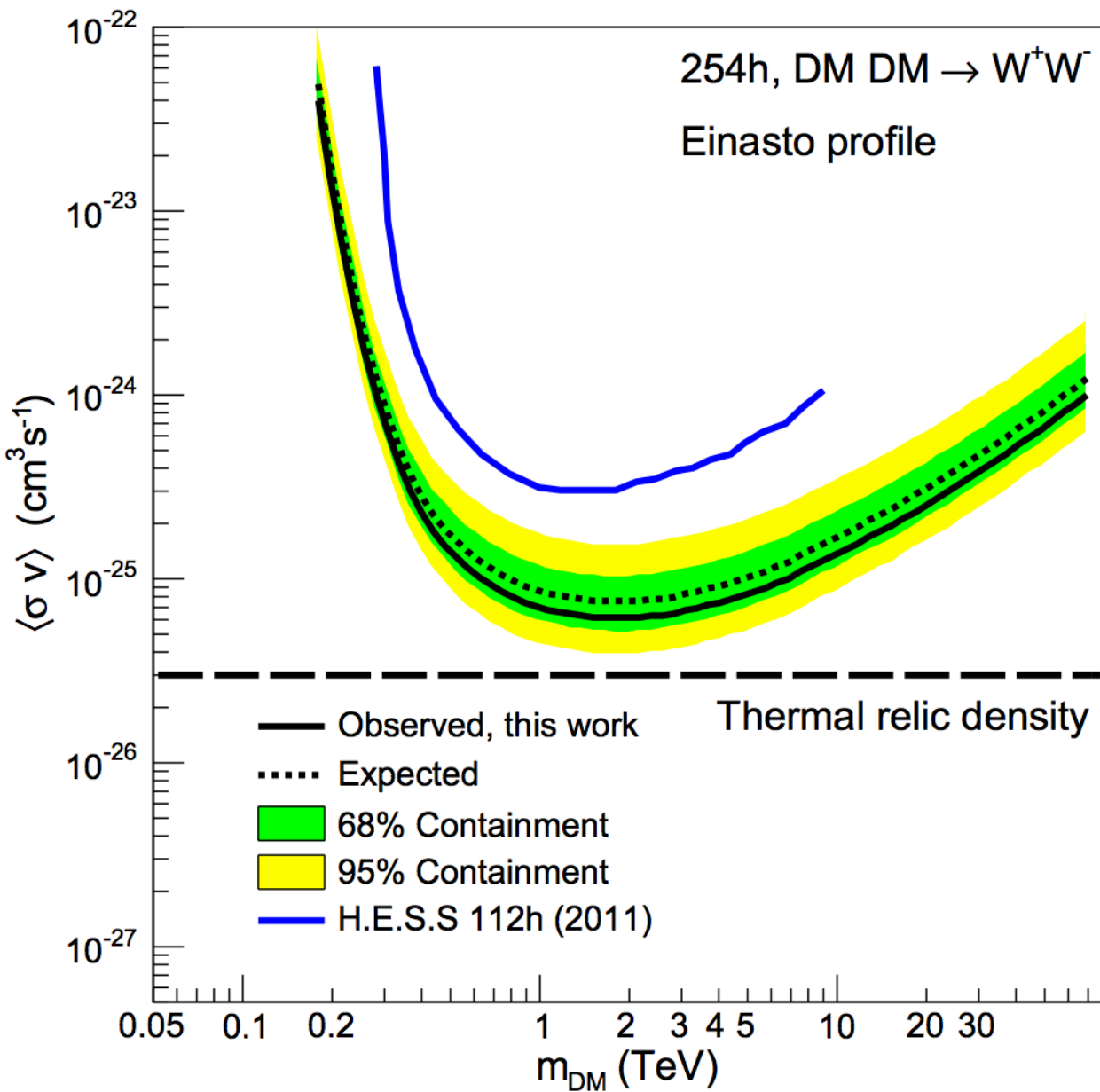
Are Dwarf Constraints in Tension with the Galactic Center Excess?

Keeley et al. (1710.03215)

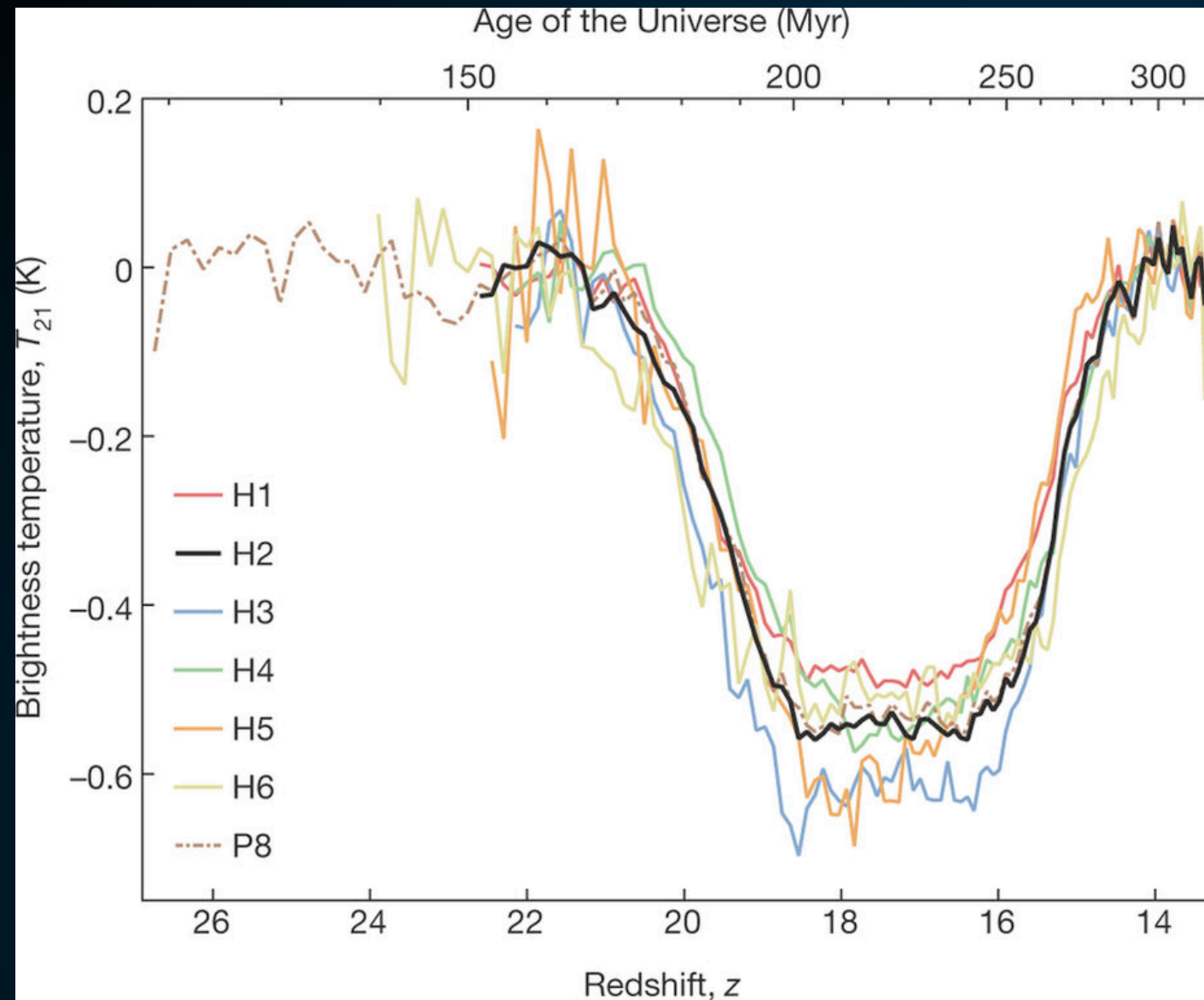


HESS Galactic Center Constraints

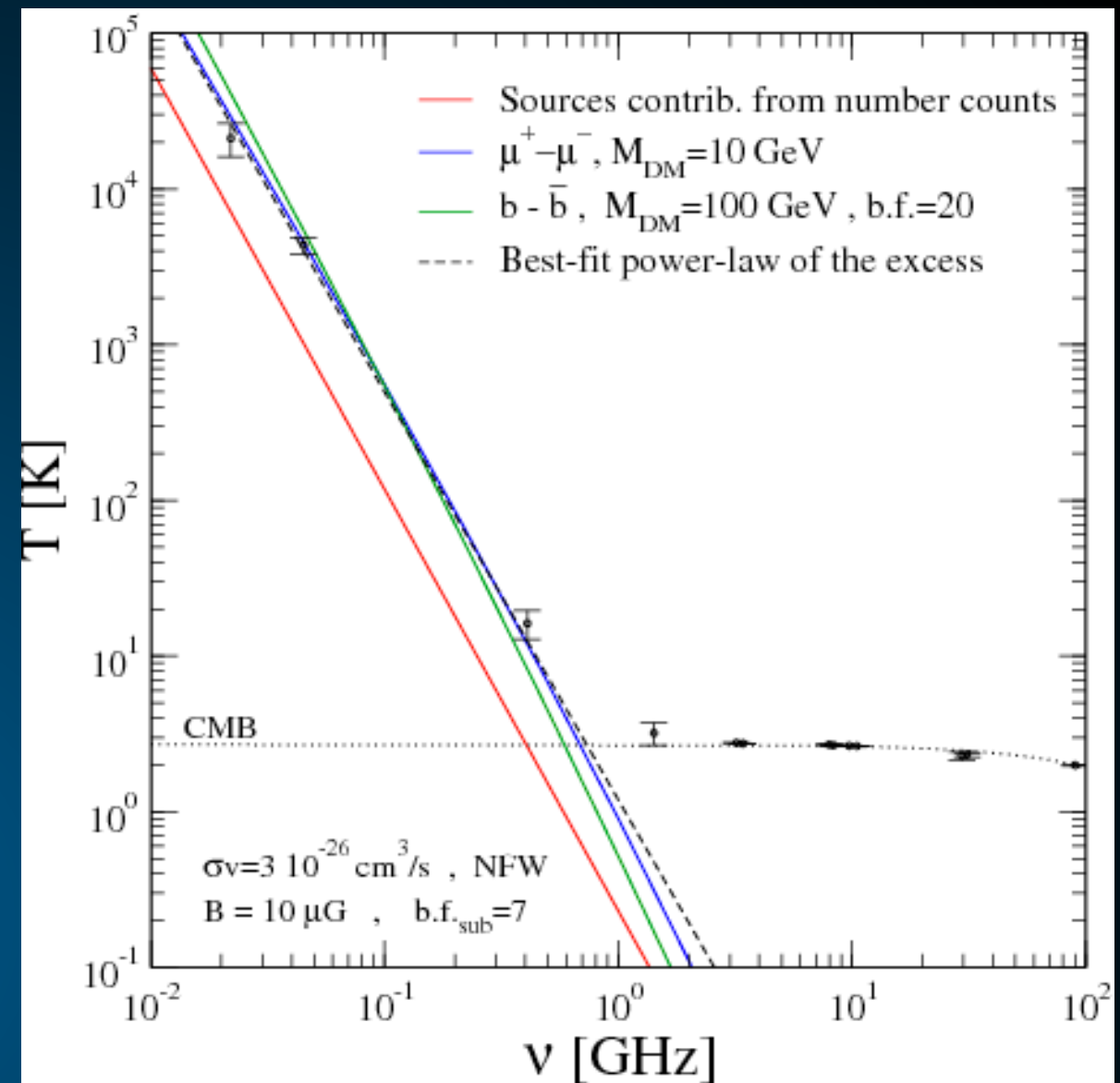
HESS Collaboration



EDGES/ARCADE Observations



EDGES - Strong constraints

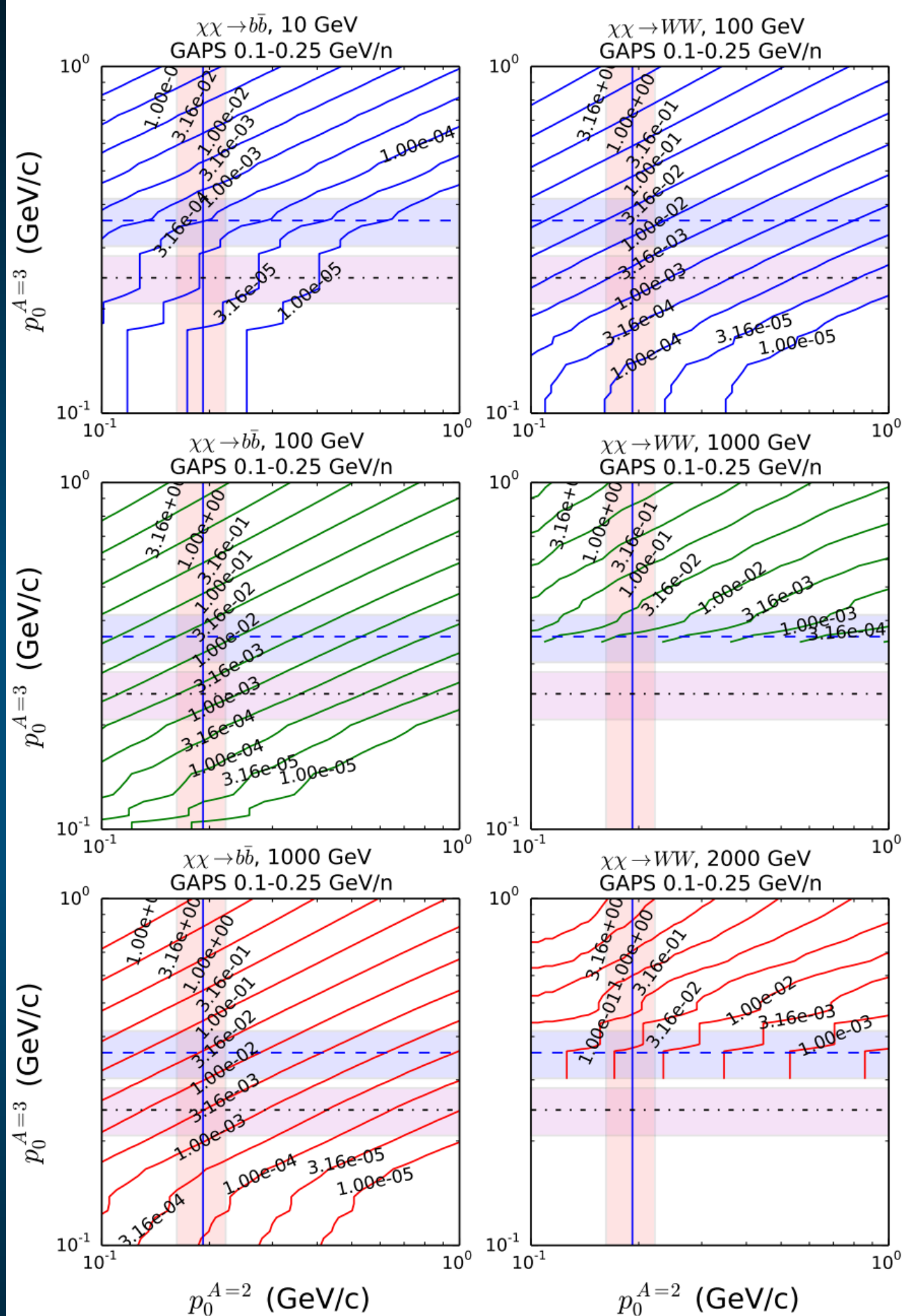


ARCADE - A Possible Signal?

Coalescence Models

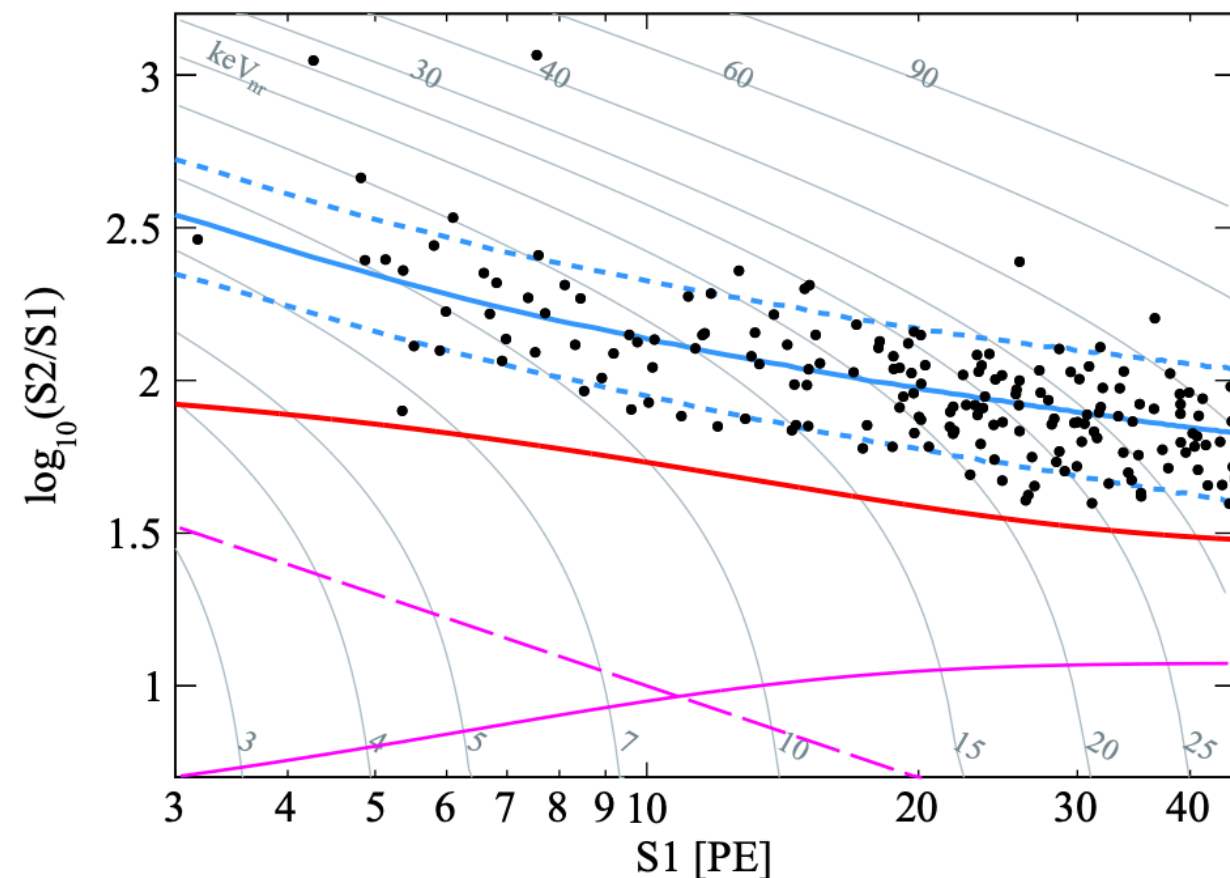
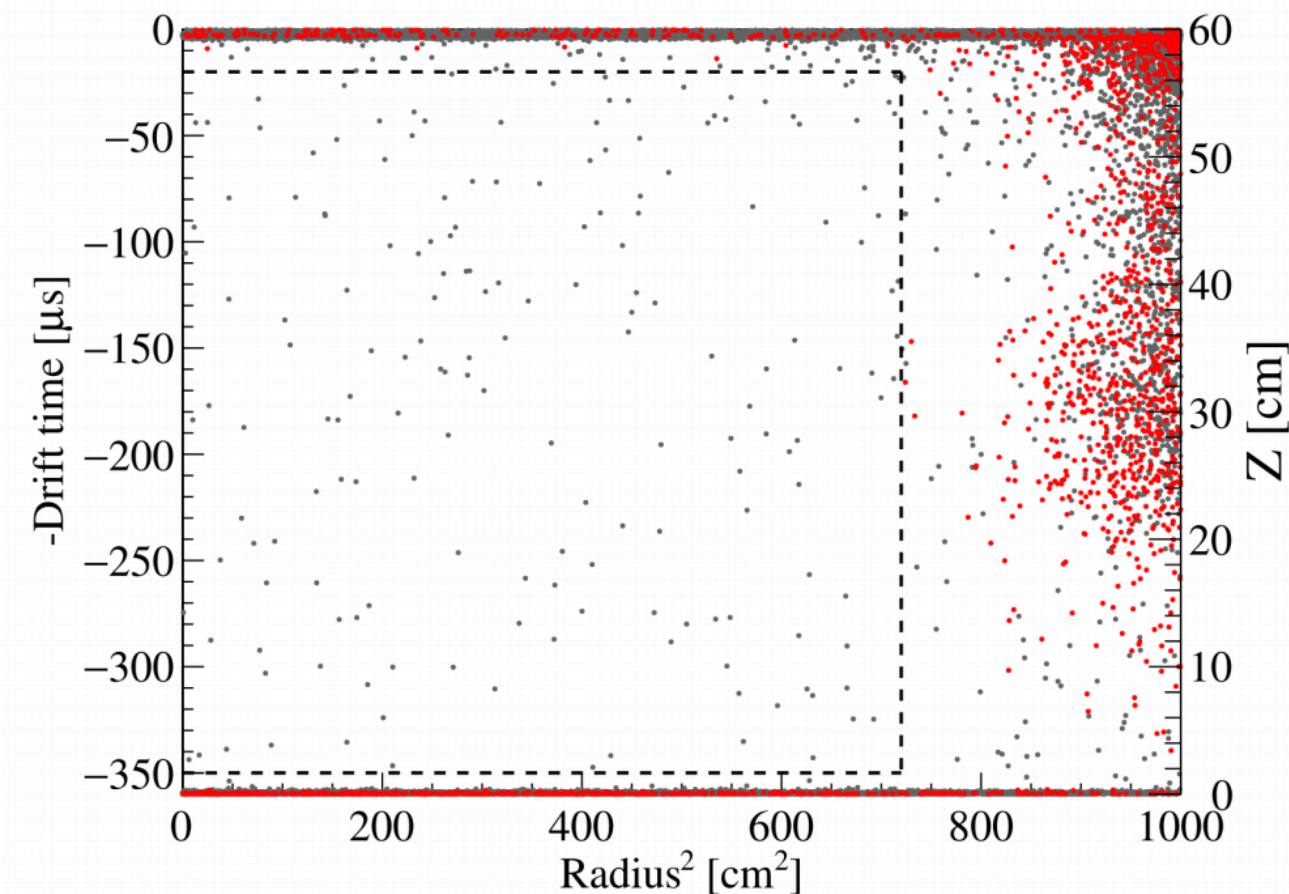
Formation of anti-particles determined from “coalescence momentum” that associates the binding energy of the atom with its formation probability.

Binding energy of anti-Helium > anti-deuterium, leads to larger formation probability.



Advantages of Direct Detection

Panda-XII Collaboration (1708.06917)



- Direct detection offers superb background rejection.
- Can set constraints on dark matter through the observation of 0 events.